

# New equations of state in simulations of core-collapse supernovae

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Matthias Hempel, Basel University  
Waseda, 8.3.2012

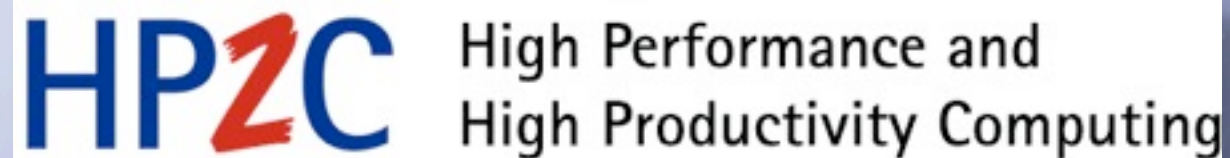
in collaboration with:

Tobias Fischer (GSI Darmstadt)

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Matthias Liebendörfer (Basel)

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# New equations of state in simulations of core-collapse supernovae

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## Outline:

- 1.) introduction
- 2.) EOS model & constraints
- 3.) EOS aspects in simulations
  - a.) sub-saturation densities
  - b.) supra-saturation densities
- 4.) conclusions

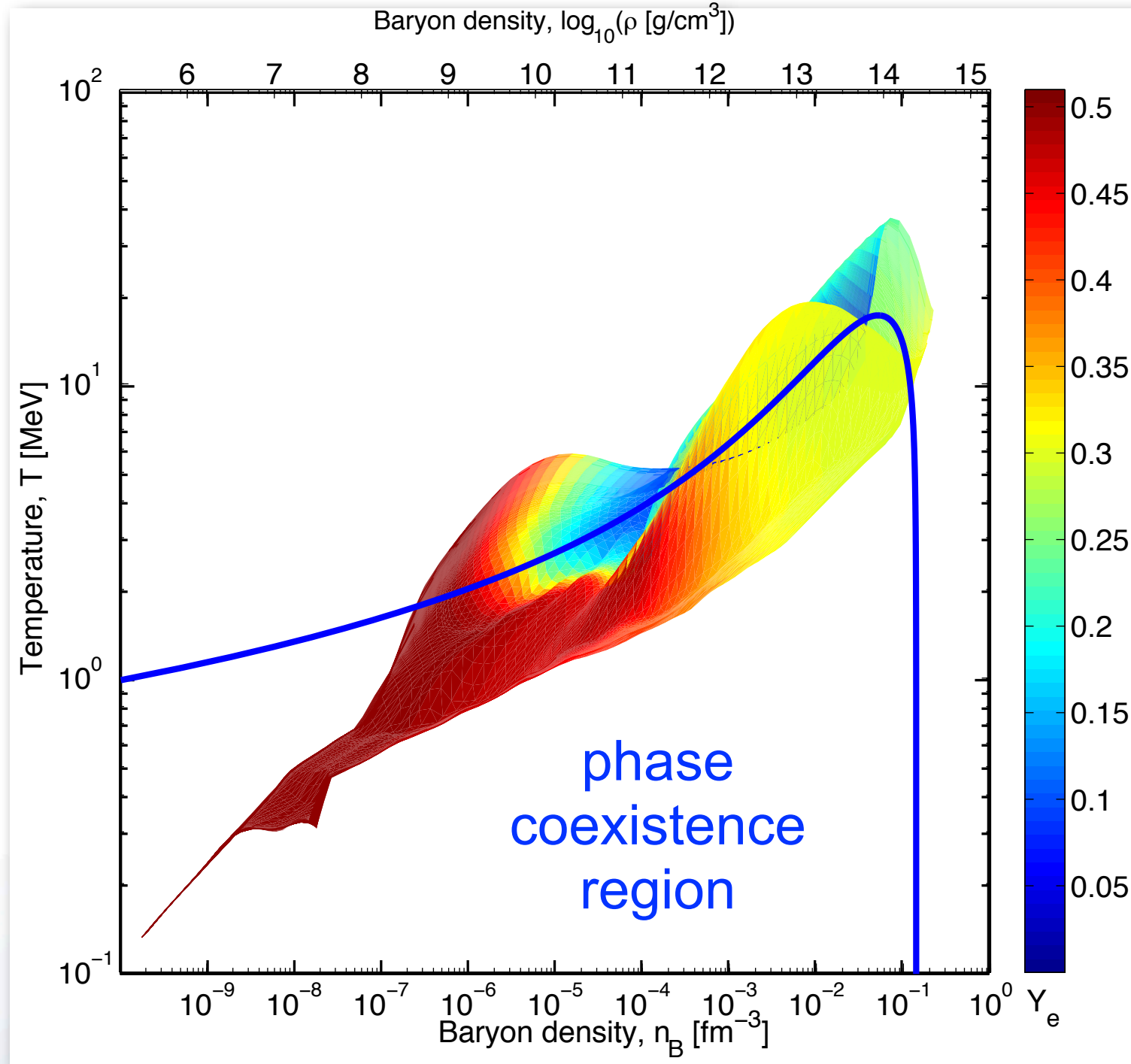
# Introduction

- EOS provides the crucial nuclear physics input for astrophysical simulations: thermodynamic quantities and nuclear composition
- plenty of EOSs for cold neutron stars
- challenge of the supernova (SN) EOS:
  - finite temperature:  $T = 0 - 100 \text{ MeV}$
  - no weak equilibrium:  $Y_p = 0 - 0.6$
  - wide density range:  $\rho = 10^4 - 10^{15} \text{ g/cm}^3$
- SN EOS: multi-purpose EOS, also applicable to neutron stars and their mergers, proto-neutron stars, heavy-ion collisions, ...



# Introduction

## state of matter in core-collapse supernovae



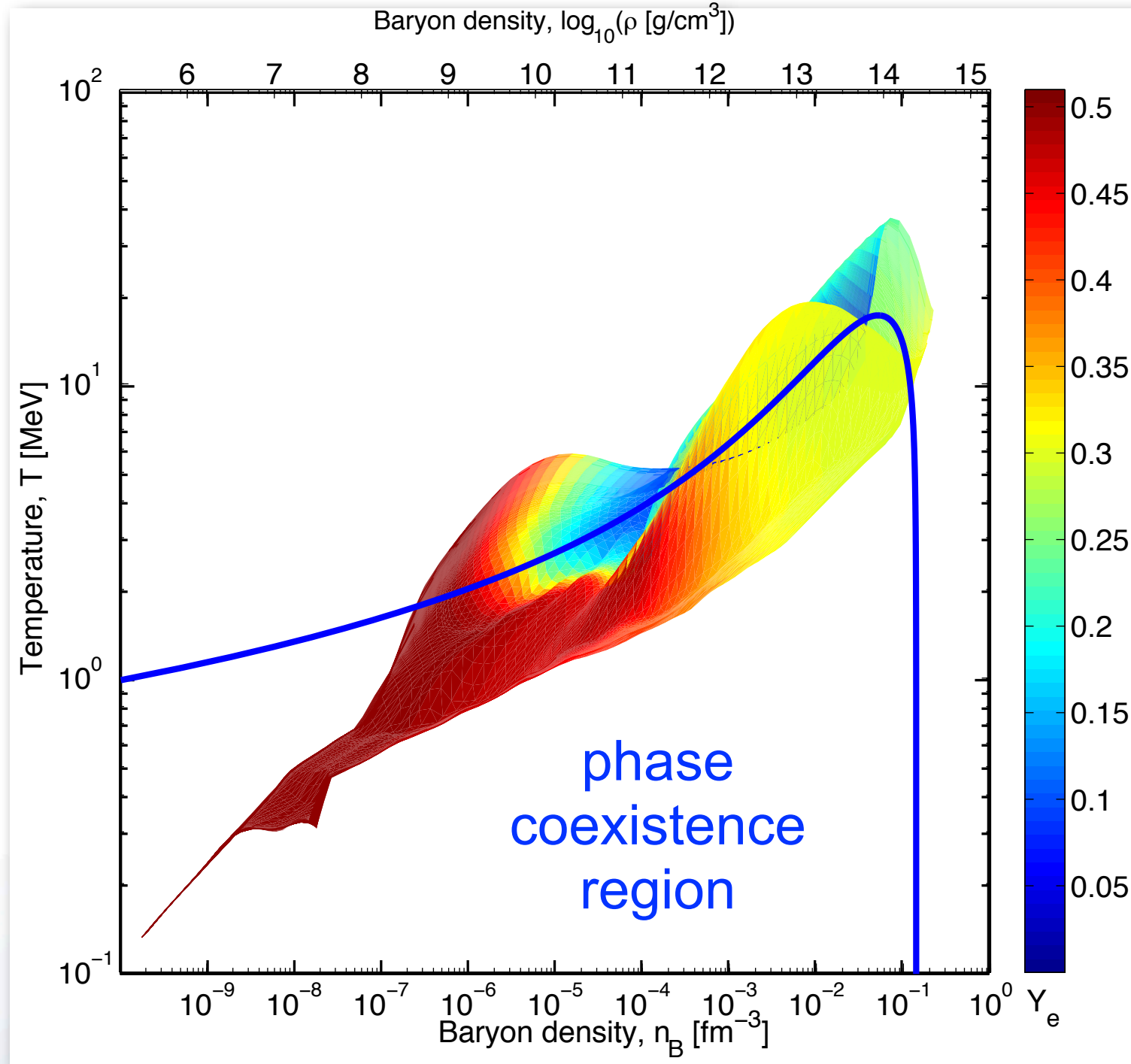
- without Coulomb, „bulk“:  
first order liquid-gas phase transition below  $T_c \sim 15$  MeV
- phase coexistence region
- with finite size effects:  
→ non-uniform nuclear matter, mixture of nuclei and nucleons

based on: [Fischer et al., ApJS 2010]



