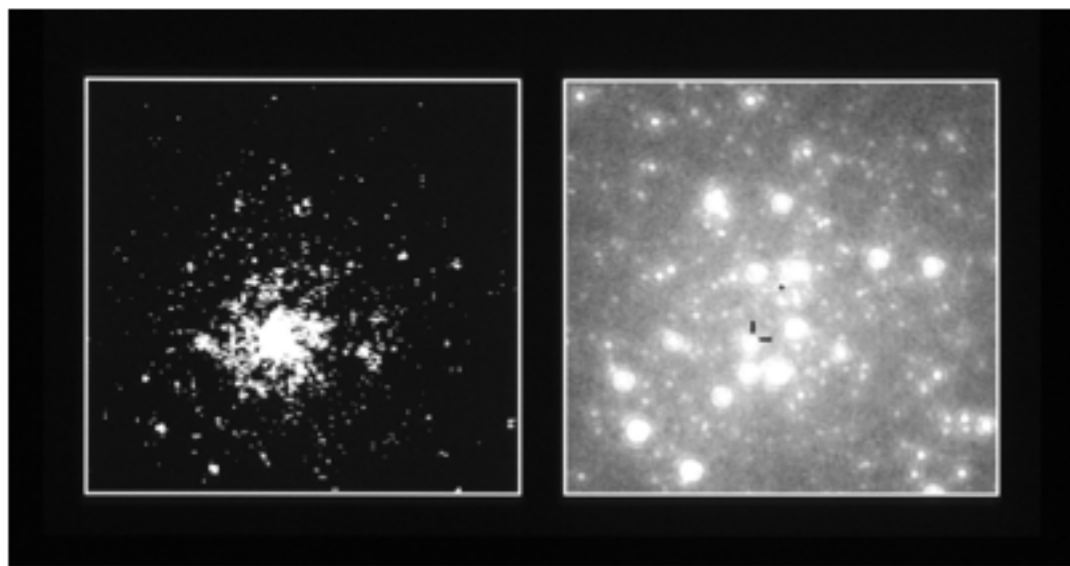


# Neutron Stars and the Equation of State of Dense Matter



HST observation of 4U1820

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March 8, 2012

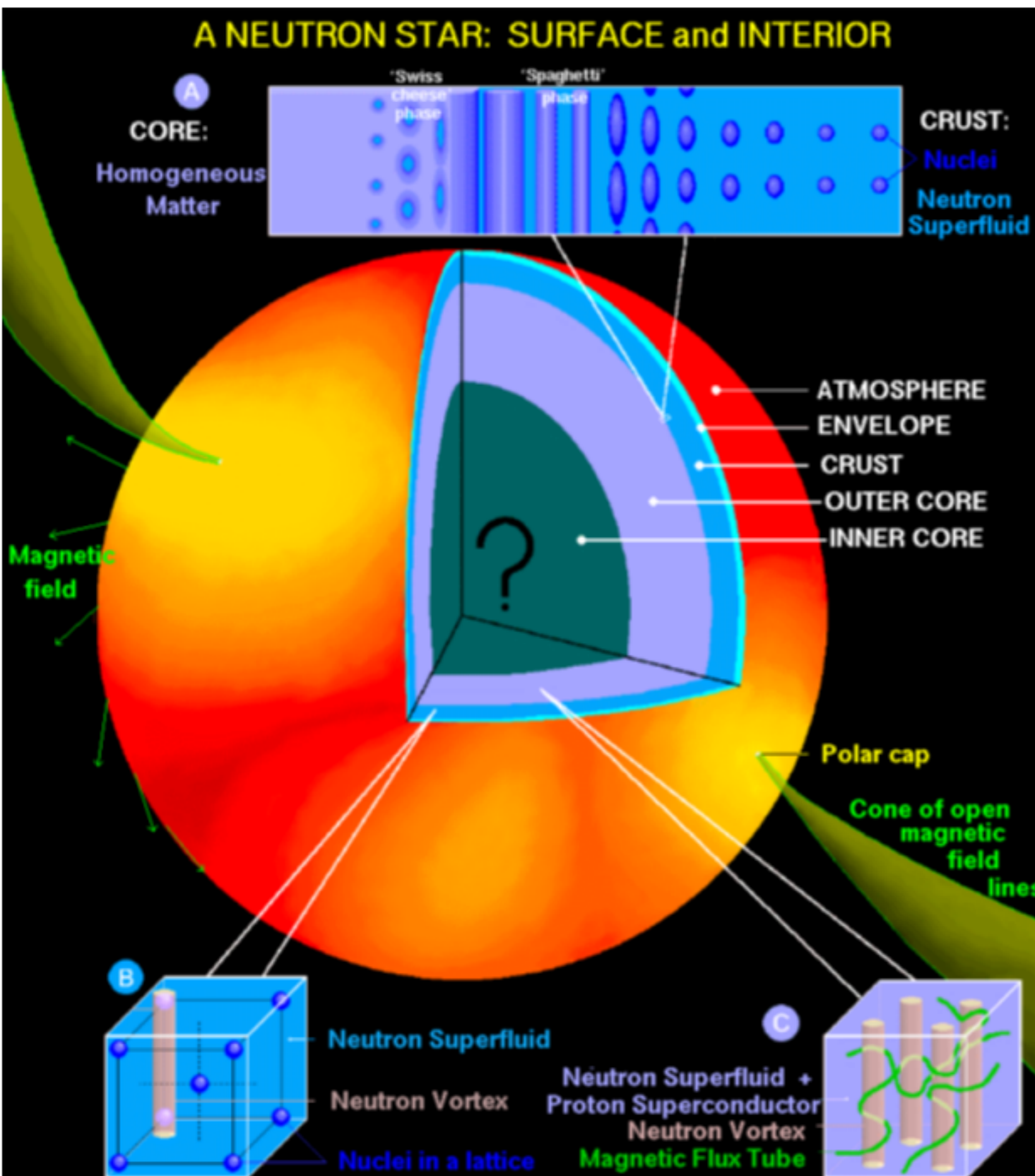
With: Edward F. Brown (Michigan State Univ.), Tobias Fischer (GSI)  
Stefano Gandolfi (Los Alamos), Matthias Hempel (Basel)  
James M. Lattimer (Stony Brook Univ.), Dany Page (UNAM),  
Madappa Prakash (Ohio Univ.)

# Outline

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- Fundamental questions about the equation of state of dense matter
  - What is the composition of neutron stars?
  - What is the nature of superfluidity at high densities?
  - What observables can help us determine the nature of dense matter?
- Fundamental neutron star questions are connected to fundamental nuclear physics questions, i.e. *What is the nuclear symmetry energy?*
- Background
- Mass and Radius Measurements
- Neutron Star Cooling
- Core-collapse Supernovae

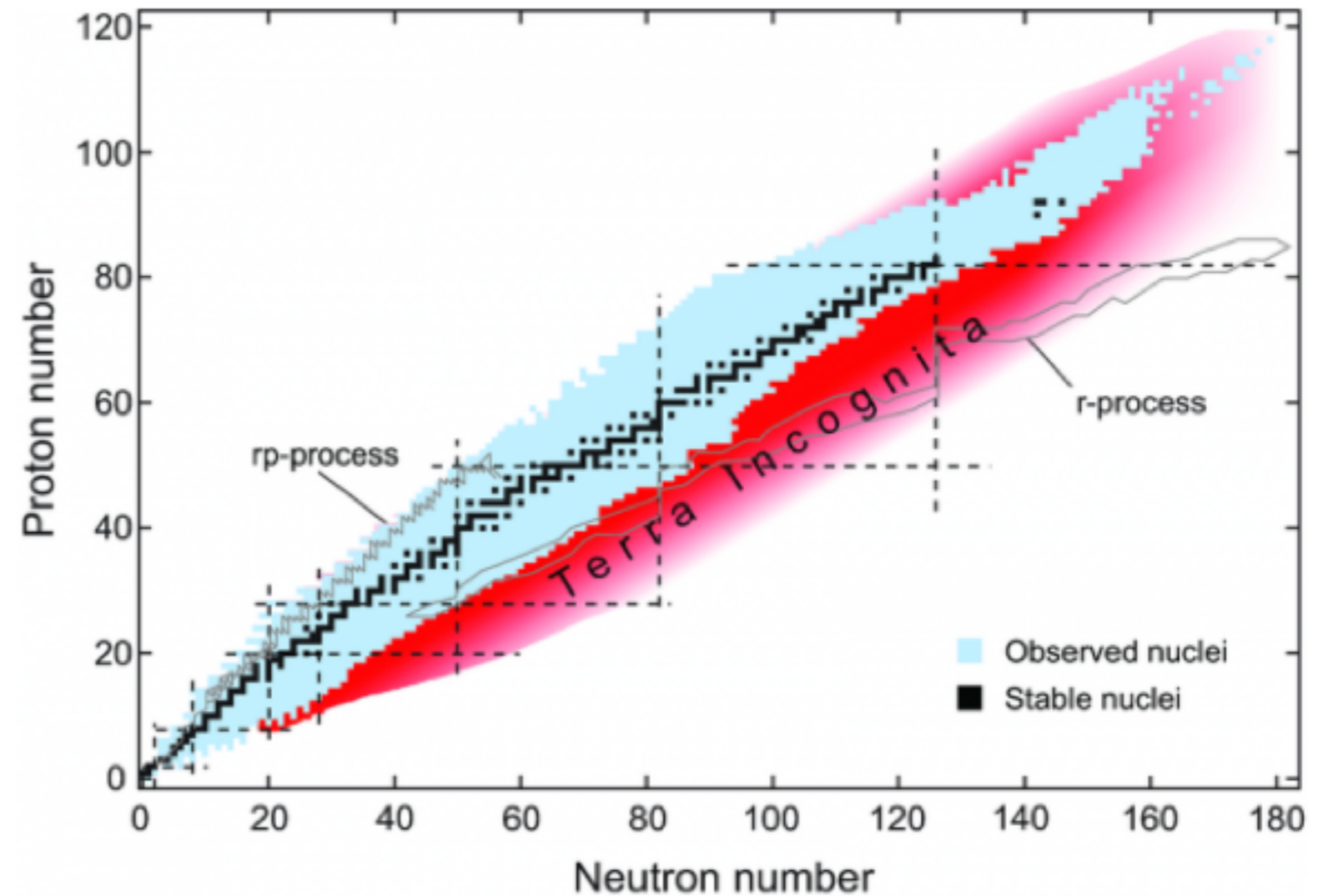
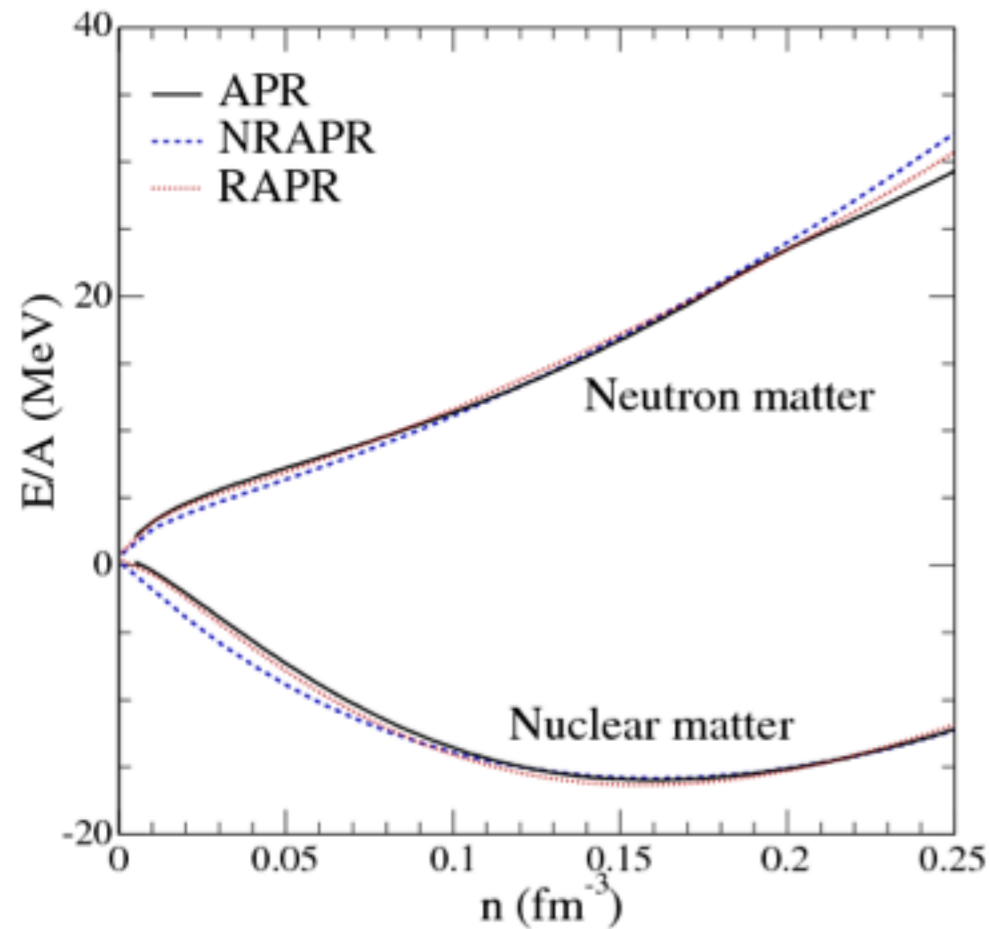
# Neutron Star Composition