# Introduction

## state of matter in core-collapse supernovae



- 99% of the energy emitted in neutrinos
- detailed knowledge of thermodynamic quantities and nuclear composition crucial
- at  $\rho \sim 10^{11} - 10^{13}$  g/cm $^3$ : neutrino spheres
- stall of the shock front

based on: [Fischer et al., ApJS 2010]

# Introduction

commonly used SN EOSs:

- Lattimer & Swesty 1991 (LS): non-relativistic liquid drop model
- H. Shen, Toki, Omayatsu and Sumiyoshi 1998 (STOS/Shen): relativistic mean-field (RMF), Thomas-Fermi approximation
- both models:
  - one representative nucleus: “single nucleus approximation”
  - no shell effects
  - only  $\alpha$ -particles of light clusters

# Recent developments

new available SN tables:

Nakazato et al. 2008	quark matter with large $n_c$ added to STOS
Ishizuka et al. 2008	Hyperons added to STOS
Sagert et al. 2009	quark matter with low $n_c$ added to STOS → explosions in 1D
H. Shen et al. 2010	Lambdas added to STOS

strange physics  
at large densities

G. Shen, Horowitz, Teige (2010)	tables for NL3 and FSUgold: virial expansion, relativistic mean-field (RMF), Hartree
MH & Schaffner-Bielich (2010)	tables for NL3, TM1, TMA, FSUgold, DD2: NSE, RMF, excluded volume
Typel et al. (2010)	table for DD2 in preparation: gRMF
Furusawa et al. (2011)	table for TM1 in preparation: liquid drop, NSE, RMF

new hadronic EOS



# Recent developments

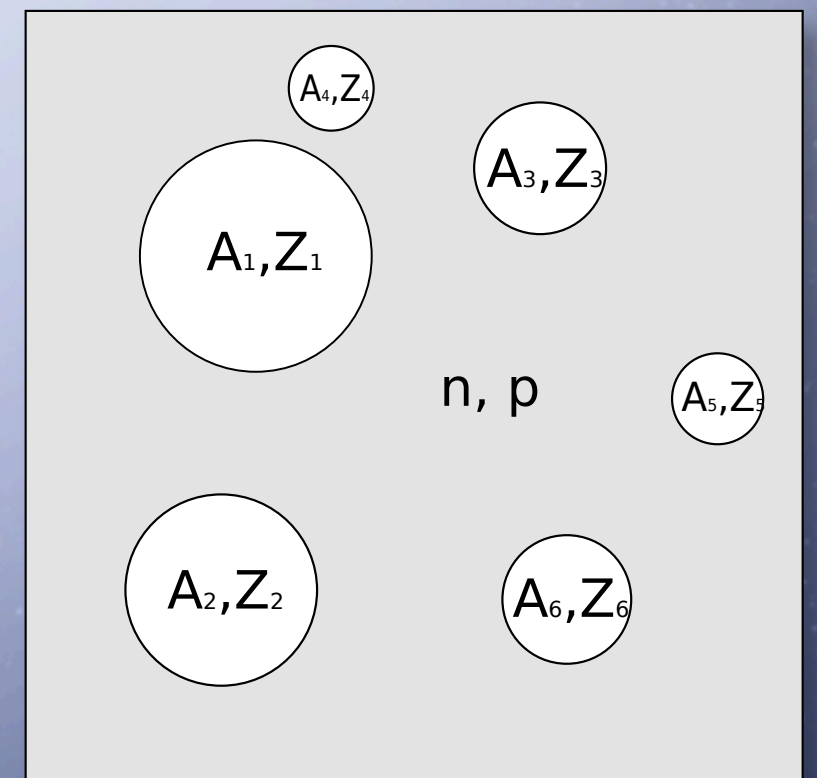
- pasta phases: Horowitz et al. 2004, Newton & Stone 2009
- light nuclei in the (SN) medium:  
Horowitz & Schwenk 2006, Röpke 2009, Typel et al. 2010
- weak reactions:  
Langanke 2003, Arcones et al. 2008
- multifragmentation/NSE models in the supernova context:  
Ishizuka et al. 2003, Botvina & Mishustin 2004, 2010, Blinnikov et al. 2009
- ...

# EOS model & constraints

# Excluded volume NSE model with interactions

MH, J. Schaffner-Bieleich; NPA837(2010) (HS)

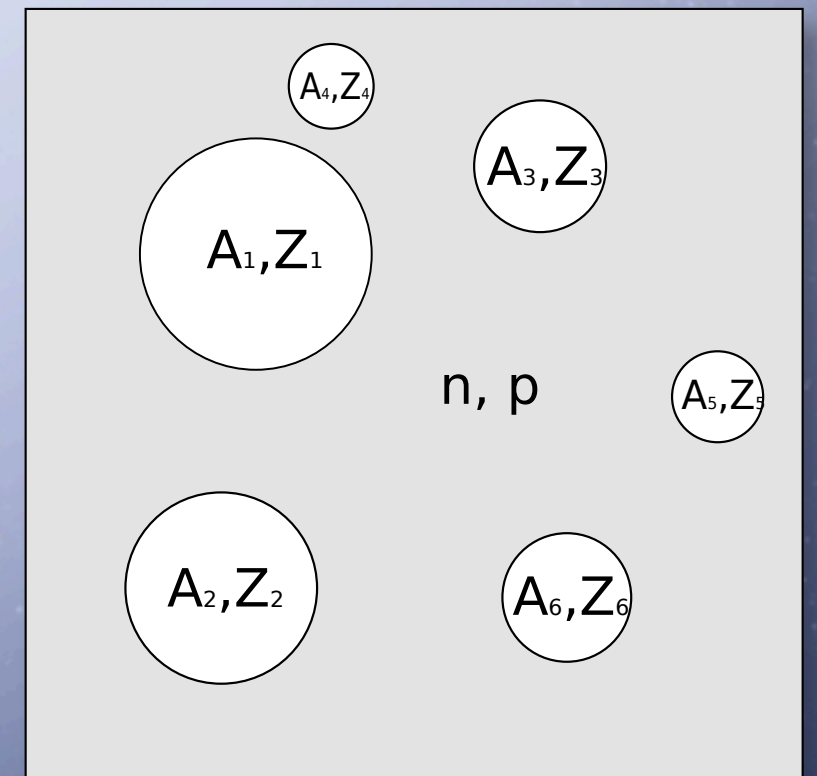
- chemical mixture of nuclei and interacting nucleons in nuclear statistical equilibrium
- model characteristics:
  - ensemble of heavy nuclei
  - inclusion of all possible light nuclei
  - nuclear shell effects
  - various nucleon interactions
- input: RMF interactions, nuclear mass tables, Coulomb energies, excited states, and excluded volume effects





# Excluded volume NSE model with interactions

- five tables for different RMF interactions:  
NL3, TM1, TMA, FSUgold, DD2
- $T = (0.1 - 200) \text{ MeV}$
- $n_B = (10^{-12} - 10) \text{ fm}^{-3}$
- $Y_p = 0.01 - 0.6$
- 1.6 million grid-points
- full nuclear composition



# Nucleons – non-linear RMF (TM1, TMA, FSUGold, NL3)

- relativistic mean-field model (RMF)
- interactions mediated via exchange of mesons and meson (self-) interactions

$$\begin{aligned}\mathcal{L} = & \bar{\psi}(i\gamma^\mu\partial_\mu - M)\psi \\ & + \frac{1}{2}\partial_\mu\sigma\partial^\mu\sigma - \frac{1}{2}m_\sigma^2\sigma^2 - \frac{1}{3}g_2\sigma^3 - \frac{1}{4}g_3\sigma^4 - g_\sigma\bar{\psi}\sigma\psi \\ & - \frac{1}{4}\Omega_{\mu\nu}\Omega^{\mu\nu} + \frac{1}{2}m_\omega^2\omega_\mu\omega^\mu + \frac{1}{4}g_4(\omega_\mu\omega^\mu)^2 - g_\omega\bar{\psi}\gamma^\mu\psi\omega_\mu \\ & - \frac{1}{4}R^a{}_{\mu\nu}R^{a\mu\nu} + \frac{1}{2}m_\rho^2\rho_\mu^a\rho^{a\mu} - g_\rho\bar{\psi}\gamma_\mu\tau^a\psi\rho^{\mu a} - \Lambda\omega_\mu\omega^\mu\rho_\nu^a\rho^{a\nu}\end{aligned}$$

- alternative: density-dependent coupling constants (DD2)
- coupling constants fitted to experimental data
- well-established description of finite nuclei and nuclear matter