- Crust is a lattice of neutron-rich nuclei
- Composition of the crust different in accreting systems
- Outer core is homogeneous nucleonic matter
- Inner core may contain phase transitions:  
 $[\Lambda, \Sigma, \Xi], [\pi, K], [u, d, s]$
- Superfluidity is important

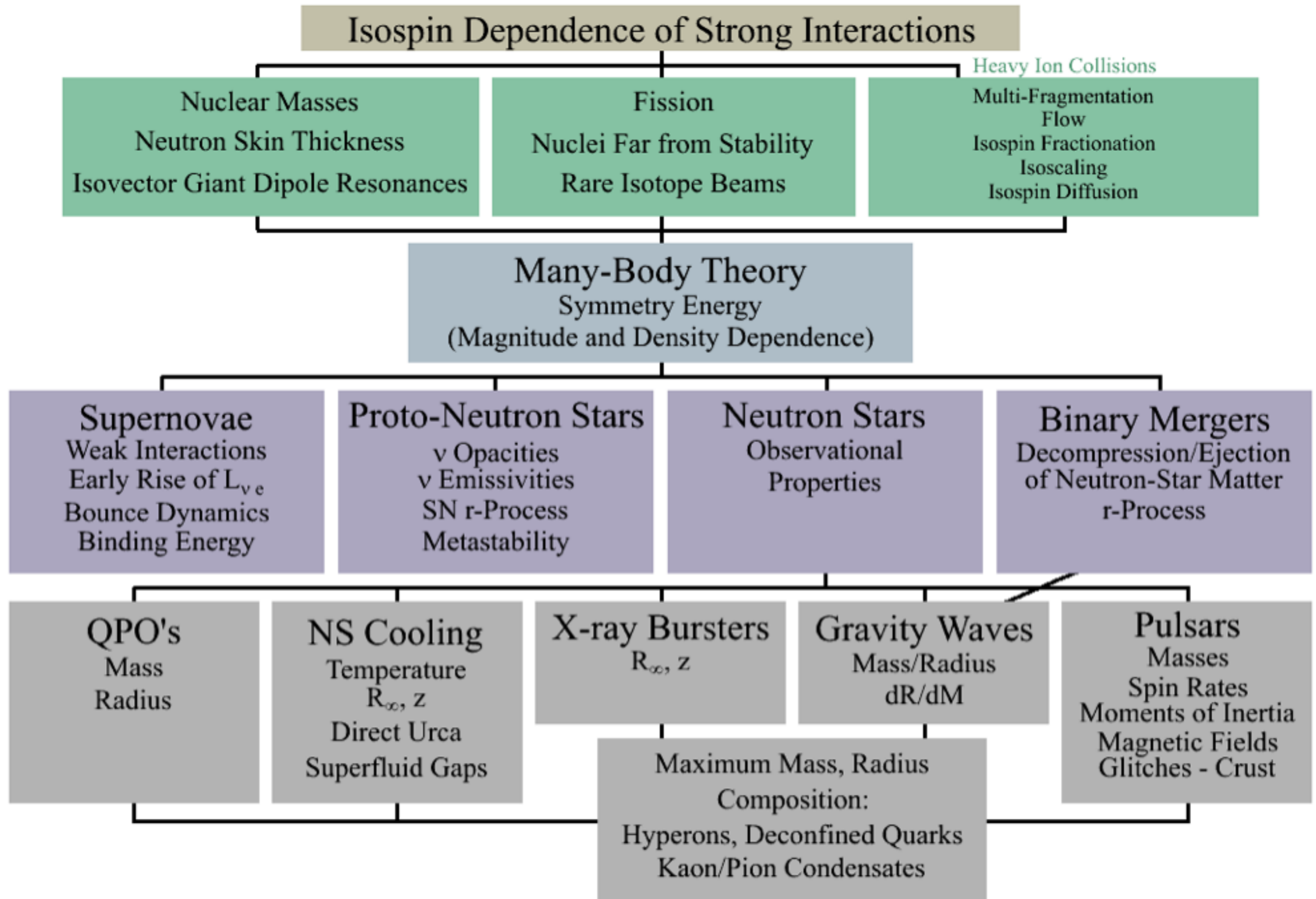
Figure by Dany Page

# The Nuclear Symmetry Energy

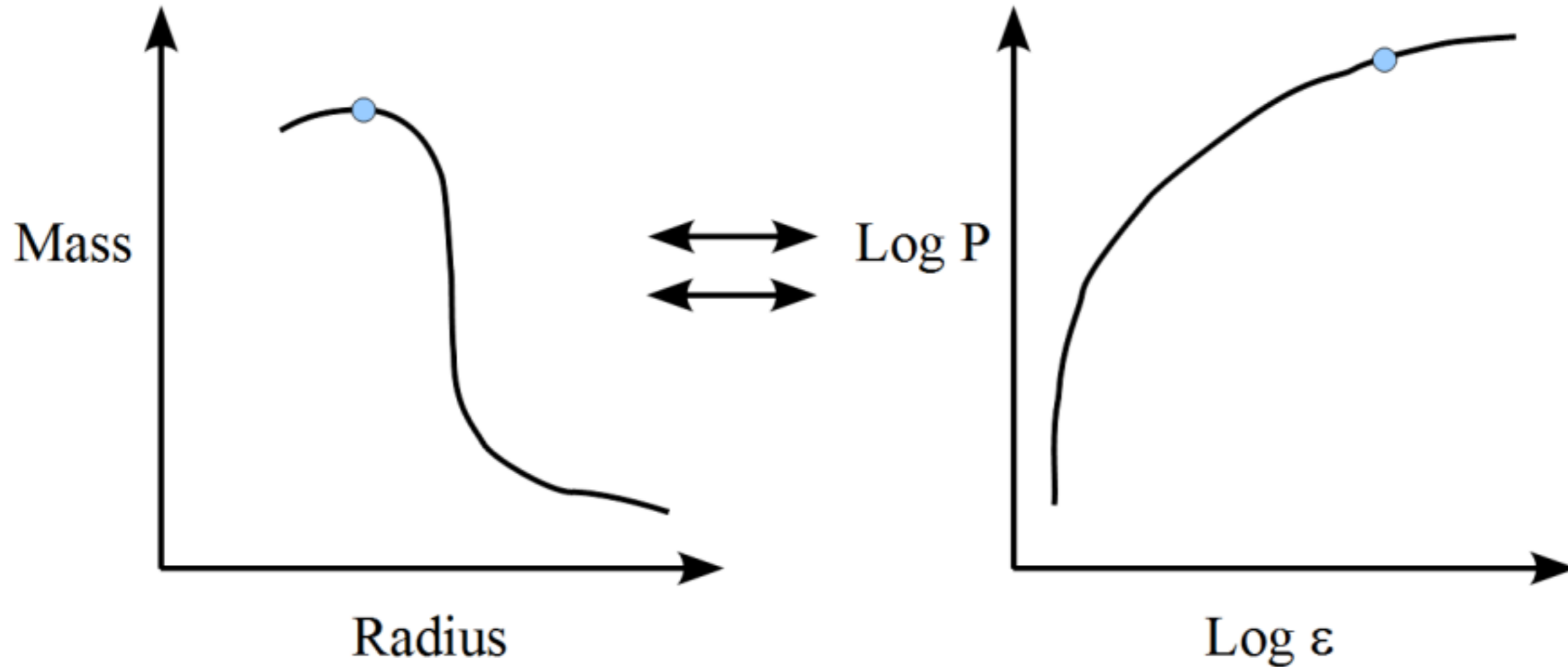


- The symmetry energy is the energy cost to create an isospin asymmetry
- The origin of the 'valley of stability'
- One of the largest uncertainties in the nucleon-nucleon interaction
- $S$  is the value at the nuclear saturation density  $S = S(n_0)$
- $L$  is the derivative,  $L = 3n_0 S'(n_0)$

# Connections to the Symmetry Energy



# M vs. R and the EOS of Dense Matter



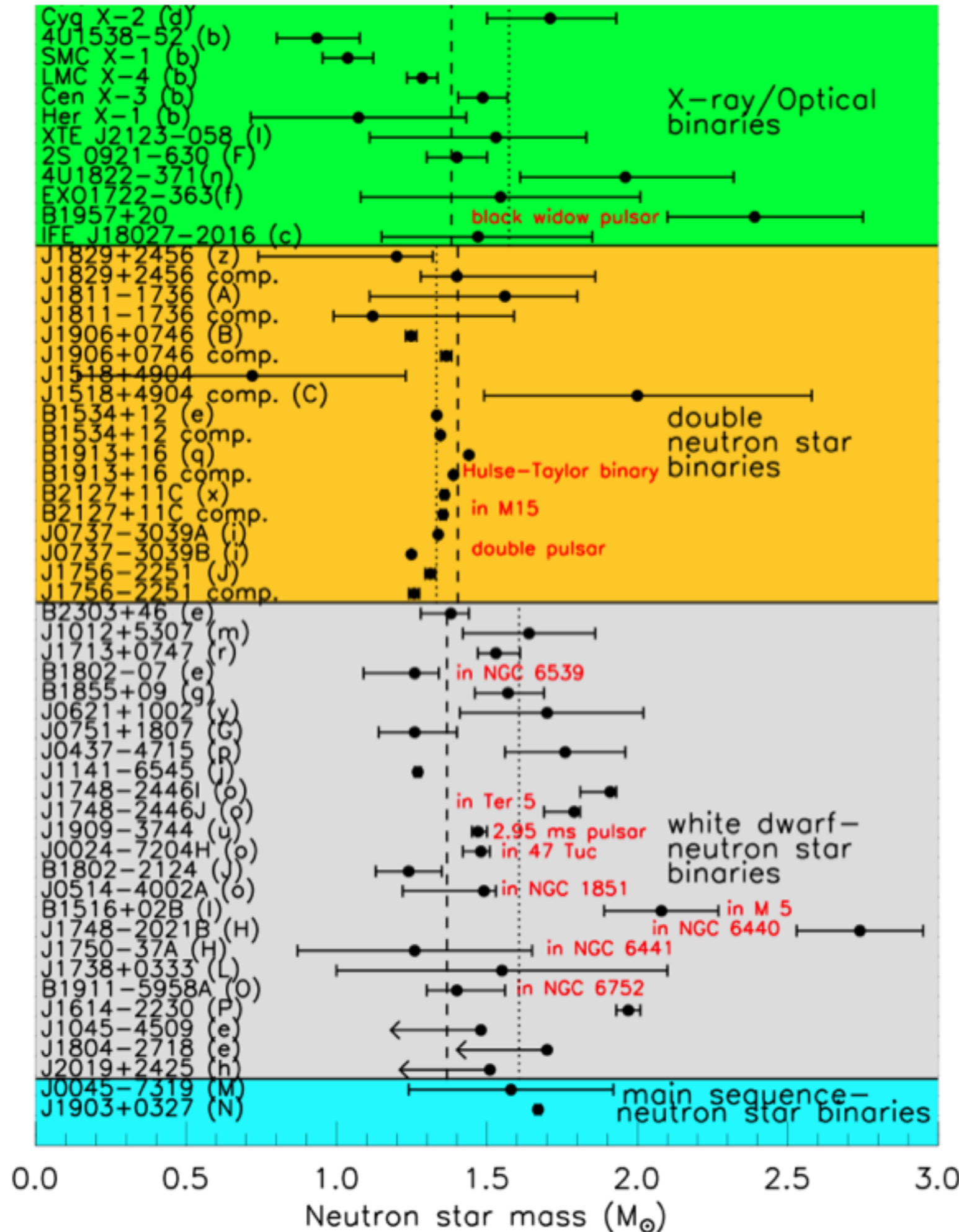
## As of 5 years ago:

- Accurate mass measurements from double pulsars (e.g. Hulse-Taylor pulsar)
- Limited radius information for a few sources (e.g. Rutledge et al.)
- A few limited constraints from pulsar spins and pulsar glitches

## Now:

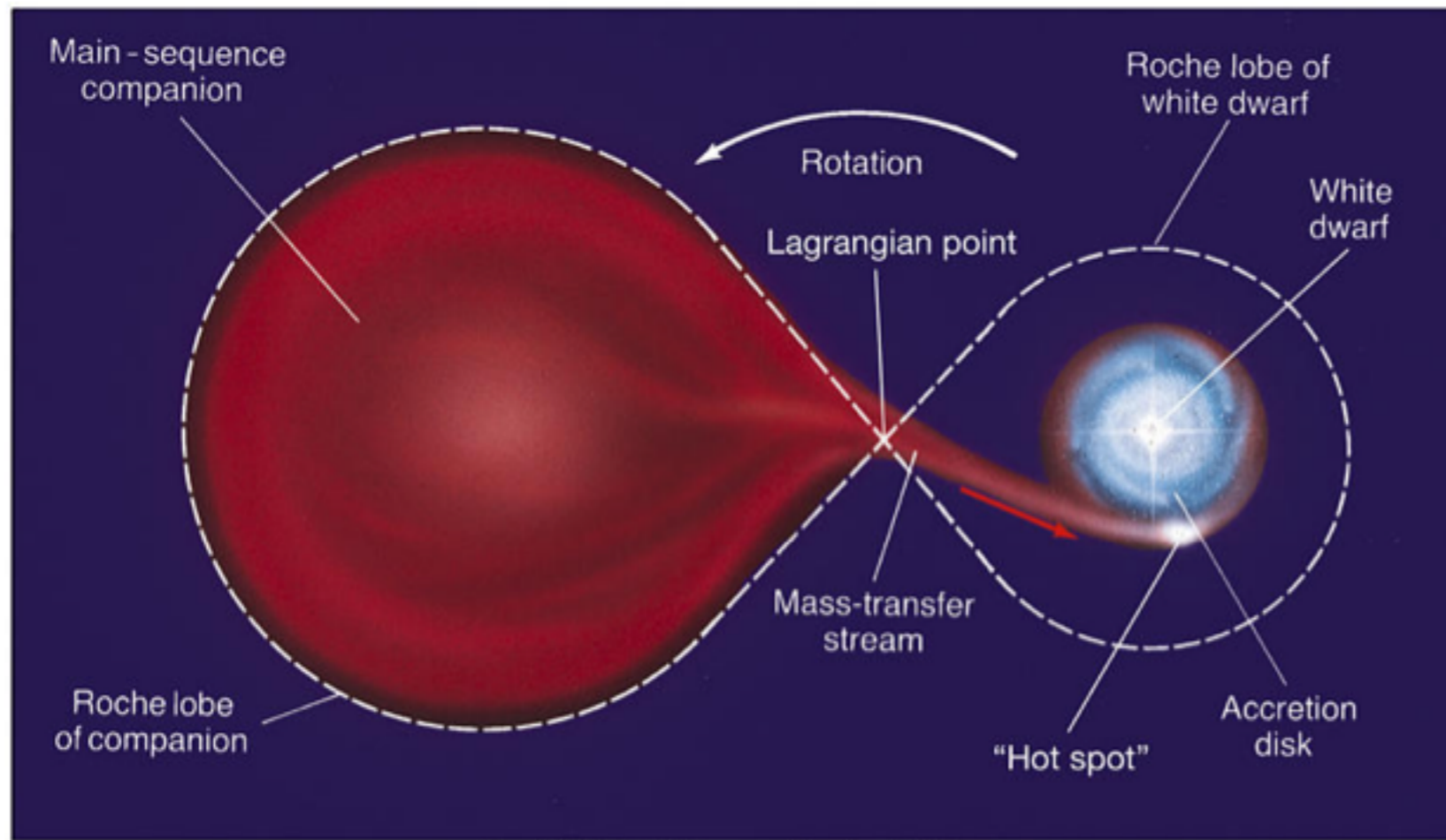
- 10-15 percent measurements of M and R for the same object
- A 2 solar mass neutron star (Demorest et al. 2010)

# Neutron Star Mass Measurements

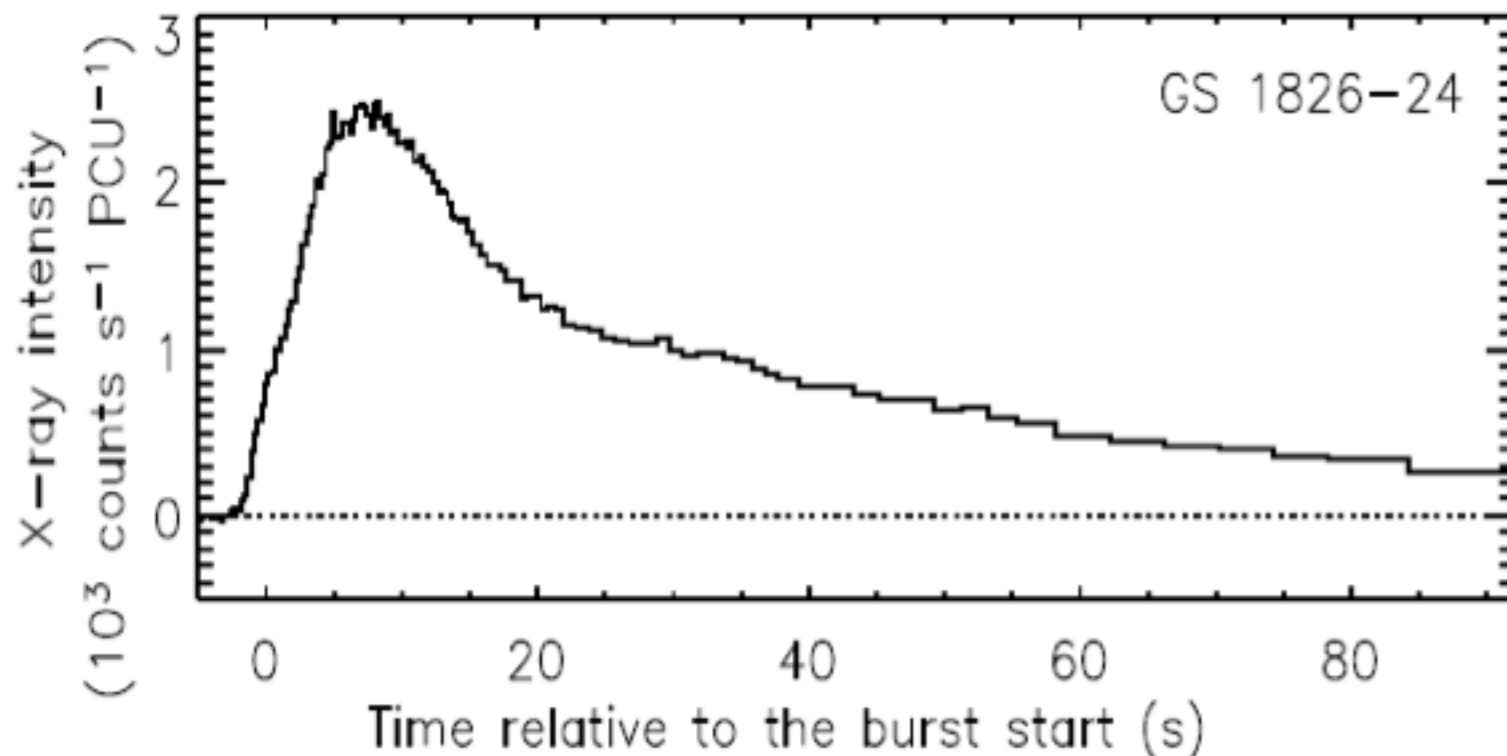


- Neutron star mass distribution is an important unknown quantity
- Friere et al. (2008):  $2.74 \pm 0.21$ , but depends on assumptions about the inclination angle
- Demorest et al. (2010) measurement:  $1.97 \pm 0.04 M_{\odot}$

# Accreting Neutron Stars: LMXBs



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- From a main-sequence (normal) star or a white dwarf
- Overflowing the Roche lobe
- Most often accrete a mix of hydrogen and helium, sometimes heavier elements
- At high enough density, light elements are unstable to thermonuclear explosions



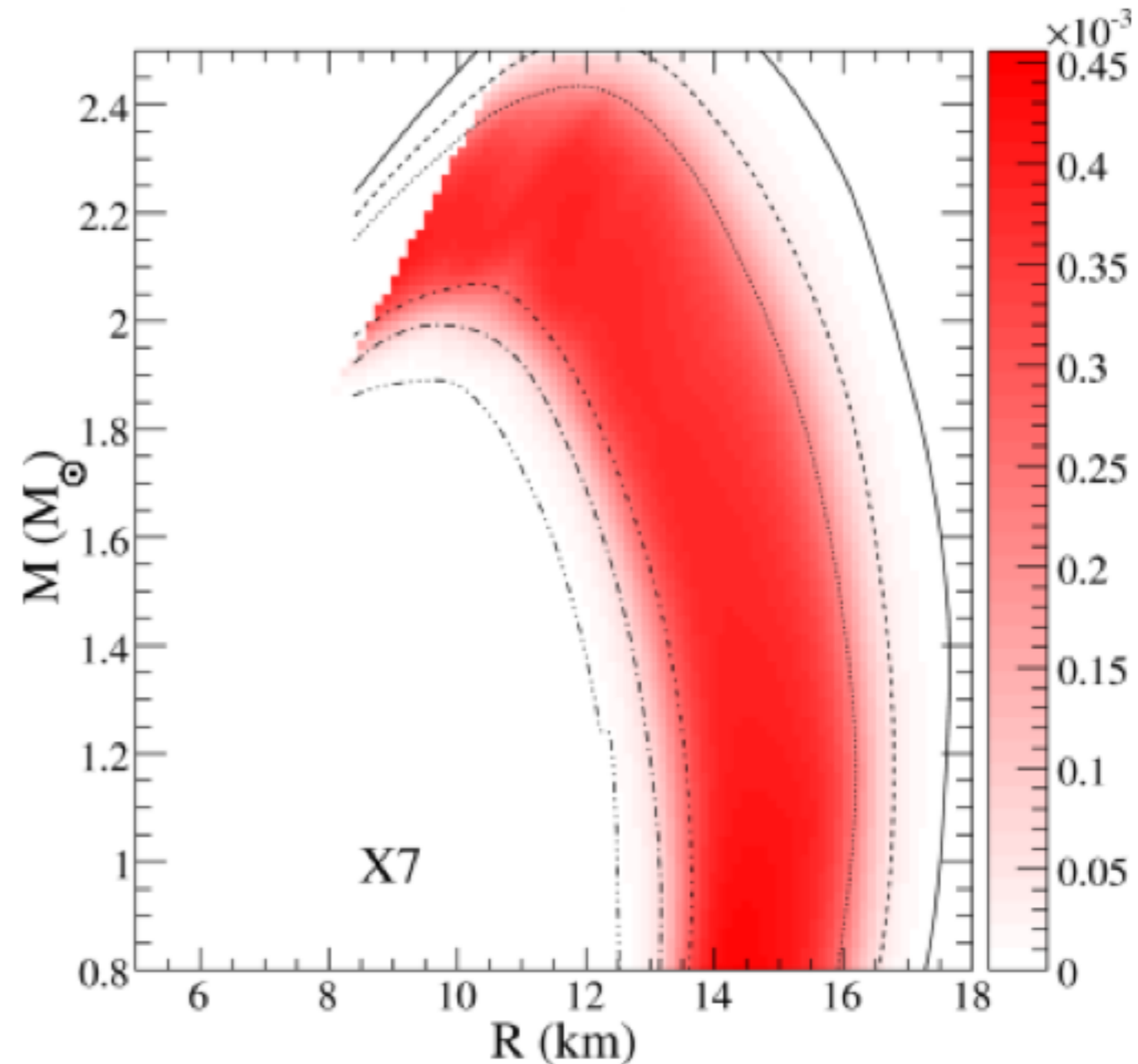
# Mass Measurements and QLMXBs

- *Quiescent LMXBs in globular clusters:*

- H atmosphere
- Known distance
- Small magnetic field
- Measure radius:

$$F \propto T_{\text{eff}}^4 \left( \frac{R_{\infty}}{D} \right)^2$$

[Rutledge et al. (1999), Heinke et al. 2006, Webb and Barrett (2007), Guillot et al. (2010)]