*Sugahara & Toki 1994*

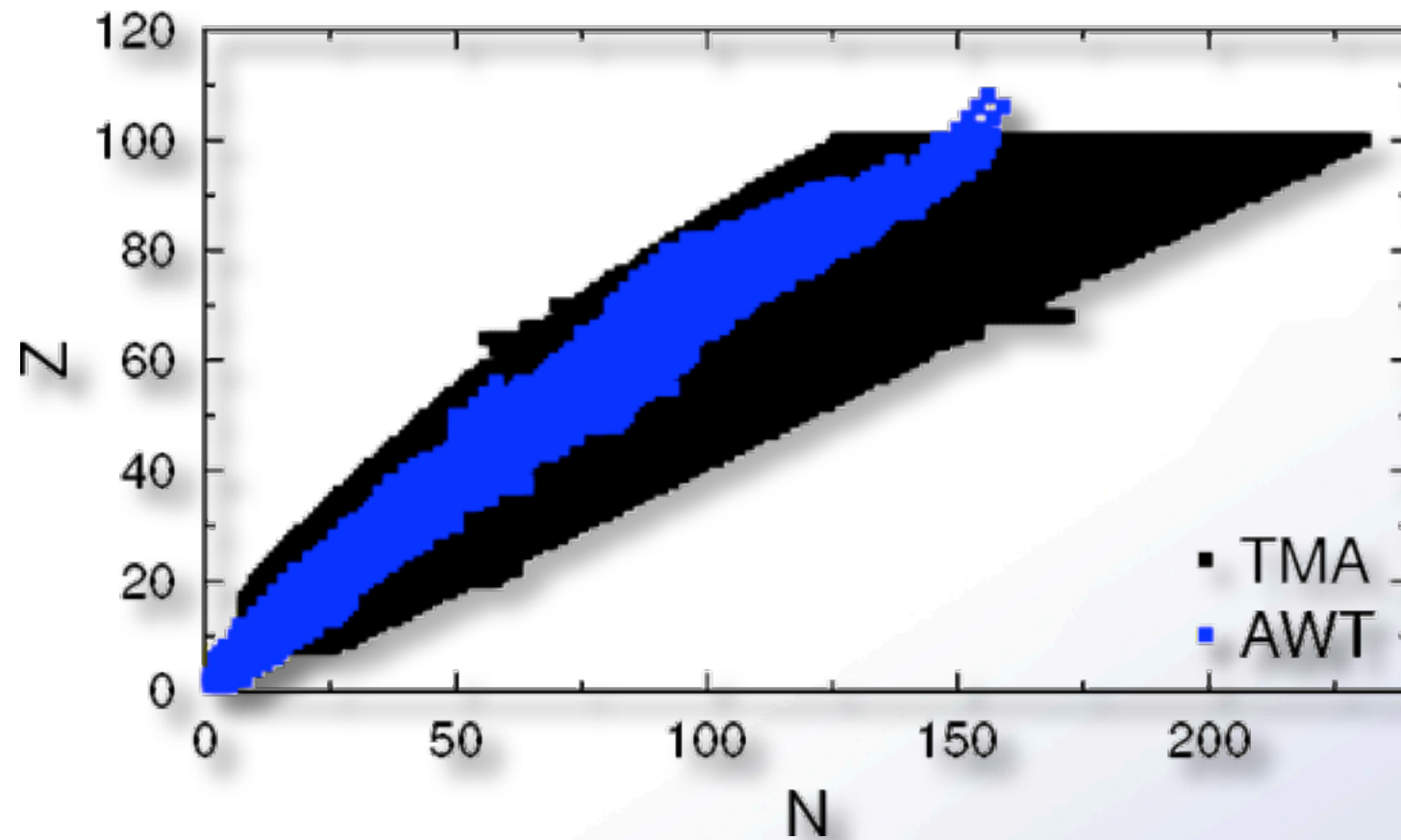
*Toki et al. 1995*

*Todd-Rutel & Piekarewicz 2005*

*Lalazissis et al. 1997*

*Typel 2005, Typel et al. 2010*

# Nuclei - mass tables



- experimental data of AWT (2003): over 2000 precisely measured binding energies

*[Audi et al.; 2003NPA729]*

unknown nuclei: mass table generated with theoretical nuclear model

- TMA: BCS  $\delta$ -force pairing, axial deformations, mass-number dependent coupling constants *[Geng et al.; 2005PTP113]*
- FSUgold: spherical, BCS pairing with constant matrix elements *[Roca-Maza, private communication] [Roca-Maza, Piekarewicz; 2008PRC78]*
- NL3: BCS  $\delta$ -force pairing, axial deformations *[Lalazissis et al.; 1999ADNDT71]*
- FRDM: finite range droplet model *[Möller et al.; 1995ADNDT59]*



# Nuclei - medium effects

- semi-empirical formula for intrinsic partition function of hot nuclei

$$g_i(T) = g_i^0 + \frac{c_1}{A_i^{5/3}} \int_0^{B_i} dE e^{-E/T} \exp(2\sqrt{a_i E}) \quad [Fai, Randrup; 1982NPA381]$$

- Coulomb energies in T=0 Wigner-Seitz approximation

$$E_{Coul} = -\frac{3}{5} \frac{Z_i^2 \alpha}{R_i} \left( \frac{3}{2} x - \frac{1}{2} x^3 \right)$$

- possible extension: combination with liquid-drop model

→ see talk of S. Furusawa

# Excluded volume

model assumptions

- each baryon fills a volume  $1/n_B^0$ , volume of a nucleus  $V_{A,Z} = A/n_B^0$

- nucleons must not be inside of nuclei

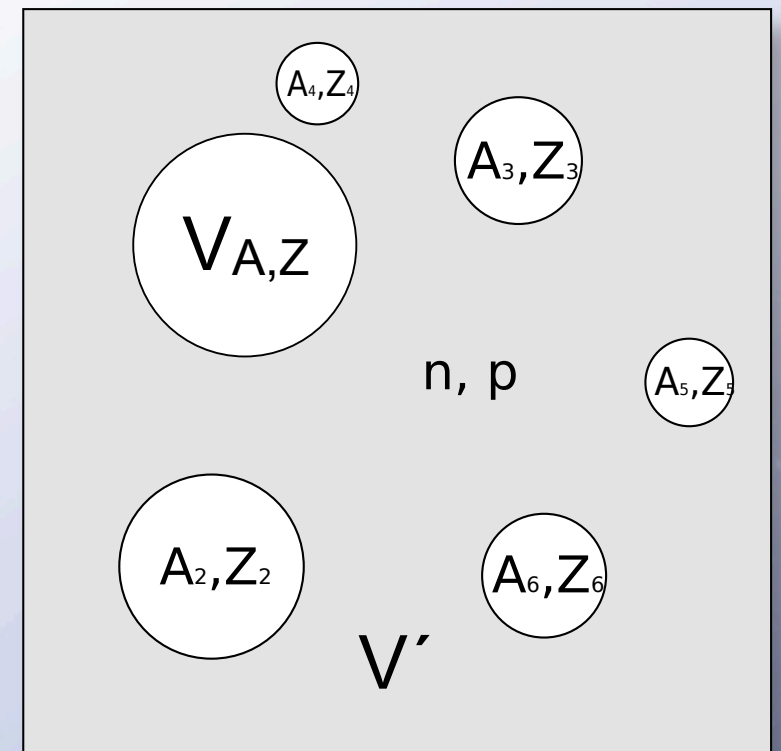
- nuclei must not overlap

• filling factor of the nucleons

$$\begin{aligned}\xi &= V'/V \\ &= 1 - \sum_{A,Z} A n_{A,Z}/n_B^0\end{aligned}$$

• effect on nuclei given by the free volume fraction

$$\kappa = 1 - n_B/n_B^0$$



# Excluded volume - asymptotics

- thermodynamic potential: free energy density

$$f = f_e^0(T, n_e) + \sum_{A,Z} f_{A,Z}^0(T, n_{A,Z}) + f_{Coul}(n_e, \{n_{A,Z}\}) \\ + \xi f_{RMF}^0(T, n'_n, n'_p) - T \sum_{A,Z} n_{A,Z} \ln(\kappa) ,$$

ideal gas expressions, Coulomb contribution, filling factor for nucleons, direct excluded volume contribution

- low fraction of nuclei,  $\xi \sim 1$ : unmodified RMF model ✓
- $\kappa \rightarrow 0$  for  $n_B \rightarrow n_B^0 \Rightarrow f \rightarrow \infty$  if nuclei are present  
→ dissociation of nuclei above saturation density ✓
- thermodynamic fully consistent



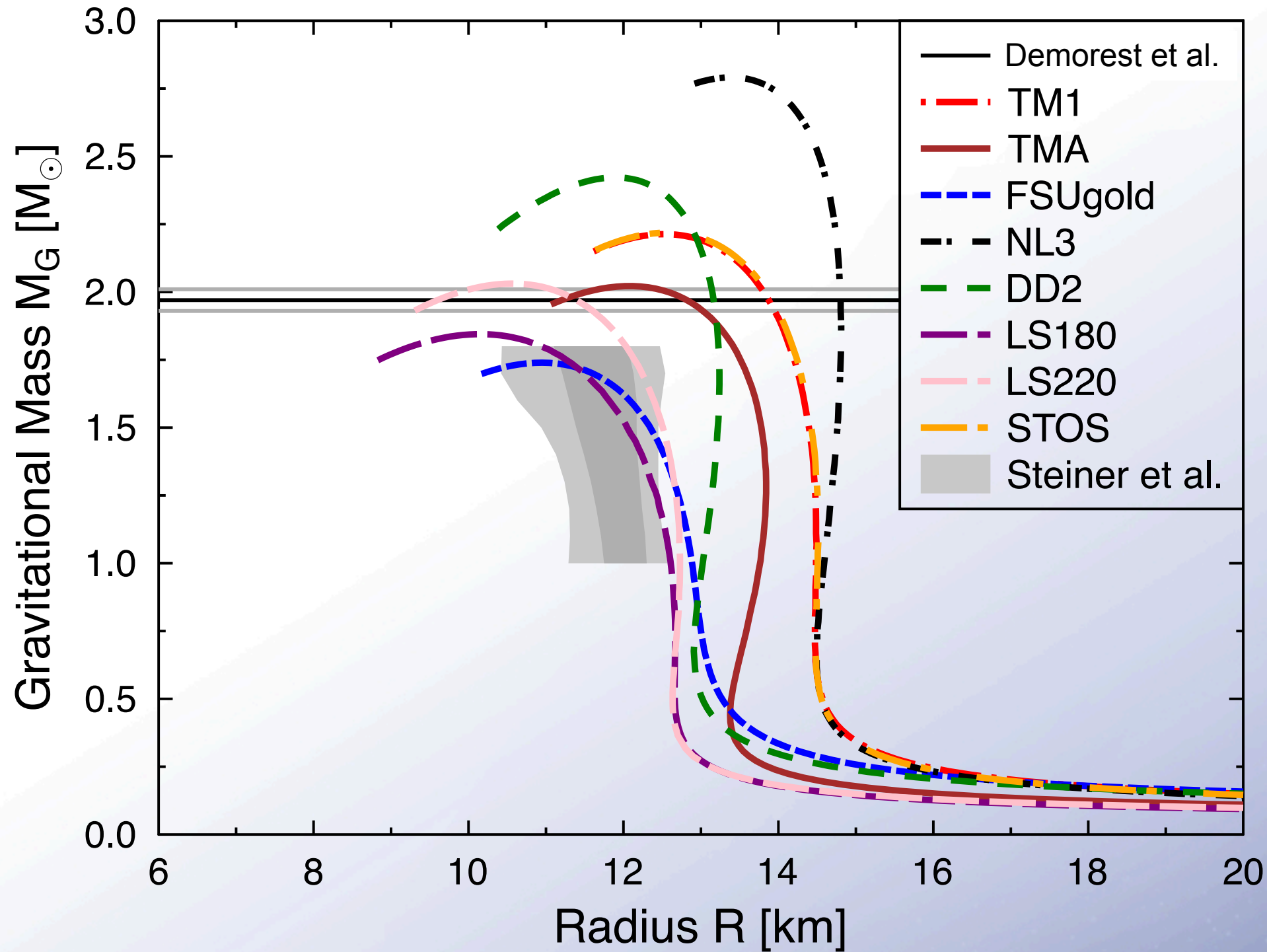
# EOS constraints - saturation properties & maximum mass