Steiner et al. (2010)

# Photospheric Radius Expansion Bursts

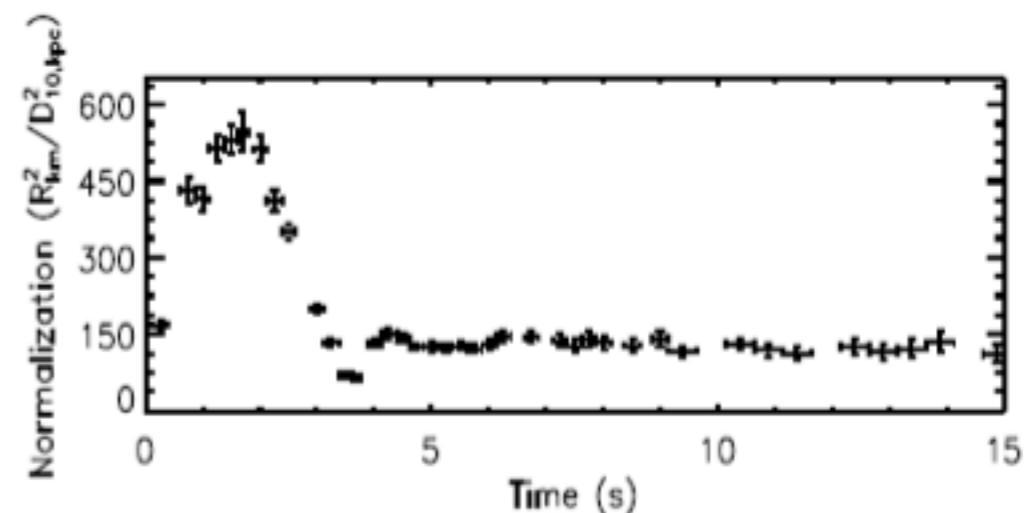
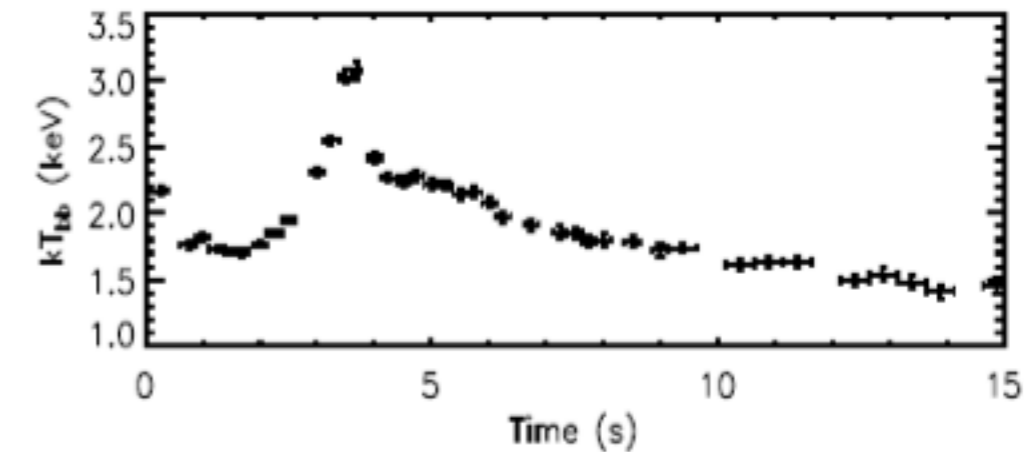
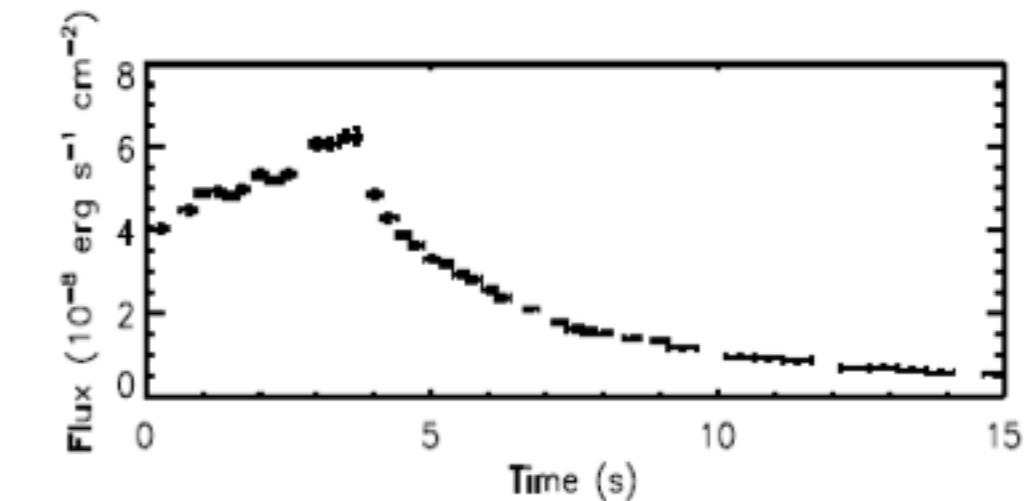
- X-ray bursts sufficiently strong to blow off the outer layers - radiate at the Eddington limit
- Flux peaks, then temperature reaches a maximum, "touchdown"

$$F_{TD} = \frac{GMc}{\kappa D^2} \sqrt{1 - 2\beta(r_{ph})}$$

- Normalization during the tail of the burst:

$$A \equiv \frac{F_{\infty}}{\sigma T_{bb,\infty}^4} = f_c^{-4} \left( \frac{R}{D} \right)^2 (1 - 2\beta)^{-1}$$

- If we have the distance, two constraints for mass and radius



Ozel et al. (2010)

(Pioneered by van Paradijs et al. 1979 and 1982 and Ebisuzuki et al. 1983)

# Photospheric Radius Expansion Bursts

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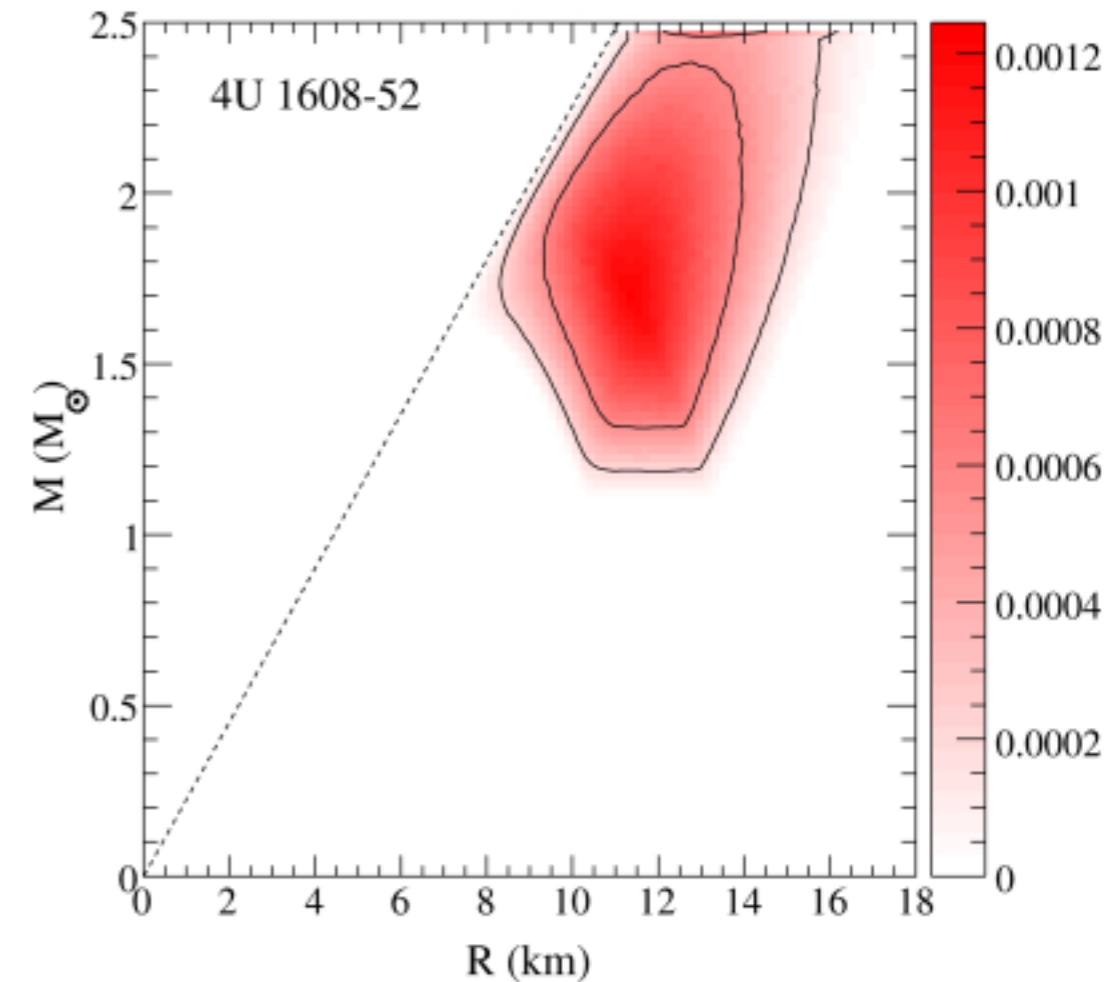
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- If we have the distance, two constraints for mass and radius
- Dimensionless parameter

$$\alpha \equiv \frac{F_{TD}\kappa D}{\sqrt{A} c^3 f_c^2}$$



Steiner et al. (2010)

# Statistical Approach

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- Well-suited to this underconstrained problem: 7-8 EOS parameters, 7-8 data points
- Bayes theorem:

$$P[\mathcal{M}_i|D] = \frac{P[D|\mathcal{M}_i]P[M_i]}{\sum_j P[D|\mathcal{M}_j]P[\mathcal{M}_j]}$$

- Different prior distributions produce different results
- Conditional probability is provided by the data

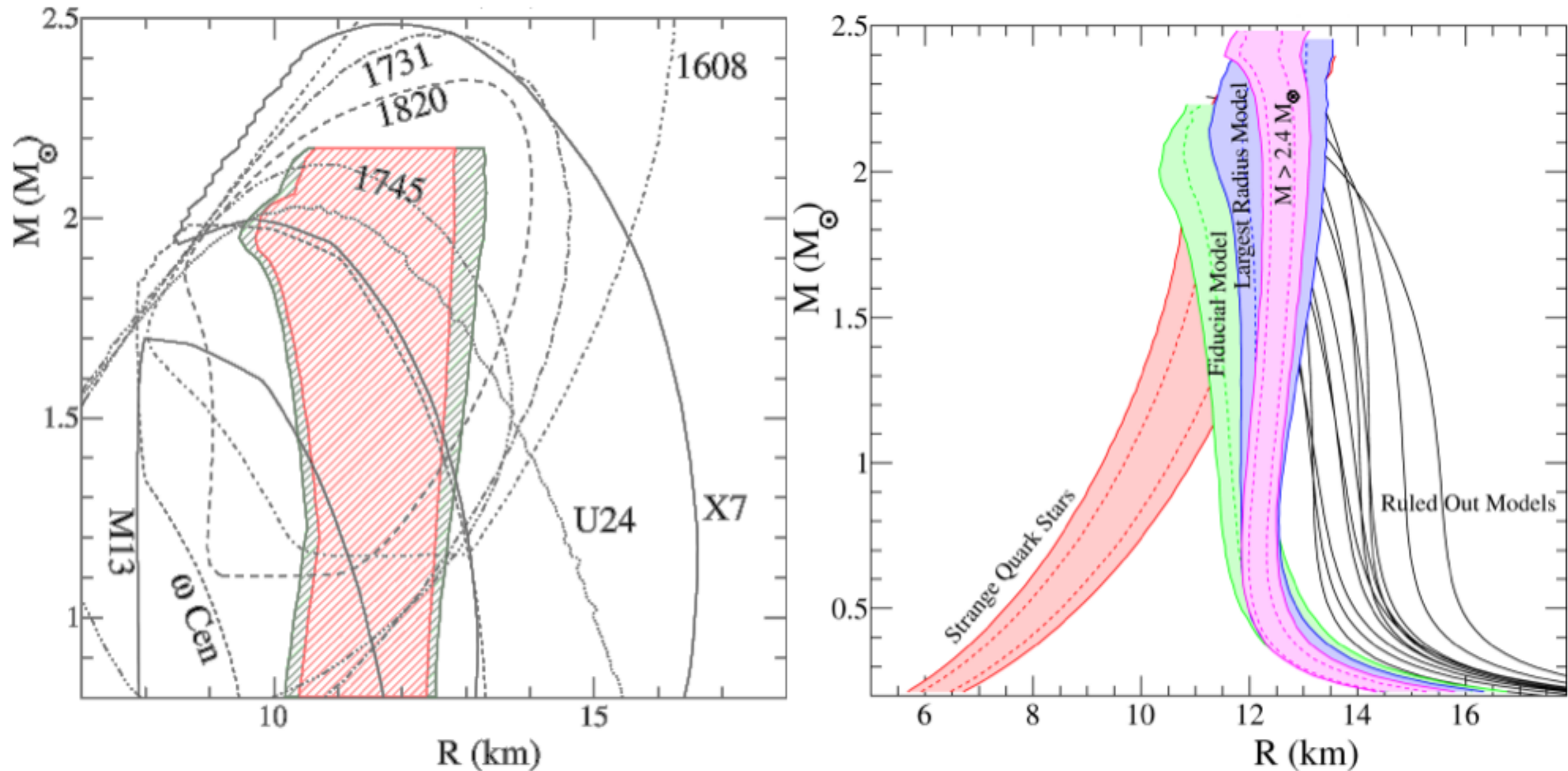
$$P[D|\mathcal{M}] = \prod_{i \in n_{\text{datasets}}} \mathcal{D}_i(M, R)|_{M=M_i, R=R(M_i)}$$

the analog of the likelihood function

- In Bayesian analysis, marginal estimation is often employed:

$$P[p_j|D](p_j) = \frac{1}{V} \int dp_1 \dots dp_{j-1} dp_{j+1} \dots dp_{N(p)} P[M|D]$$

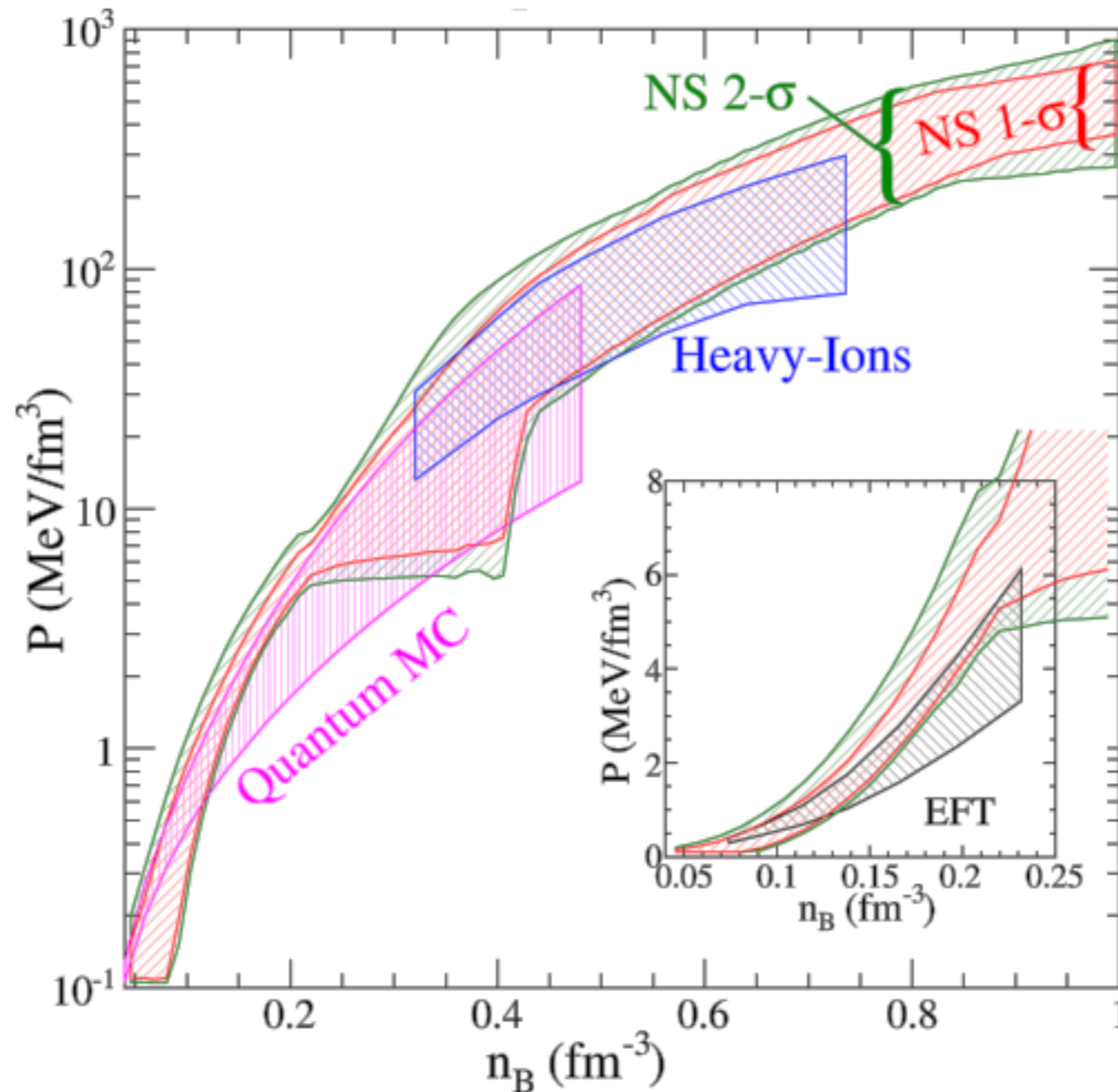
# Mass and Radius Results



Steiner, Lattimer, and Brown, in prep.

- Choose the largest range which encloses several choices in model assumptions and prior distributions
- Range of radii for a 1.4 solar mass star: 10.4 and 12.9 km

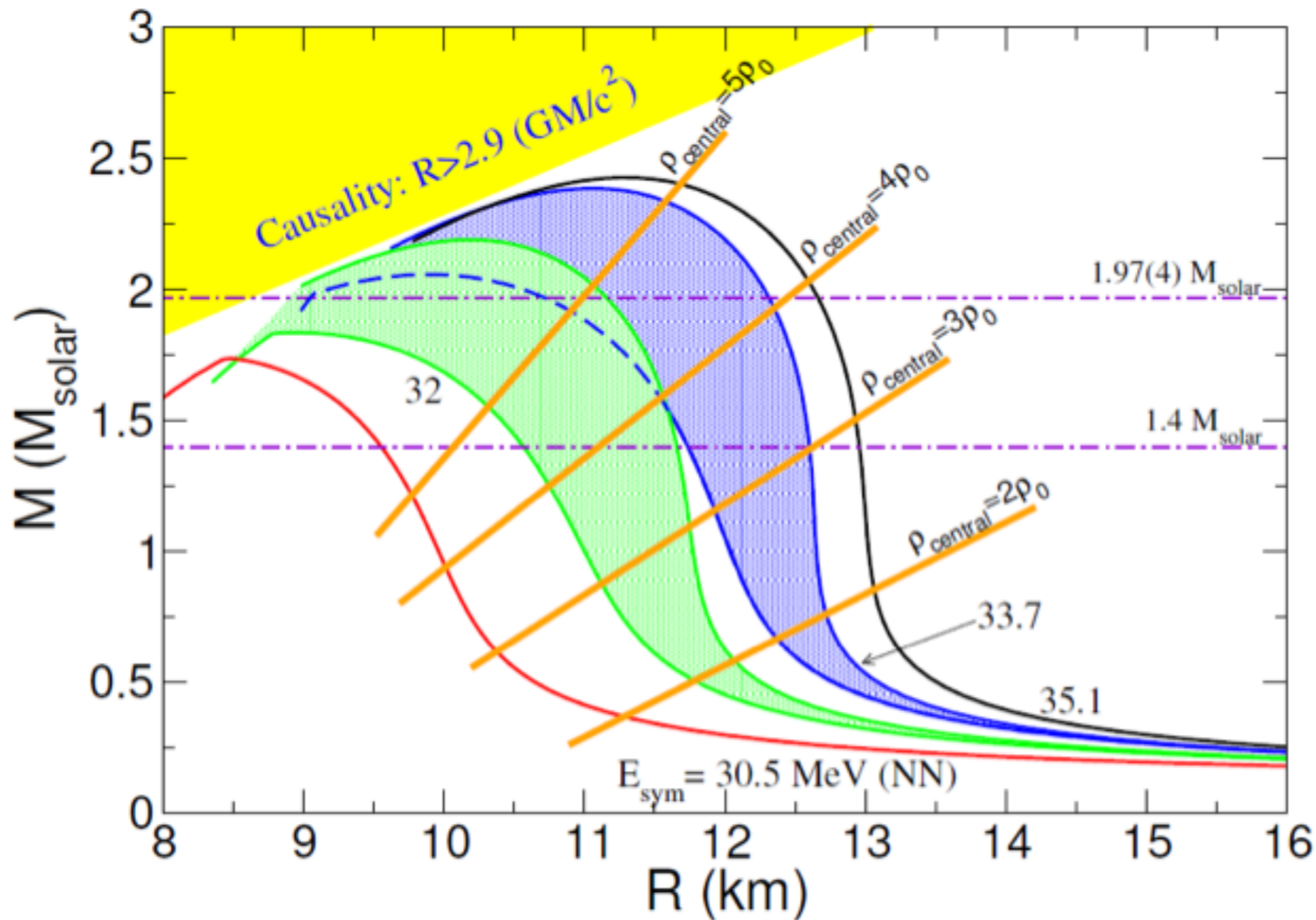
# EOS results



Steiner, Lattimer, and Brown, in prep.

- $P(\varepsilon)$  determined to within 30-50%
- $P(n_B)$  determined to within a factor of 3
- Neutron skin thickness of lead  $\delta R < 0.20$  fm

# Connection to nuclear three-body forces

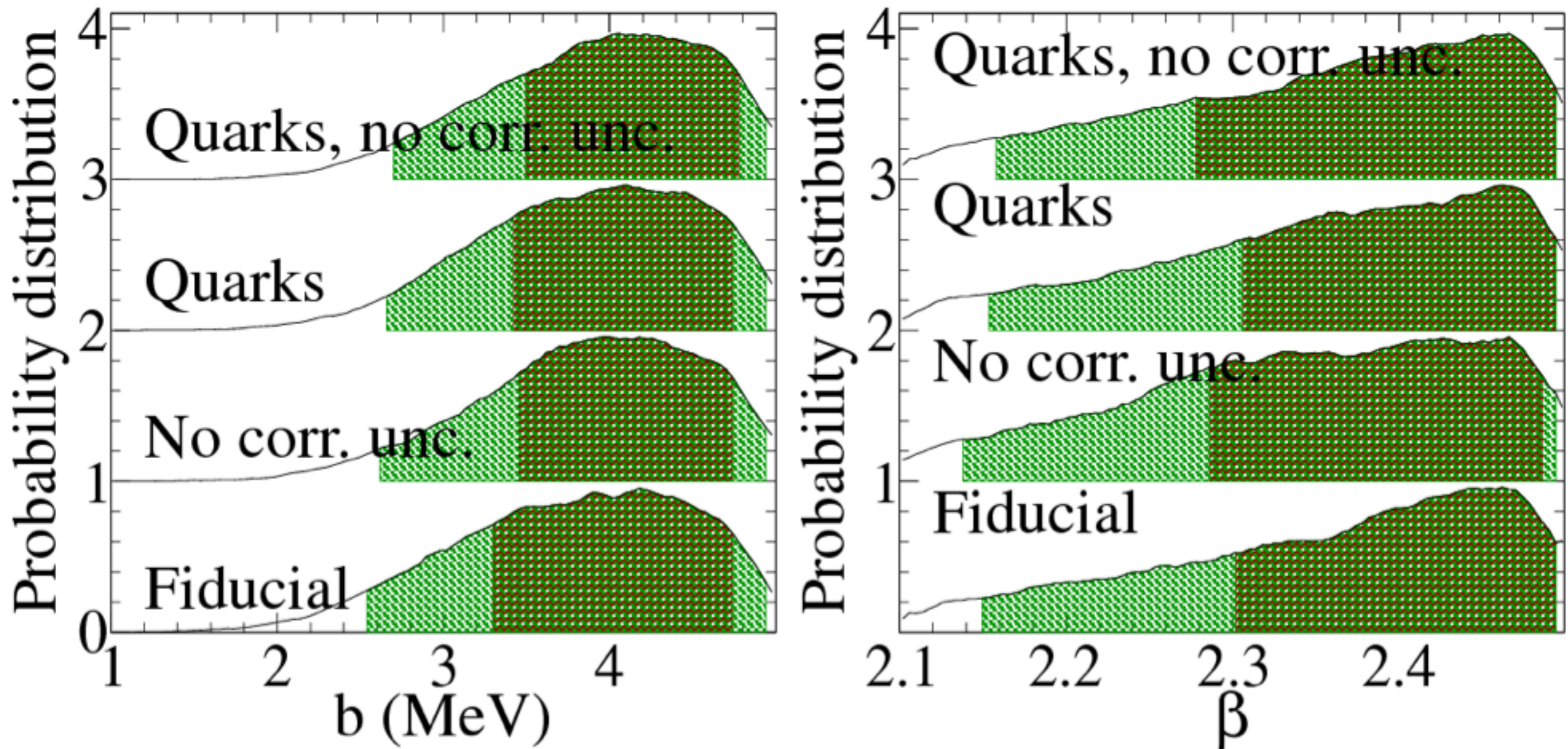


Gandolfi et al. (2012)