	$n_B^0$ [fm $^{-3}$ ]	$E_0$ [MeV]	$K$ [MeV]	$J$ [MeV]	$L$ [MeV]	$M_{\max}$ [ $M_{\odot}$ ]
TM1	0.146	-16.31	282	36.95	110.99	2.213
TMA	0.147	-16.03	318	30.66	90.14	2.022
FSUgold	0.148	-16.27	230	32.56	60.44	1.739
NL3	0.148	-16.24	271	37.39	118.50	2.791
DD2	0.149	-16.02	243	31.67	55.04	2.422
LS180	0.155	-16.00	180	28.61	73.82	1.828
LS220	0.155	-16.00	220	28.61	73.82	2.031
Exp.	$\sim 0.15$	$\sim -16$	$240 \pm 10$ [1]	$30 - 34$ [2]	$40 - 110$ [2]	$> 1.97 \pm 0.04$ [3]

- span a broad range of possible RMF models

	references	type of constraint
[1]	Piekarewicz JPG 2010	compilation of measurements of isoscalar giant monopole resonances
[3]	Tsang et al. PRL 2009 Carbone et al. PRC 2010	compilation of experiments: isospin diffusion, pygmy dipole resonance, nuclear masses, GDR, isoscaling, antiprotonic atoms
[4]	Demorest et al. Nature 2010	measurement of Shapiro delay

# EOS constraints - mass and radius measurements



- bayesian analysis of observations of six NS (X-ray burster, low -mass X-ray binaries), Steiner et al. ApJ 2010

→ see talk of A. Steiner for further discussion

# Results - low density

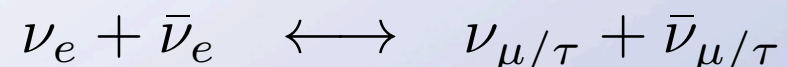
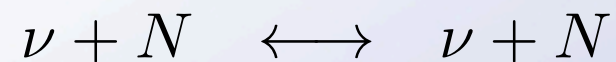
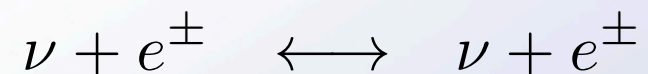
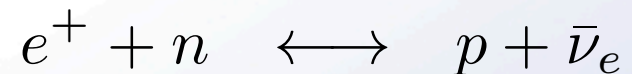
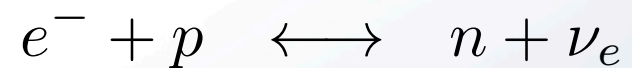
- including simulations for 15  $M_{\text{sun}}$  progenitor of Woosley & Weaver ApJS 101 (1995)
- HS TM1 in comparison with STOS TM1 and/or LS 180

MH, T. Fischer, J. Schaffner-Bielich, M. Liebendörfer, ApJ 748, 70 (2012)

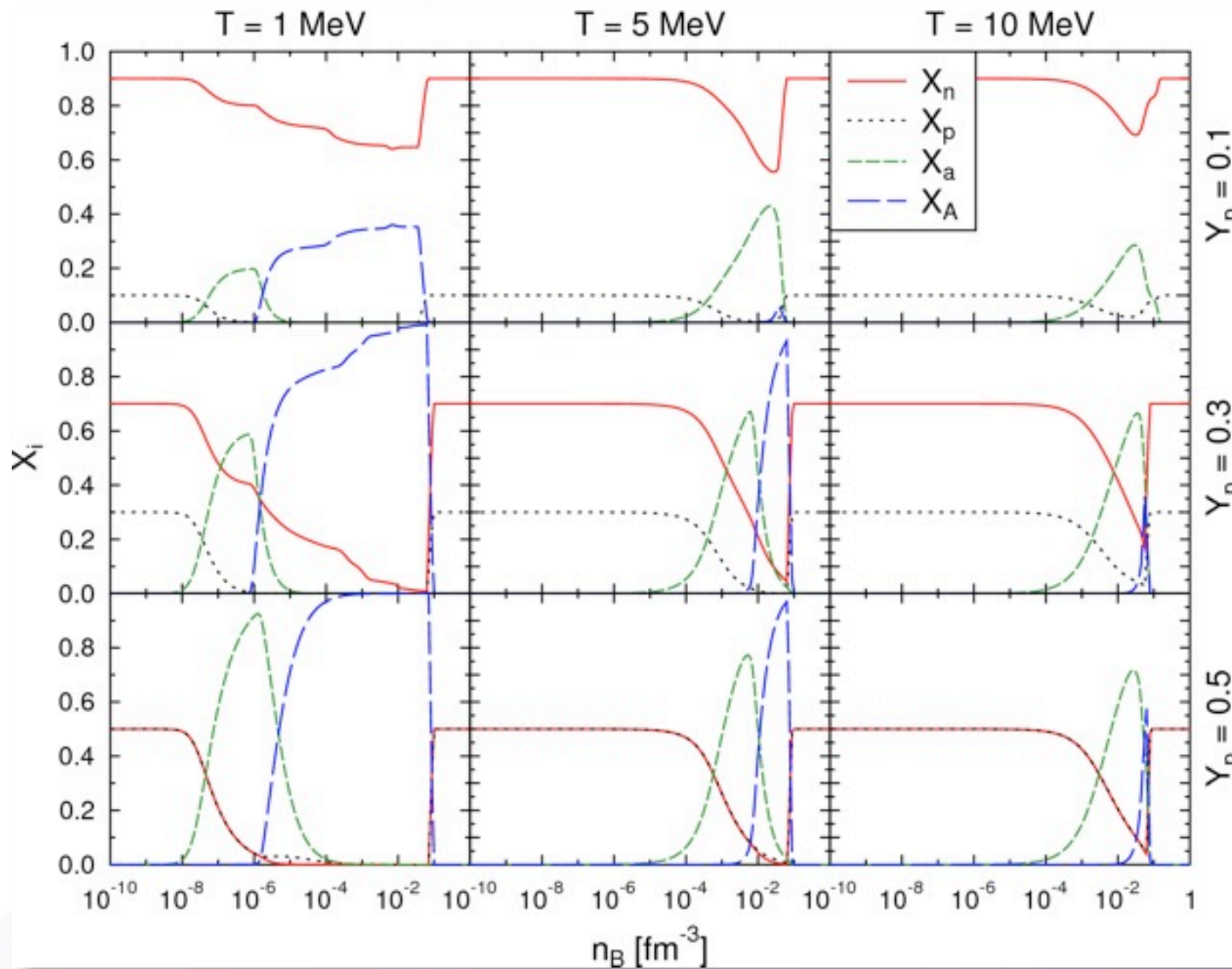


# Supernova simulations

- core-collapse SN simulations by Tobias Fischer, GSI Darmstadt
  - spherical symmetry
  - general relativistic radiation hydrodynamics
  - three flavor neutrino transport
- weak reactions:
  - all light clusters treated as alpha-particles
  - only average heavy nucleus



# Nuclear composition



- general isothermal density dependence:  
 nucleons  $\rightarrow$  light  $\rightarrow$  (heavies)  $\rightarrow$  uniform nucleon matter
- low temperature:  
 shell effects