- Red = no three-body
- Blue, Green = adjusted three-body interactions
- Black = Urbana IX
- Strong correlation between S and L

- Build up a many-body system from effective two- and three-body nucleon-nucleon interactions
- Three-body forces are also related to neutron star radii
- $E = a \left( \frac{n}{n_0} \right)^\alpha + b \left( \frac{n}{n_0} \right)^\beta$  is a convenient parameterization

# Constraints on three-body force parameters

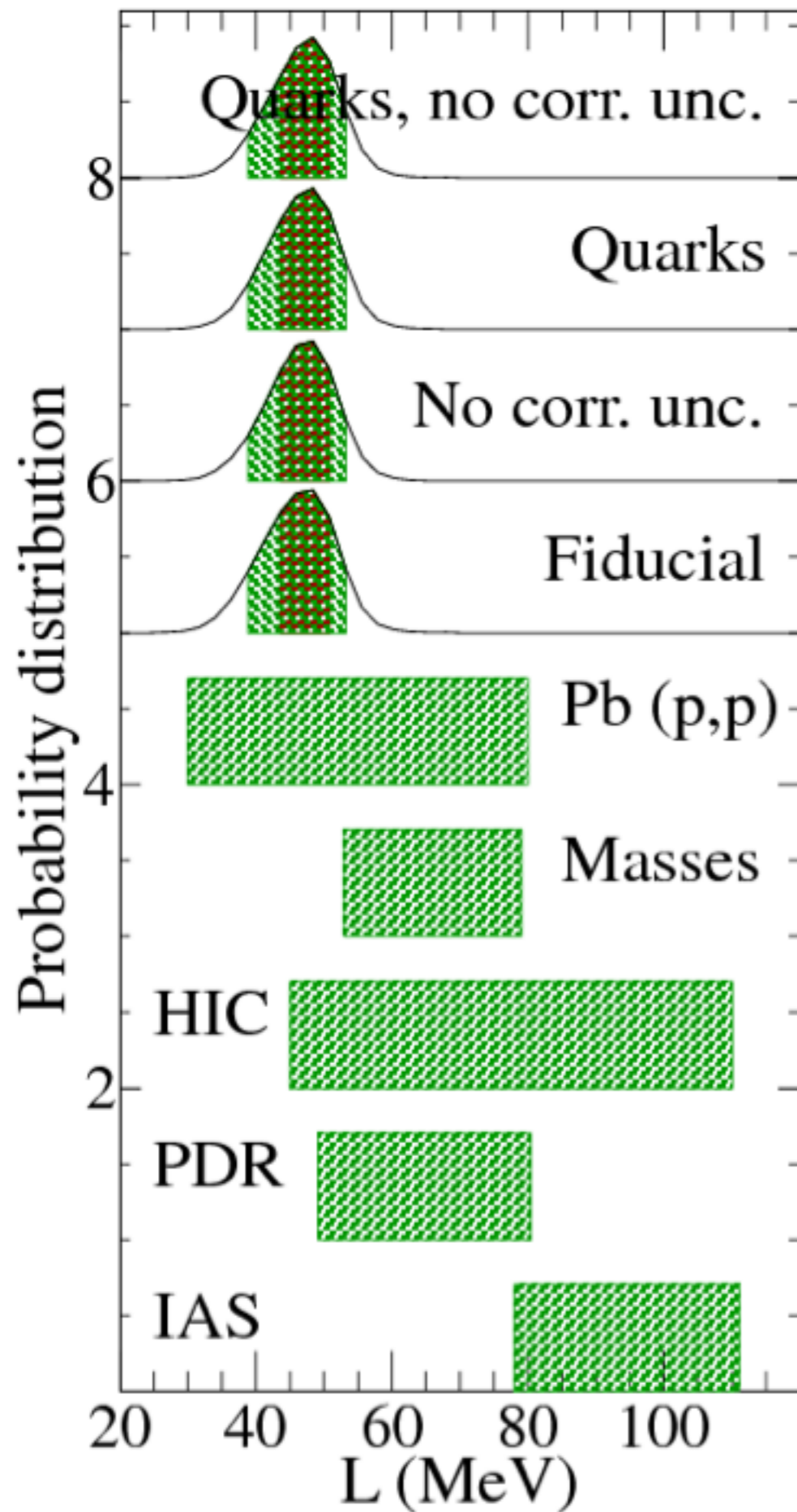


Steiner and Gandolfi (2012)

- Values of  $a$  and  $\alpha$  are unconstrained, but constraints on  $b$  and  $\beta$
- This means that neutron star radii are constraining nuclear three-body forces



## Symmetry Energy Results



- Strong constraints on the derivative of the symmetry energy
- Almost no constraint on  $S$
- Implications for neutron star cooling through the direct Urca process

# The direct Urca process

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- Neutron Star cooling processes

- Modified Urca cooling:  $n + n \rightarrow n + p + e + \bar{\nu}$ ,  
 $Q \sim T^8$

- Cooper pair breaking emissivity  $n + n \rightarrow n + n + \nu + \bar{\nu}$ ,  
 $Q \sim T^7 \log T$

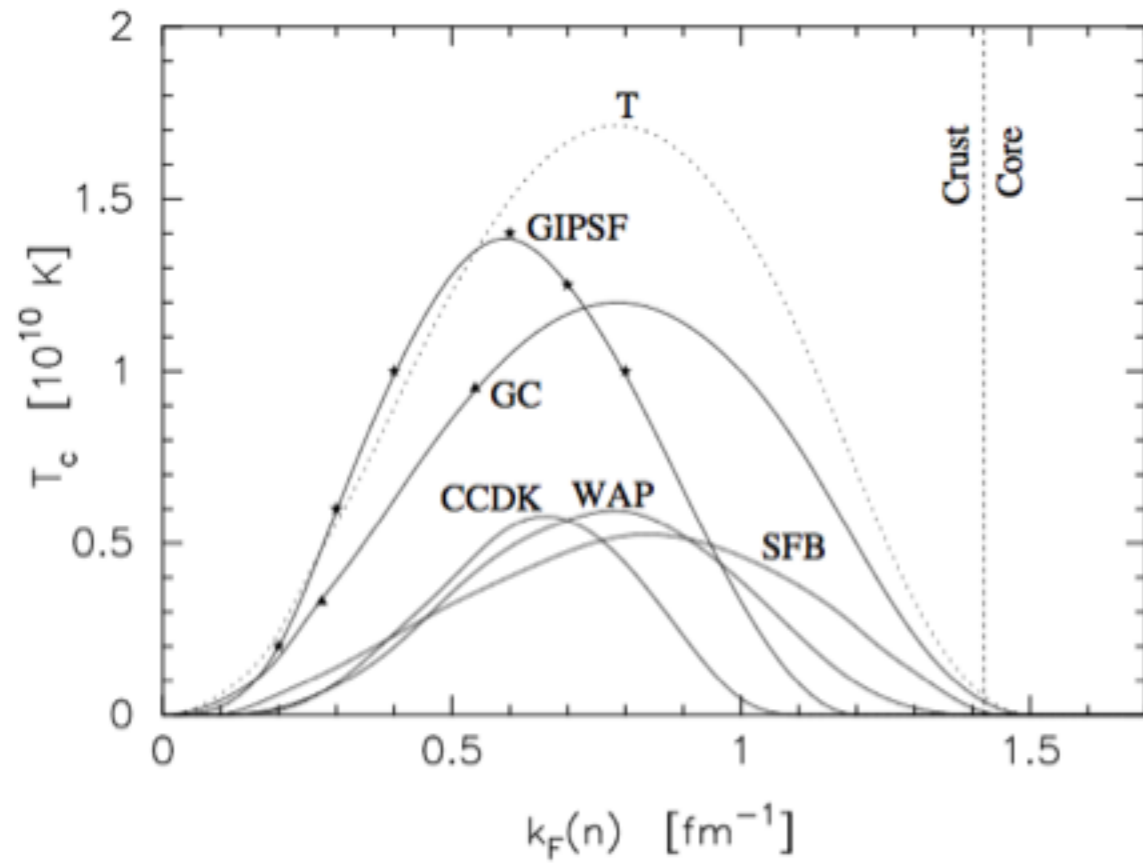
- Direct Urca:  $n \rightarrow p + e + \bar{\nu}$ ,  
 $Q \sim T^6$

- Direct Urca requires a large enough proton Fermi momentum

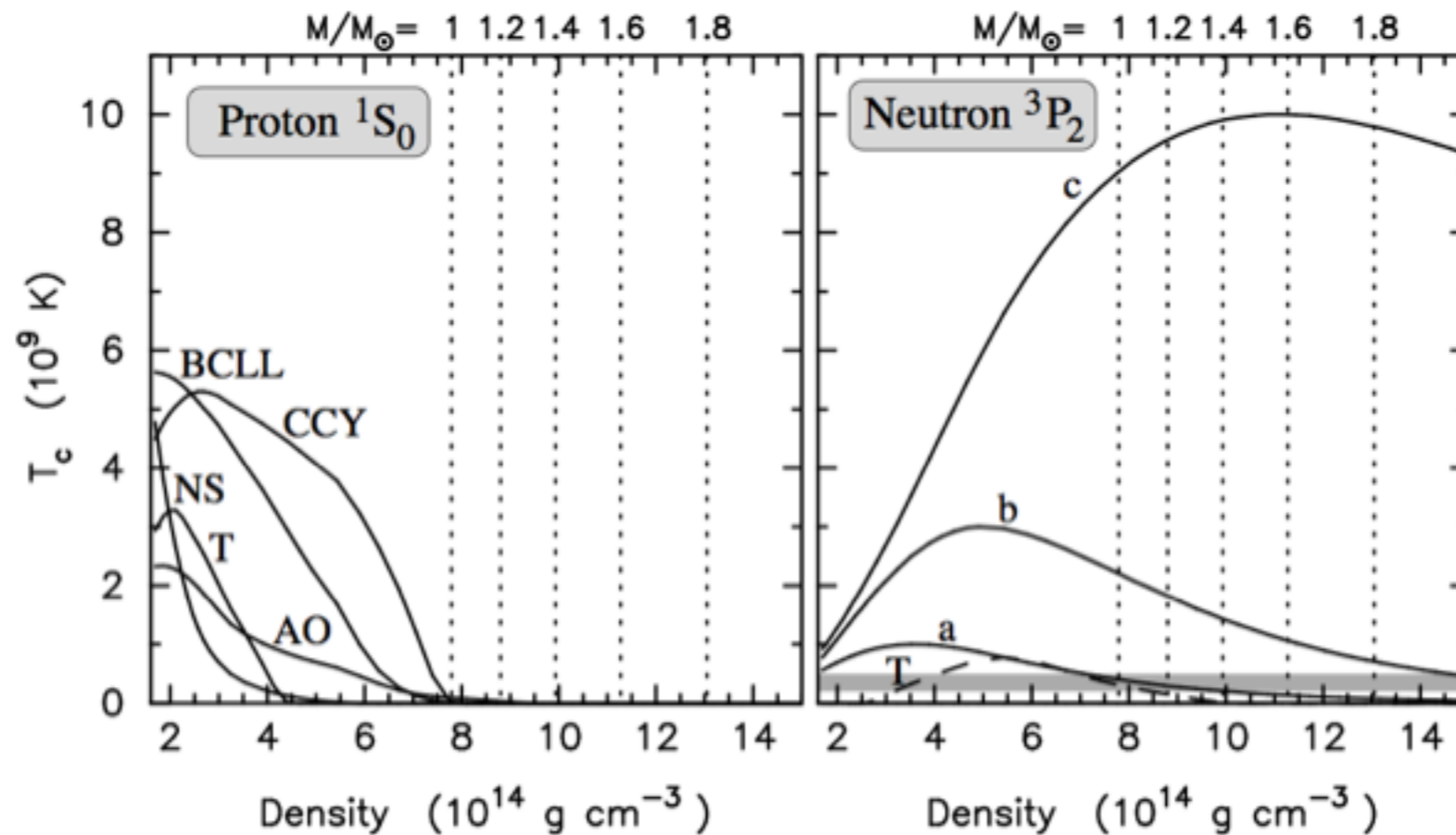
- Thus also connected to the symmetry energy

- Also quark and hyperon direct Urca processes

# Neutron Star Superfluidity



- Superfluidity can block the direct Urca process
- Also blocks quark direct Urca process

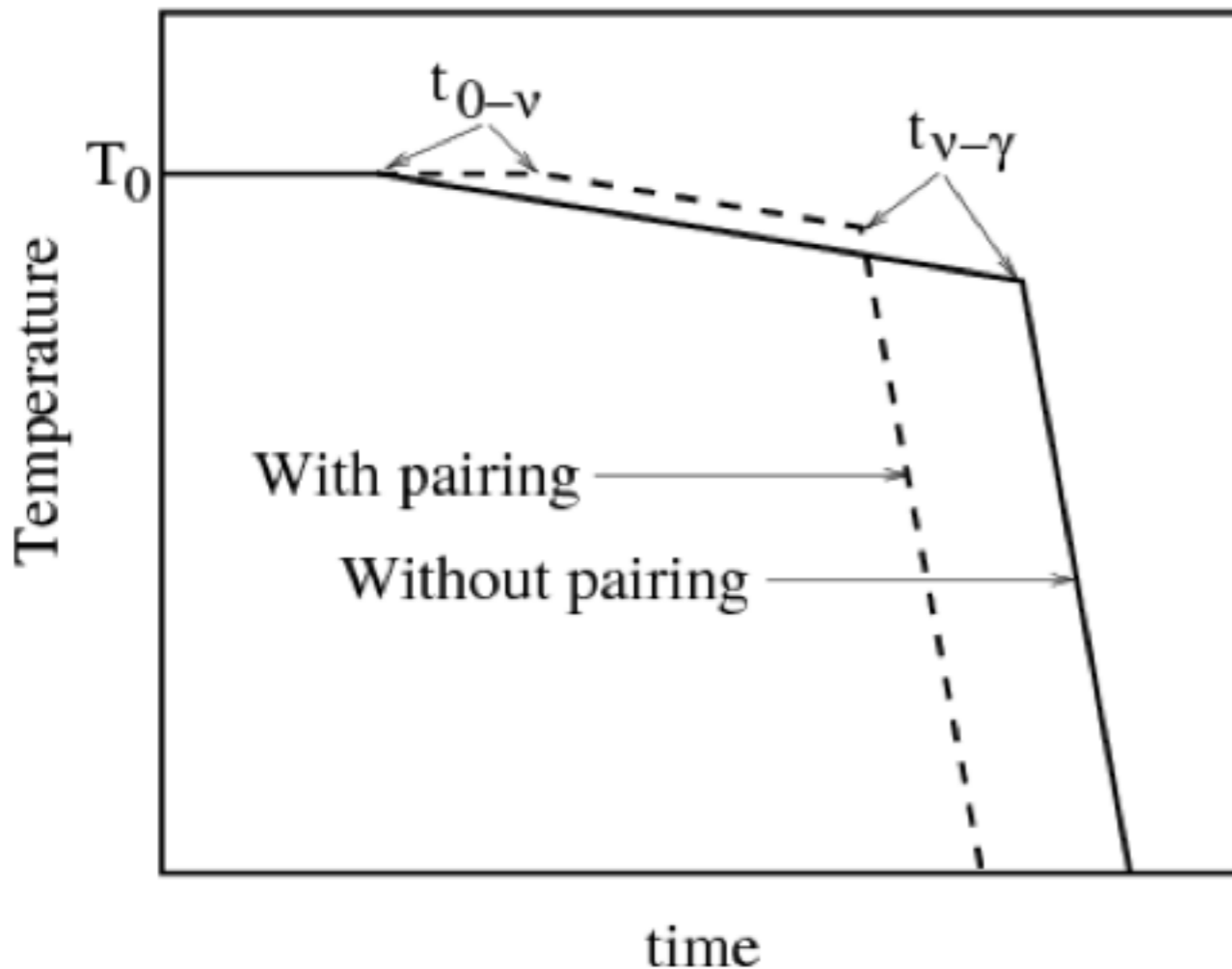


# Thermal Emission from Isolated Neutron Stars

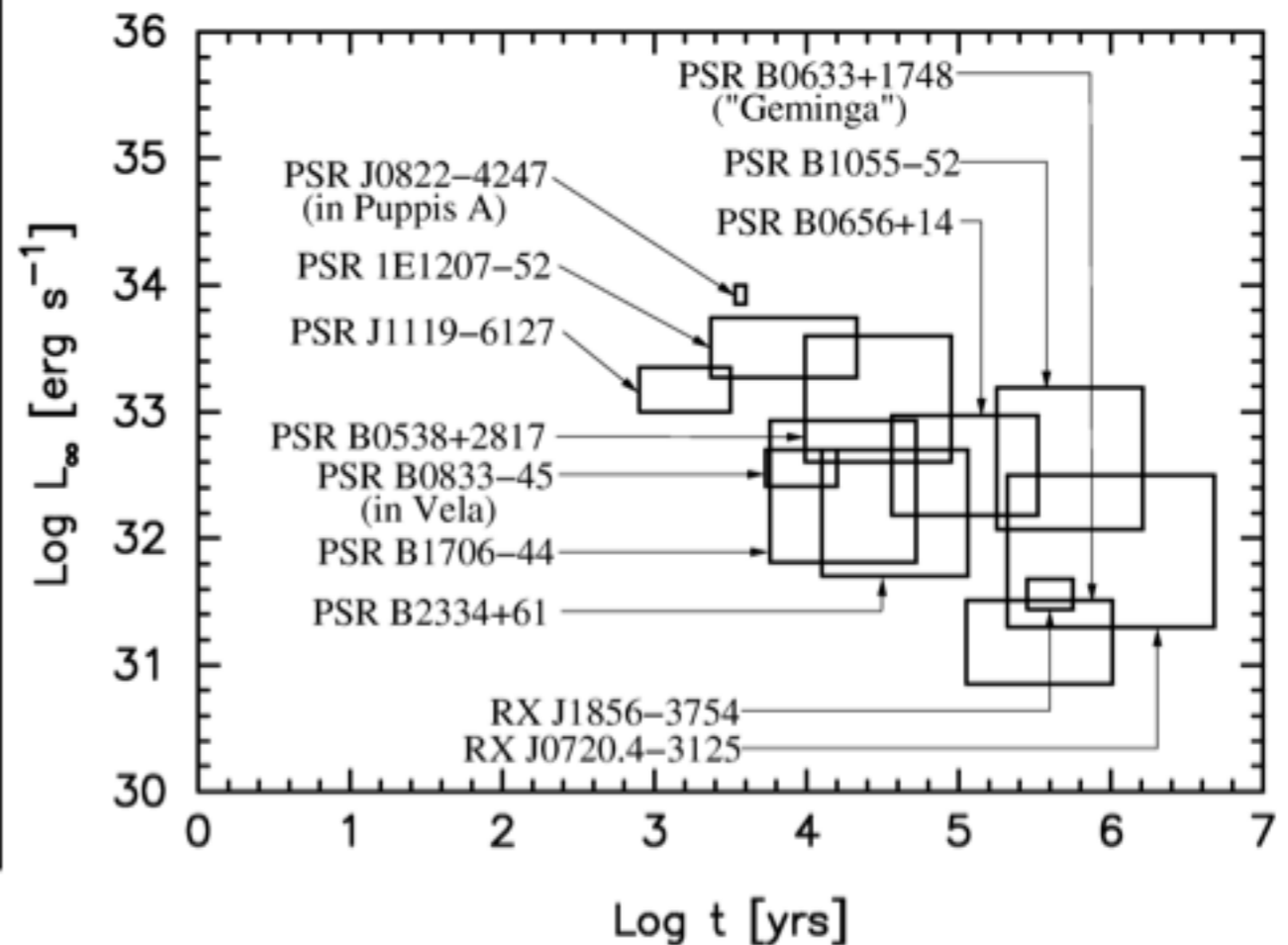
- No distance measurement required
- Requires a model of the NS atmosphere to associate the observed spectrum with a luminosity or temperature

$$C_V \frac{dT}{dt} = L_\nu + L_\gamma, \quad L_\gamma \sim T^{2+4\alpha}, \quad L_\nu \sim T^8 \text{ (Modified Urca)}, \quad C_V \sim CT$$

- Age assumed from spin-down age or associated with a supernova remnant



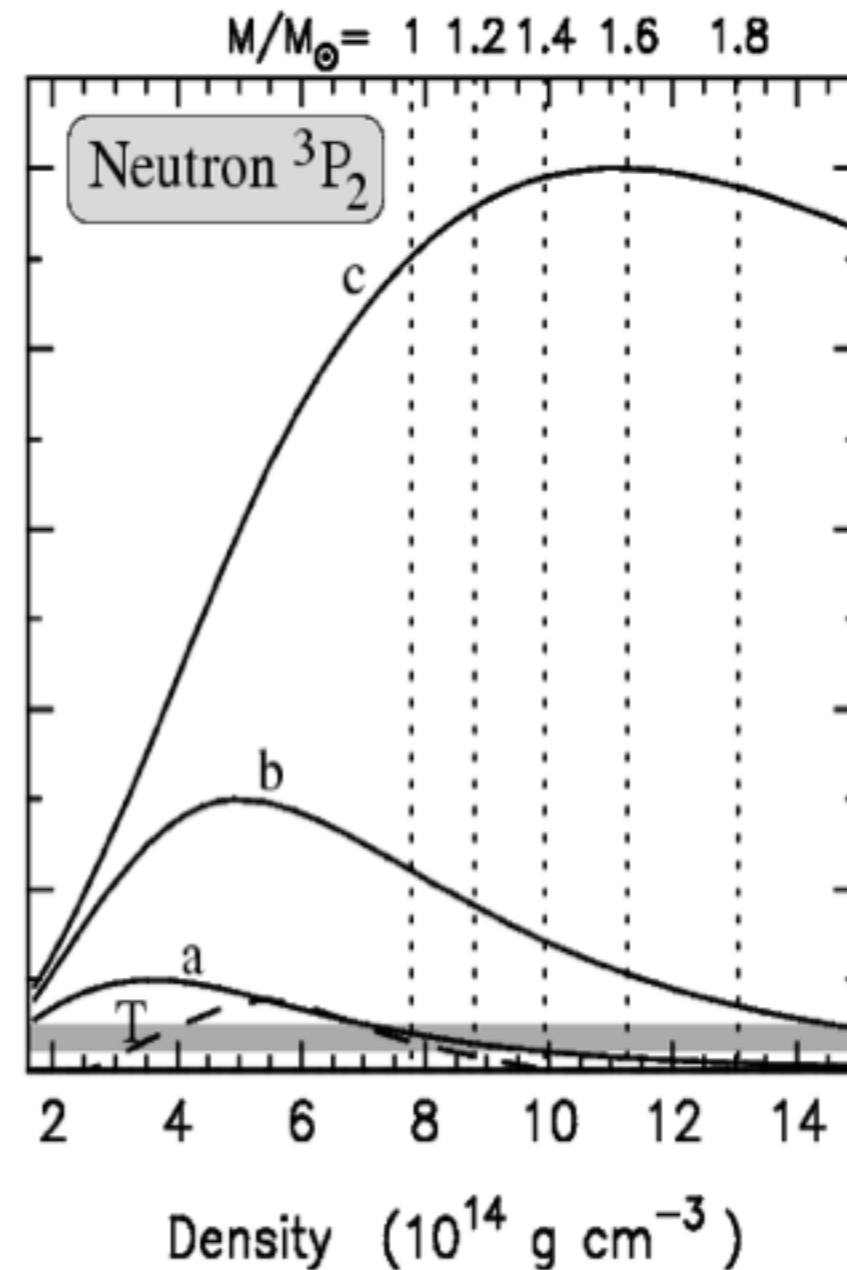
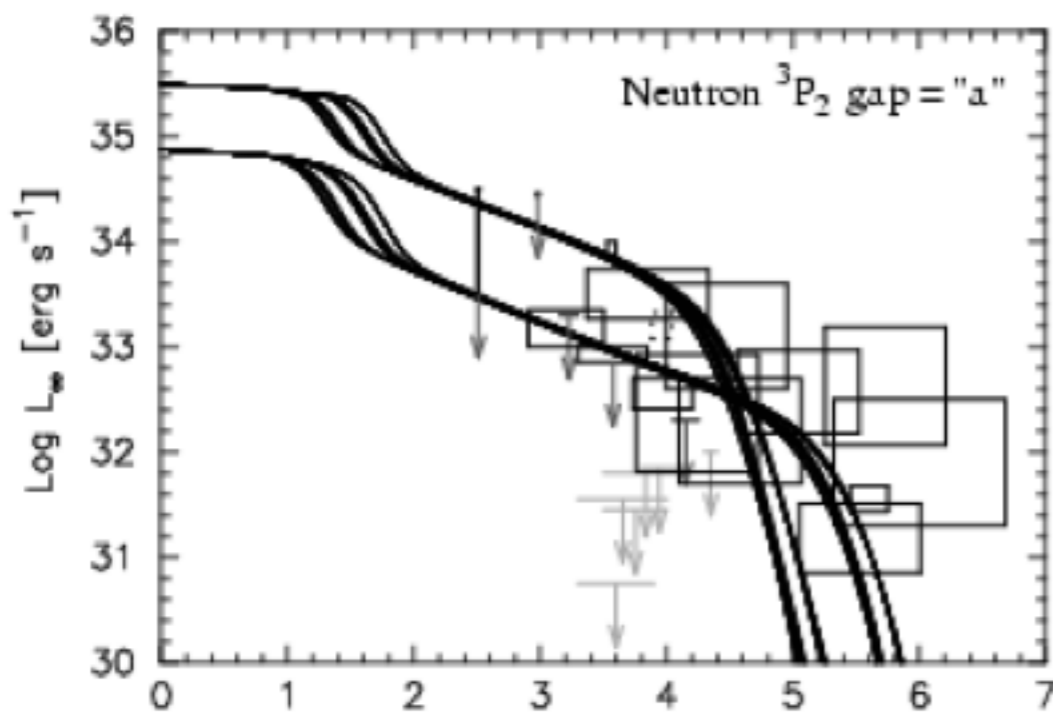
Page, et al (2004)