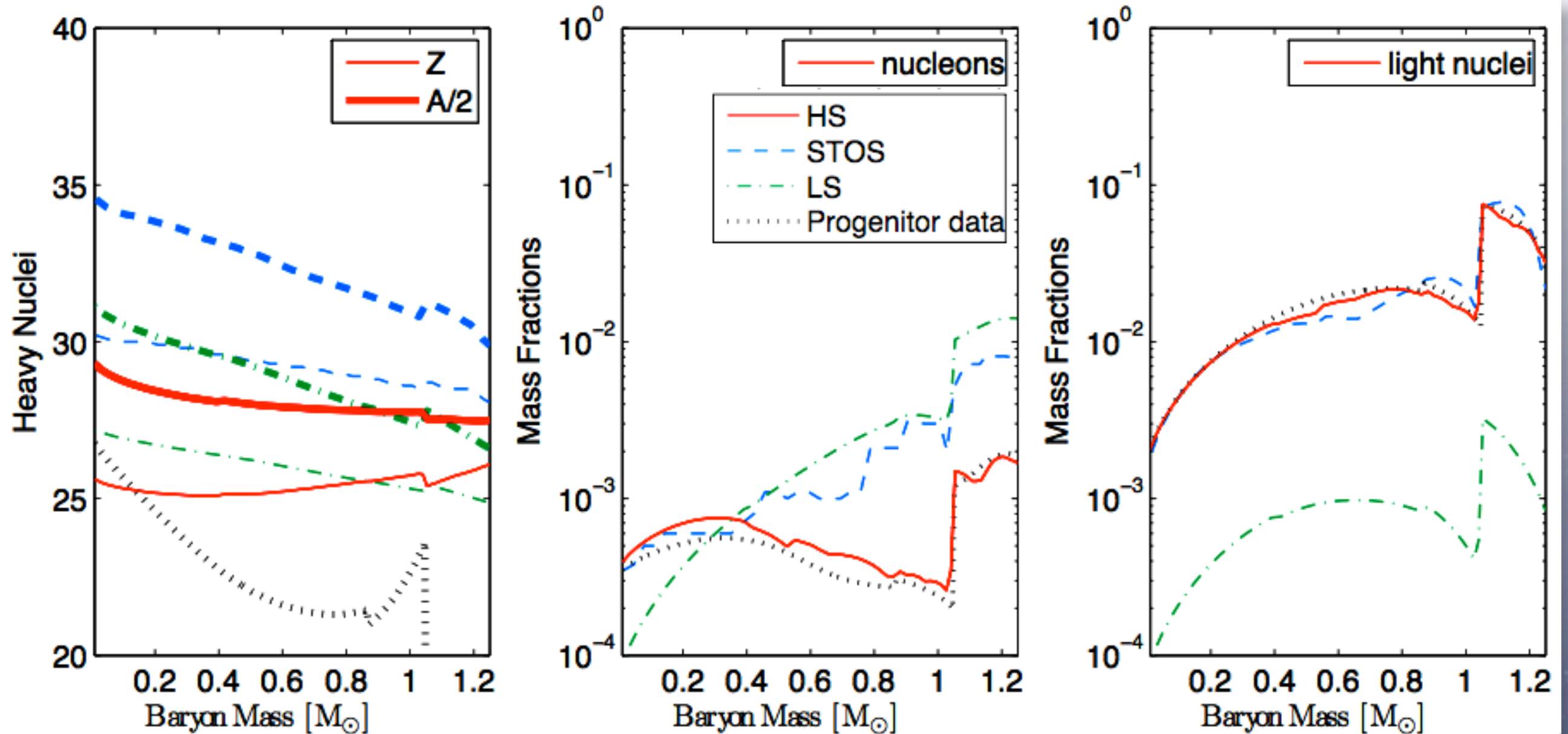
- distinction of light and heavy nuclei by charge  $Z=6$ :

$$X_a = \sum_{A,Z \leq 5} A n_{A,Z} / n_B$$

$$X_A = \sum_{A,Z \geq 6} A n_{A,Z} / n_B$$



# 15 $M_{\text{sun}}$ progenitor - composition



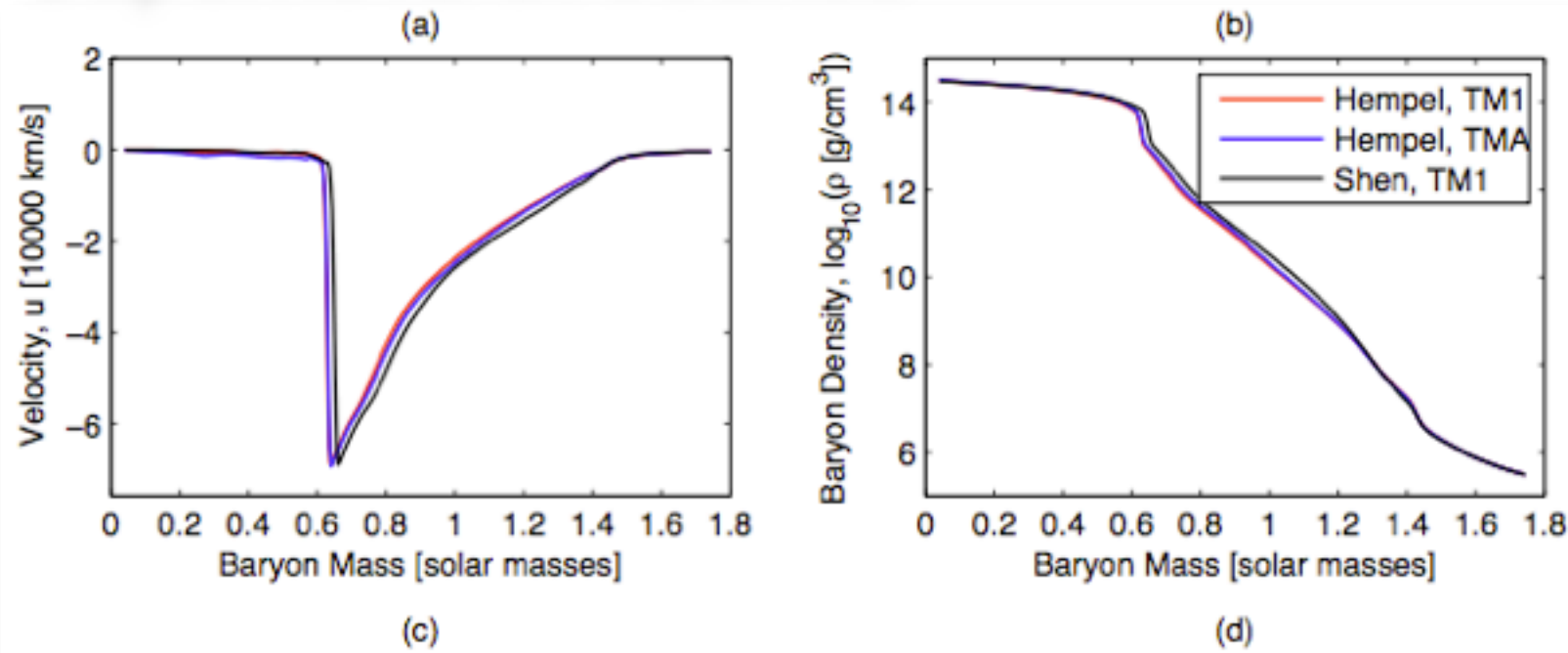
- no iron group nuclei in STOS and LS
- progenitor not in full NSE
- still good agreement of composition in HS

- HS: Hempel, Schaffner-Bielich (TM1)
- STOS: Shen, et al. (TM1)
- LS: Lattimer, Swesty (K=180 MeV)



# Supernova at bounce

15  $M_{\text{sun}}$  progenitor (Woosley & Weaver 1995)

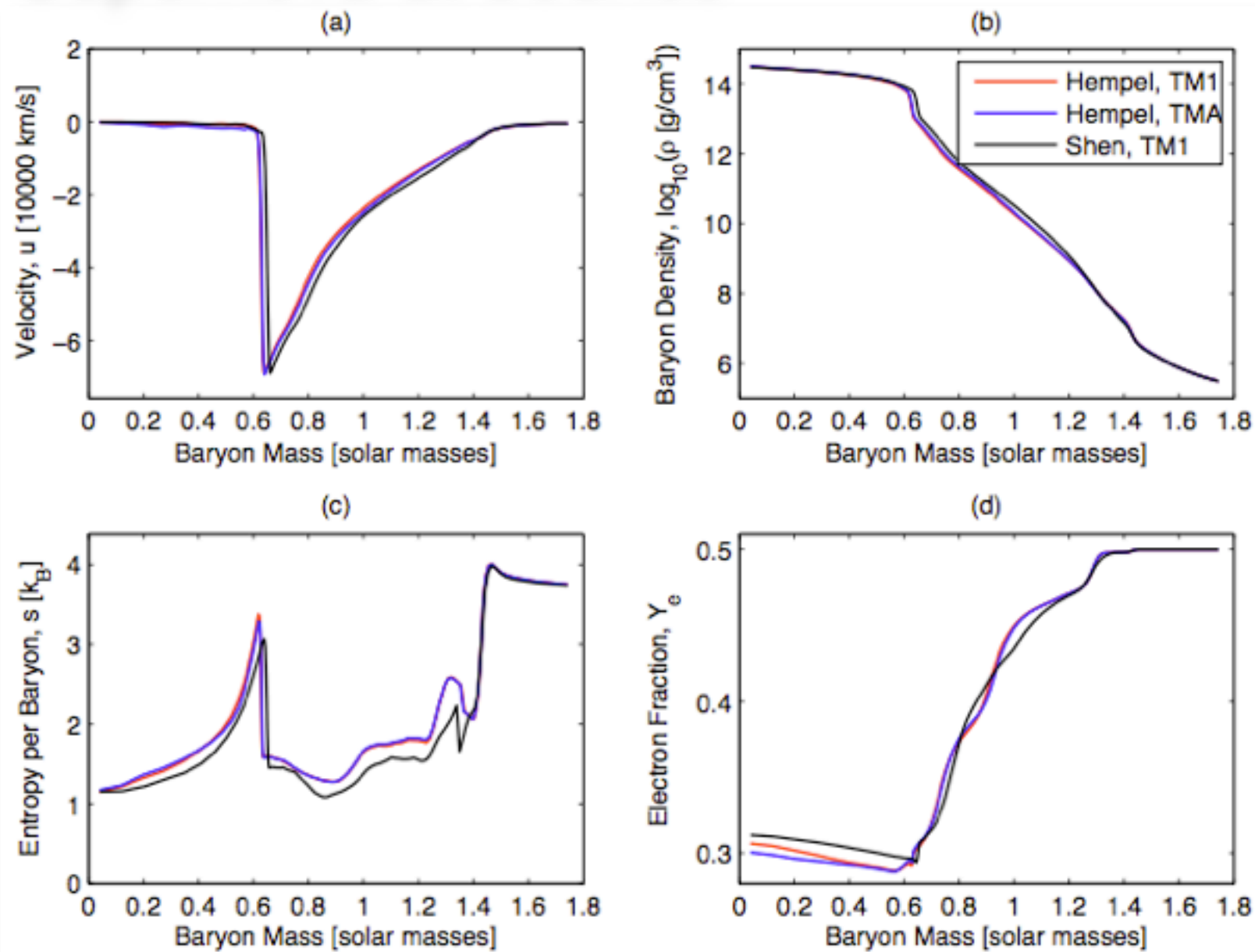


- similar evolution
- slightly more compact PNS core



# Supernova at bounce

15  $M_{\text{sun}}$  progenitor (Woosley & Weaver 1995)