Page, et al (2009)

# Using Thermal Emission to Constrain Dense Matter

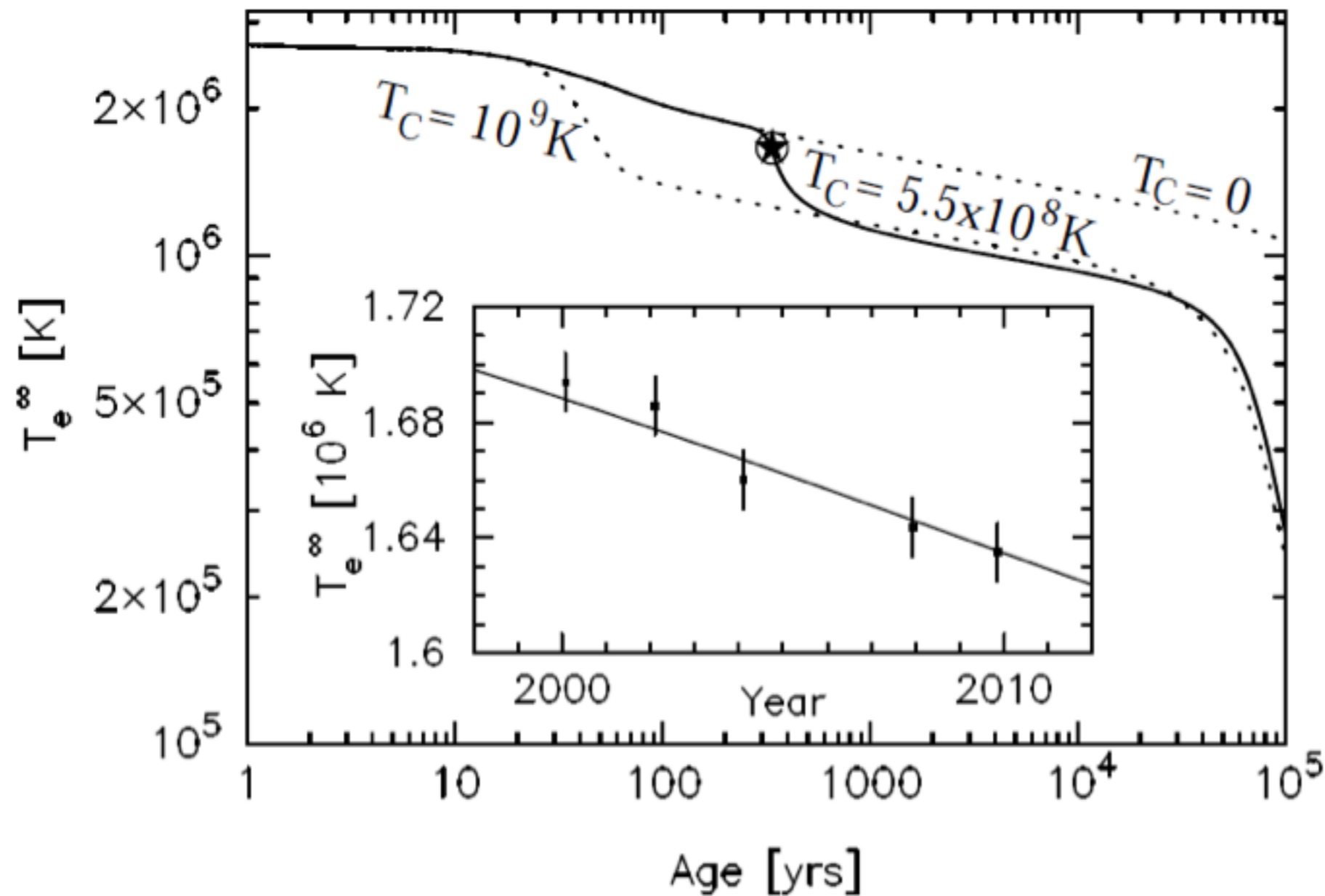
- Minimal model has only neutrons and protons and no direct Urca, but includes all emissivities for neutrons and protons
- Current observations make very stringent constraints on the neutron  ${}^3P_2$  gap



Page et al. (2009), see also Yakovlev et al. and Blaschke et al.

# Neutron Star in Cas A

- The large slope is only well reproduced by the neutron triplet superfluid transition and associated PBF emissivity
- Cas A requires a very particular triplet gap

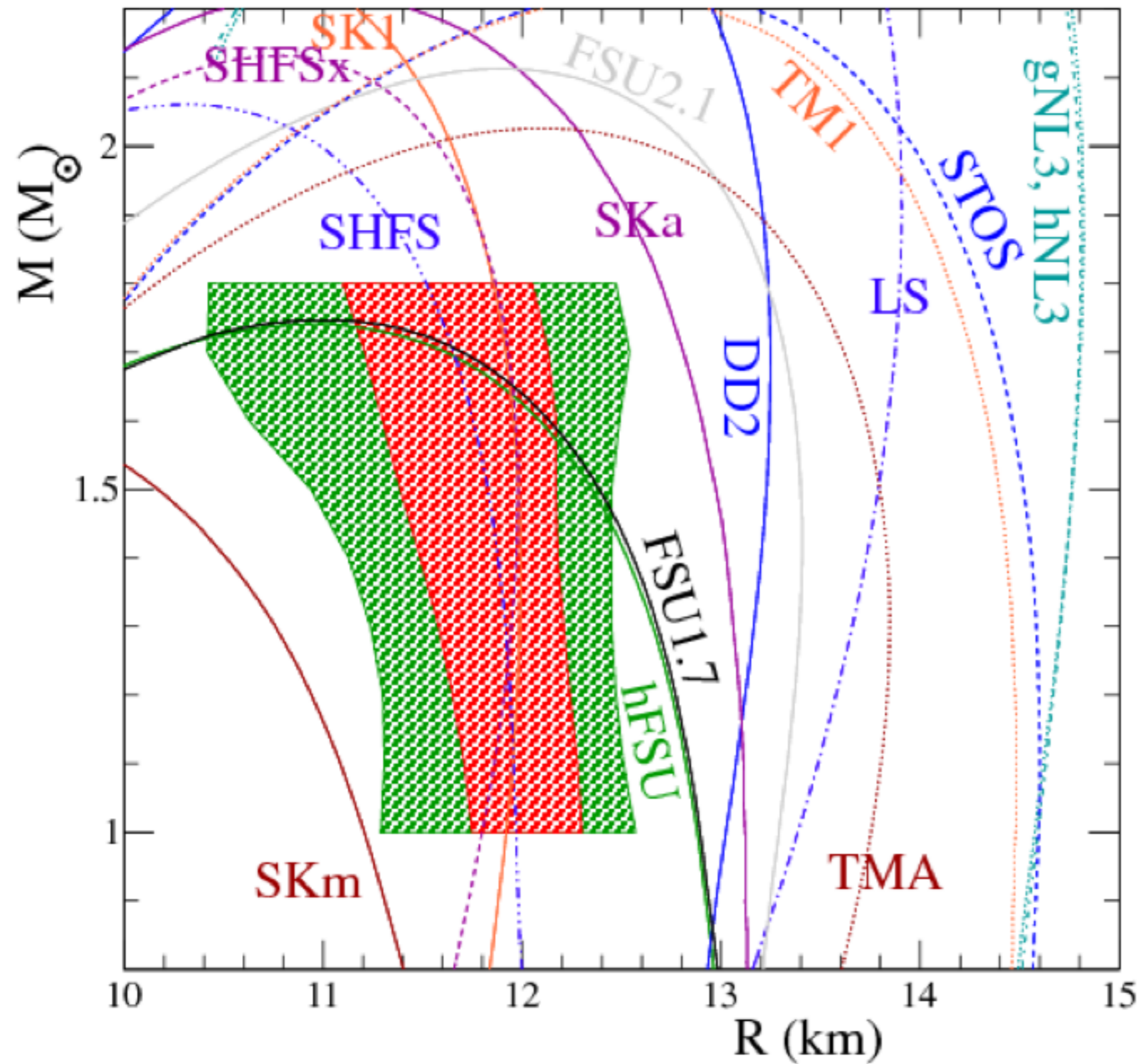


## Dense Matter from Supernovae

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- . Historically: EOSs with large compressibilities generated robust explosions!
- . Time to black hole formation is a signal of phase transition (Gentile et al. 1993, Nakazato et al. 2008)
- . Smaller symmetry energies form more compact NSs, giving larger  $\nu$  luminosities (Sumiyoshi et al. 1995)
- . Protoneutron star metastability (Prakash et al. 1997, Pons et al. 2001):  
Phase transition gives delayed collapse
- . Quark degrees of freedom can impact neutrino signal (Sagert et al. 2009)
- . and more!

# Masses and Radii Implied by Supernova EOSs





## New Supernova EOSs

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- . Two new supernova EOSs
- . RMF mean field model
- . Nuclei determined by Hartree calculations for EOS fit

Quantity	Experiment	Baseline	Extreme
$K$ (MeV)	210-250	245	239
$n_0$ (fm <sup>-3</sup> )	0.155-0.165	0.158	0.160
$-E_B$ (MeV)	15-17	16.2	16.1
$E_{\text{sym}}$ (MeV)	28-36?	31.6	28.8
$E_{\text{Pb-208}}$ (MeV)	7.87	7.81	7.79
$R_{\text{Pb-208}}$ (fm)	5.50	5.44	5.43
$E_{\text{Zr-90}}$ (MeV)	8.71	8.64	8.57
$R_{\text{Zr-90}}$ (fm)	4.27	4.19	4.20

## Status of Matter at High Densities

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- . Quark matter is neither required nor ruled out
- . Neither is absolutely stable quark matter
- . Two solar mass neutron star is a problem for hyperons!  
(see poster by Miyatsu)
- . Hyperons + quarks?
- .  $n - n - \Lambda$  interactions?



# Summary

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- *Current* mass and radius measurements, modulo some systematic uncertainties, can *quantitatively* constrain the equation of state of dense matter
- Several currently used EOSs are ruled out
- Current results imply all neutron stars have radii between 10.4 and 12.9 km
- Isolated neutron star cooling is giving us information about dense matter
- Supernova EOSs need updating!

# The Issue of the Photosphere Radius

- Dimensionless parameter

$$\alpha \equiv \frac{F_{TD} \kappa D}{\sqrt{A} c^3 f_c^2}$$