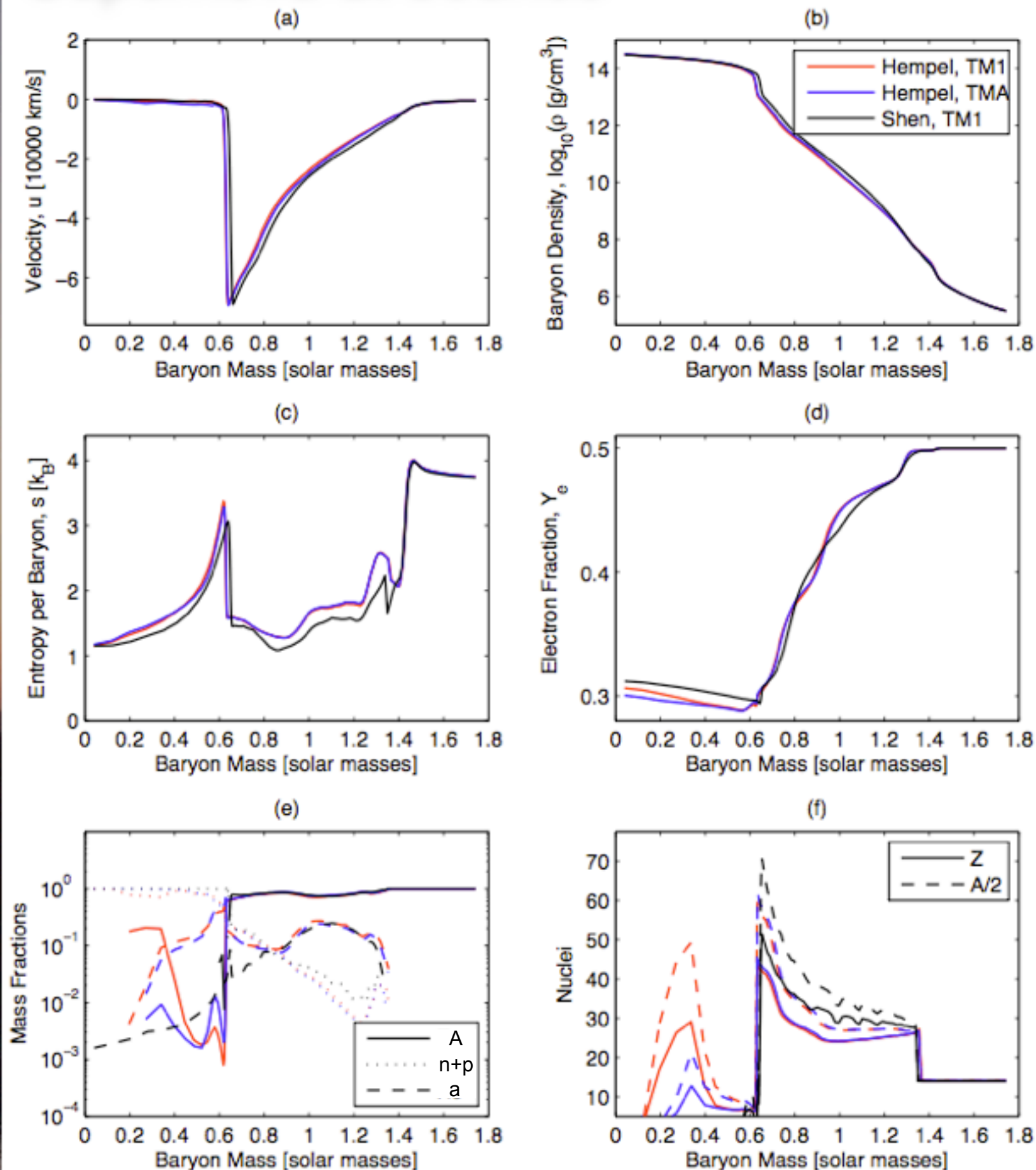


- similar evolution
- slightly more compact PNS core
- entropy difference due to missing kinetic entropy of heavy nuclei in Shen et al.
- slightly more e-captures in HS compared to Shen et al.



# Supernova at bounce

15  $M_{\text{sun}}$  progenitor (Woosley & Weaver 1995)

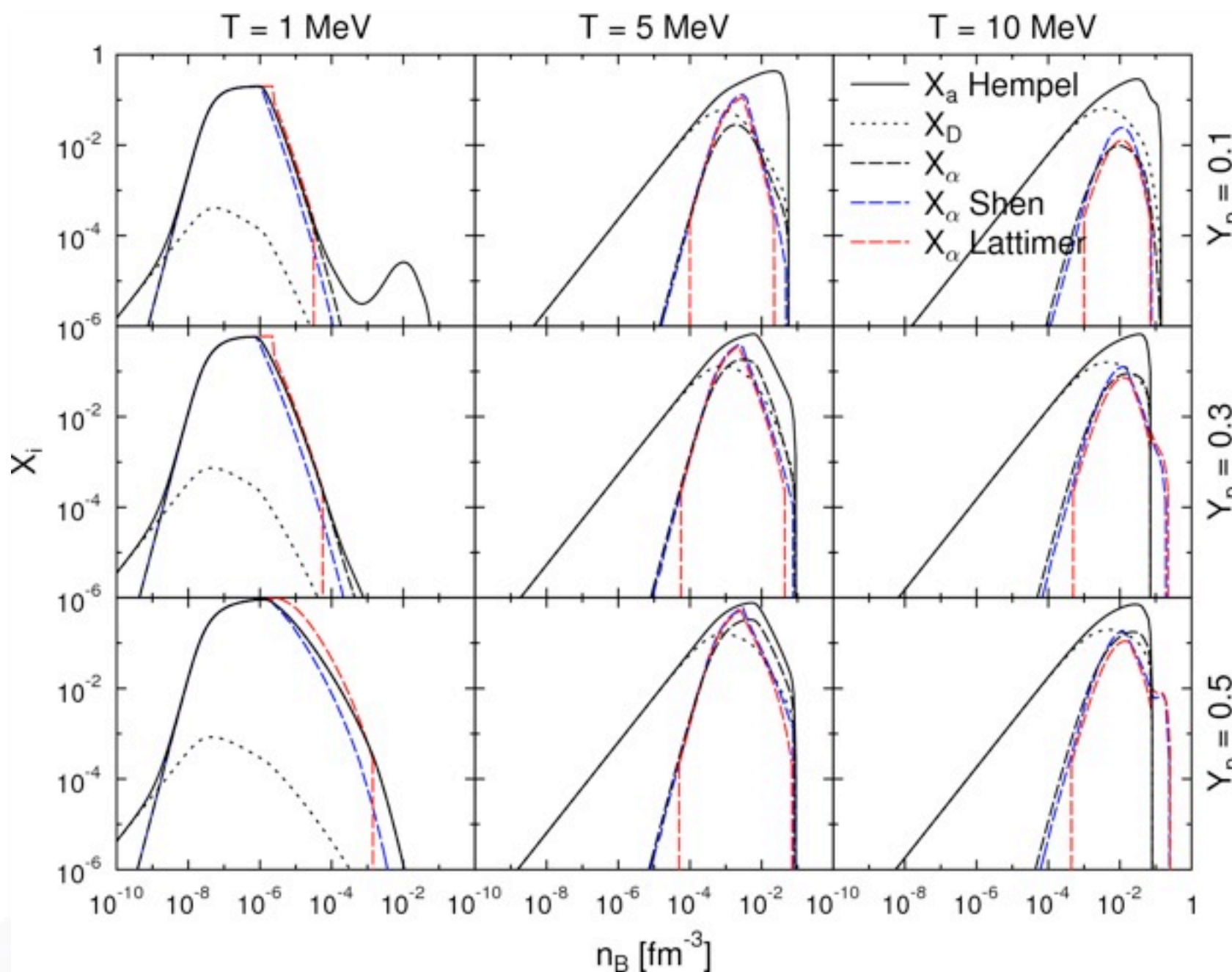


- similar evolution
- slightly more compact PNS core
- entropy difference remains
- slightly more e-captures in HS compared to Shen et al. due to smaller nuclei
- appearance of light clusters behind the shock: deuterons,  $^3\text{H}$ , Li

→ change of the low-density model more important than change TM1 ↔ TMA

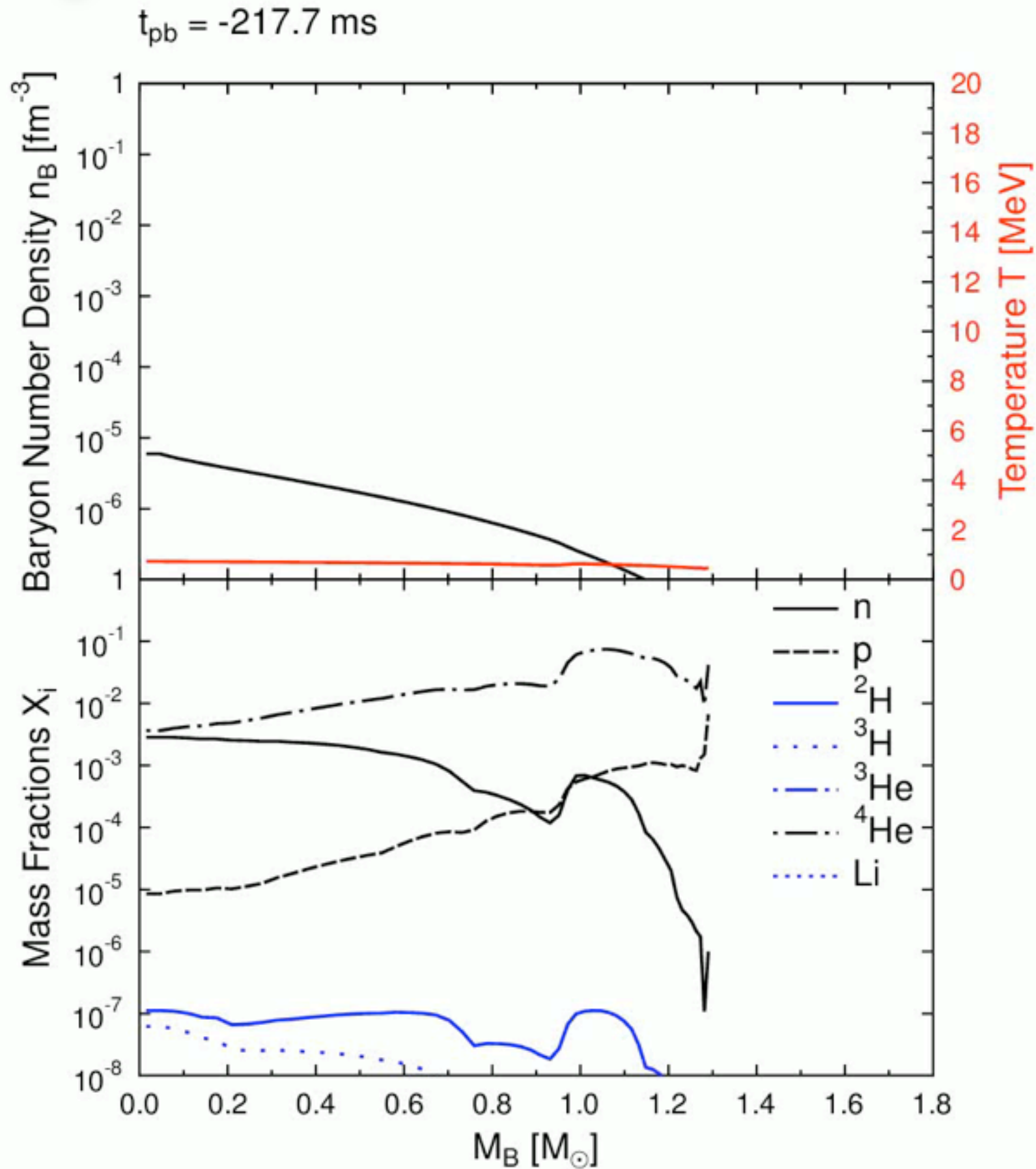


# Light nuclei - comparison with other EOS



- $T = 1$  MeV:  
light clusters well represented by alphas
- $T \geq 5$  MeV:  
large Deuteron contribution at low  $n_B$ , differences for alphas at large  $n_B$
- $n_B \sim 10^{-2} \text{ fm}^{-3}$ :  
whole distribution of light clusters important, not representable by alphas alone

# Light nuclei - animation from simulation

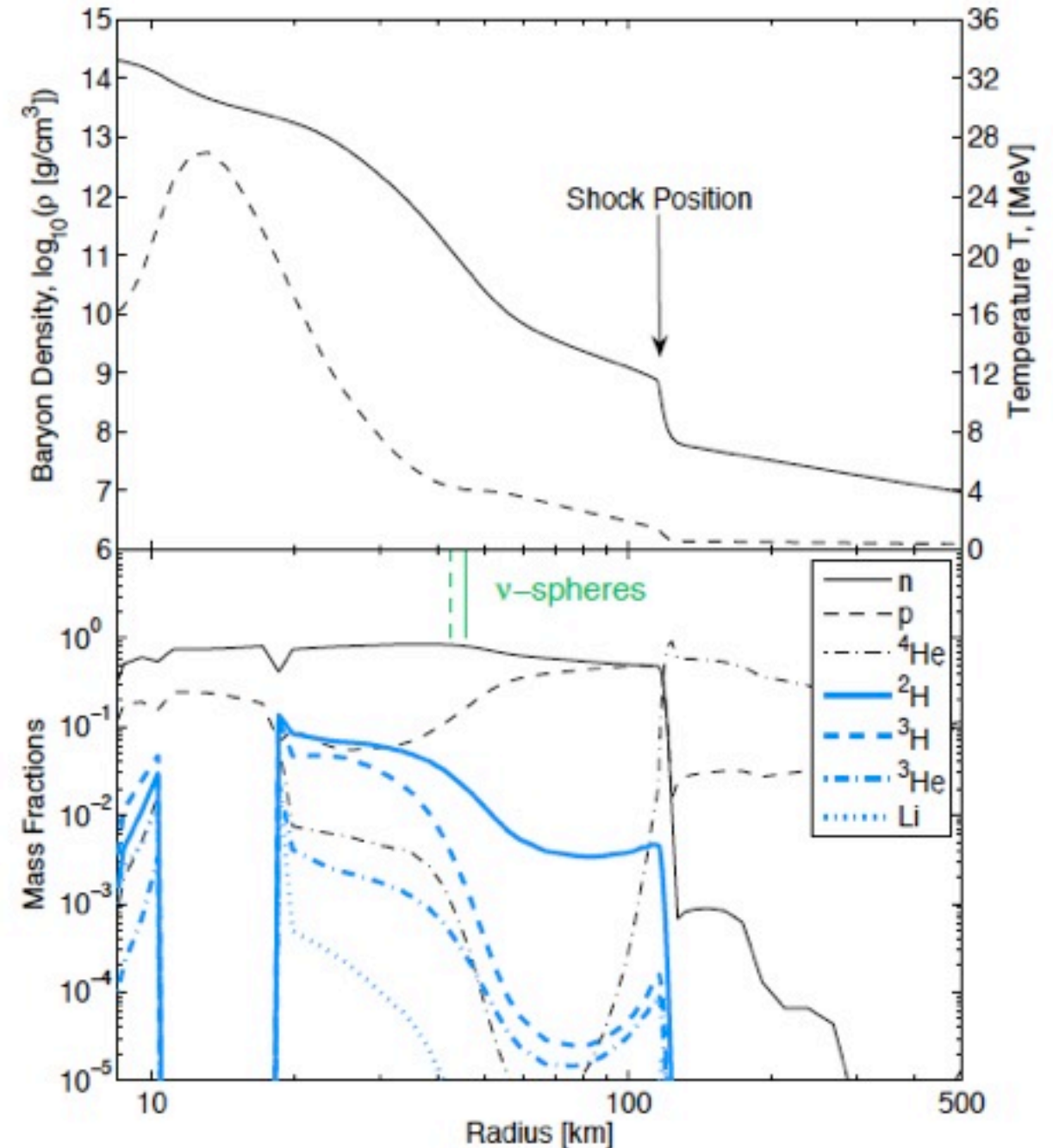
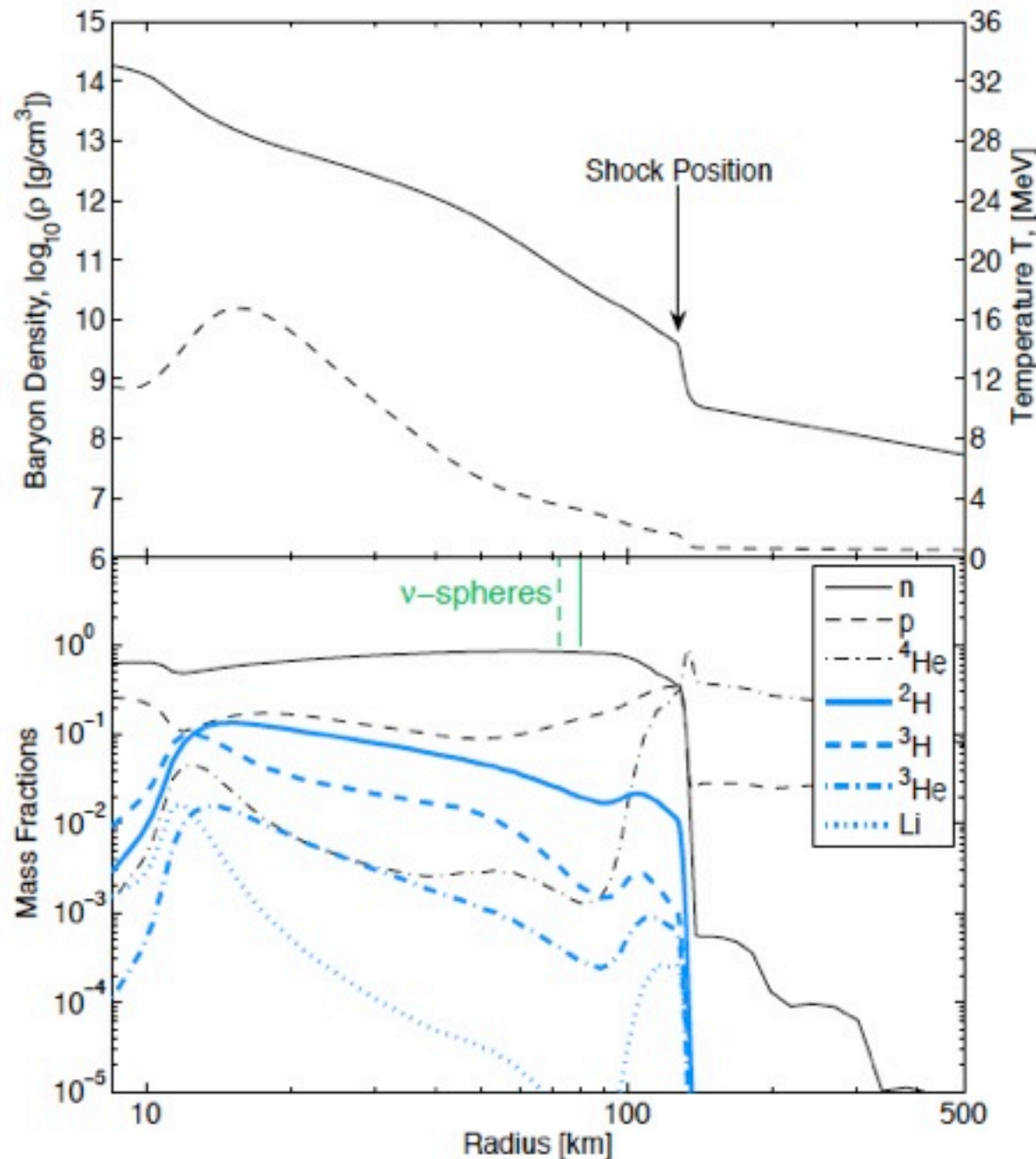




# Light nuclei - postbounce evolution

$t_{pb} = 50$  ms

$t_{pb} = 200$  ms



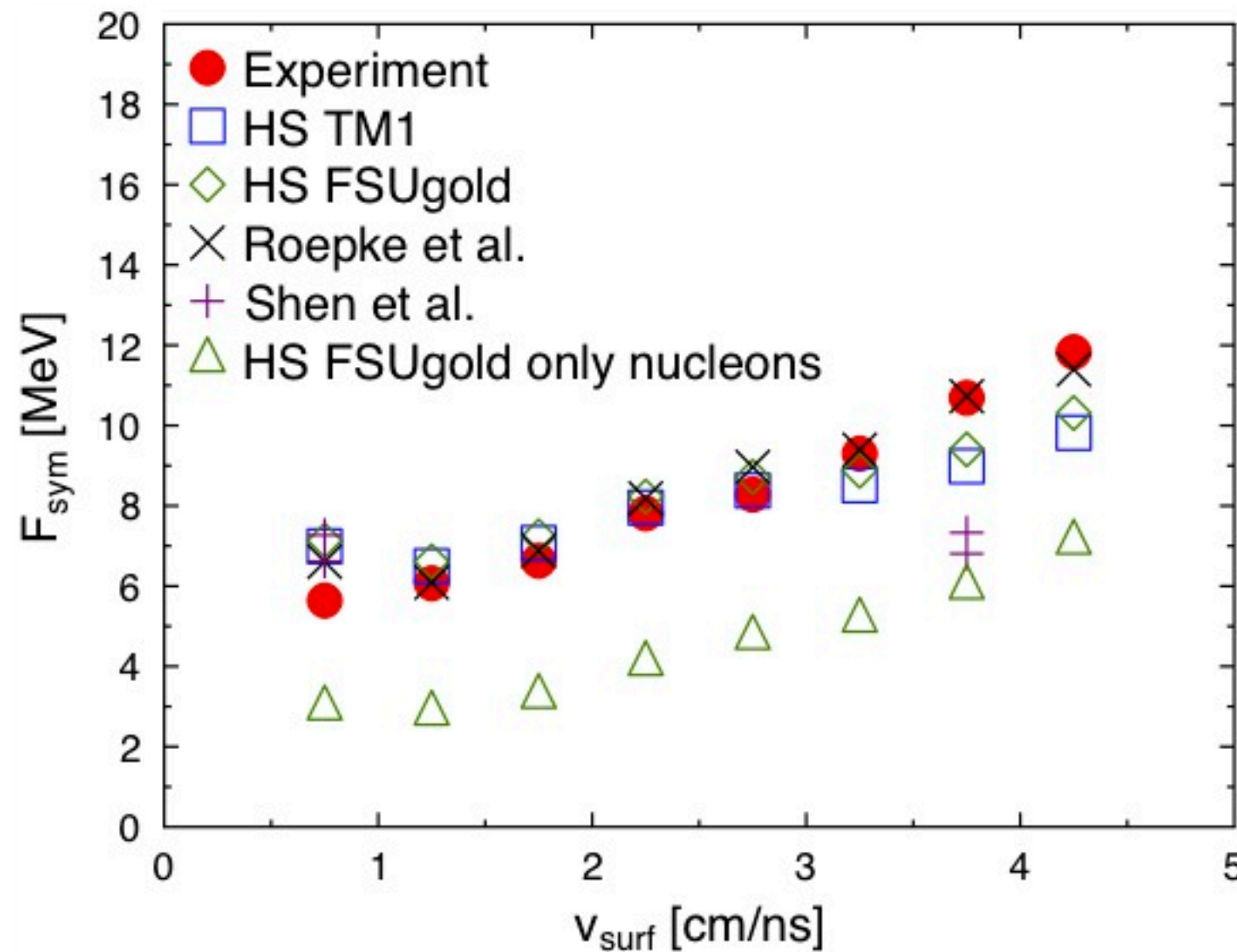
- PNS core: nucleons
- PNS envelope & shock heated matter: light nuclei
- infalling matter: heavies, alphas, nucleons

compare also:  
Sumiyoshi & Röpke PRC77 2008  
Typel et al. PRC81 2010



# Light nuclei - measurement of symmetry energy

based on: [Natowitz et al.; 2010PRL104]



- symmetry energy extracted from low-energy heavy ion collisions
- thermodynamic conditions:  $T = 3 - 7$  MeV and  $n_B = 1/100 - 1/20 n_B^0$
- cluster formation leads to increased symmetry energy

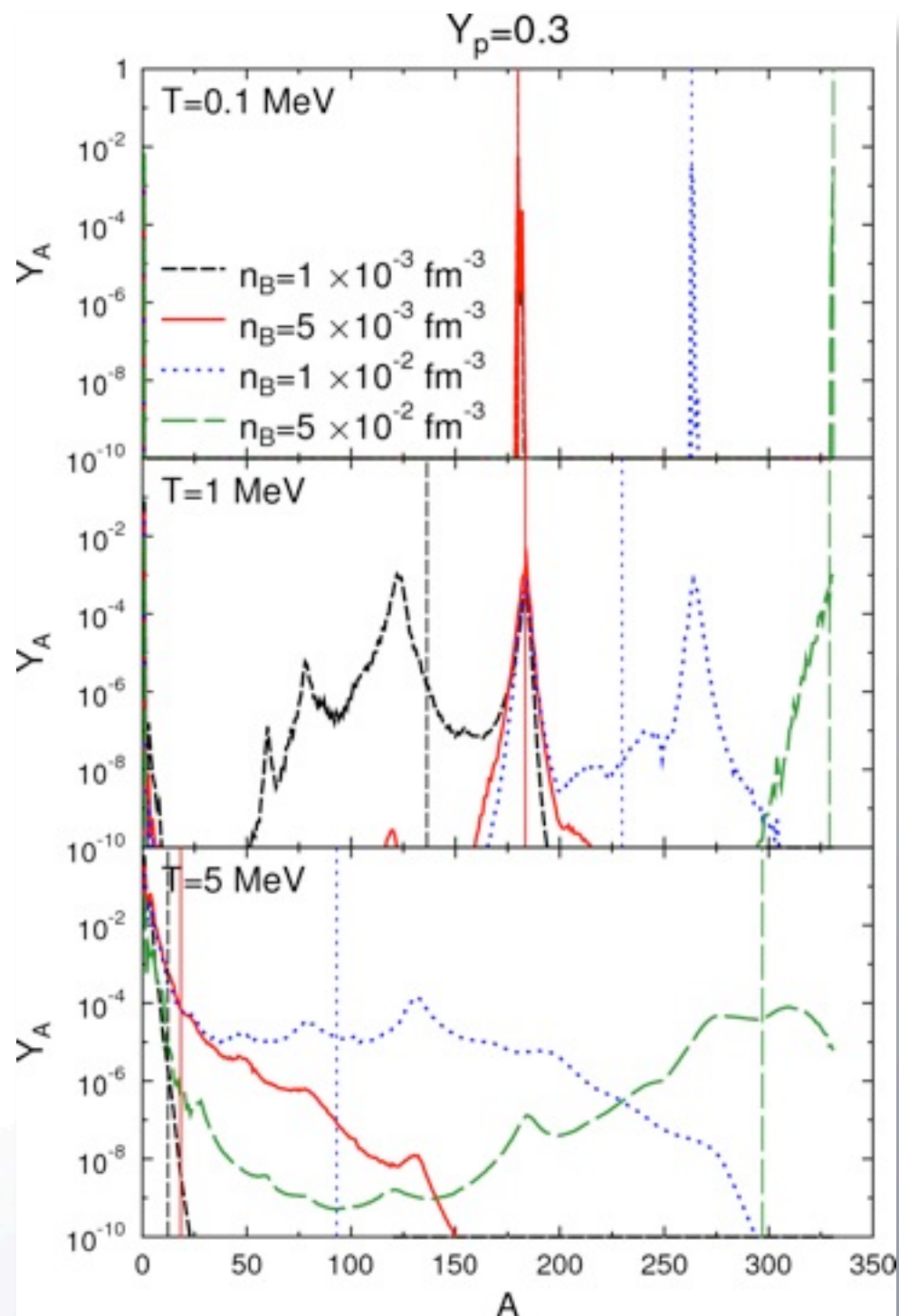
→ experimental evidence for appearance of light clusters in SN

T [MeV]	3.3	3.3	3.6	4.2	4.7	5.3	6.2	7.5
$n_B$ [ $10^{-3} \text{ fm}^{-3}$ ]	2.1	1.7	2.3	3.8	4.7	4.9	5.5	6.4

$$E_{\text{sym}}(n_B, T) = 1/2 ( E_{\text{sym}}(n_B, T, Y_p = 1) + E_{\text{sym}}(n_B, T, Y_p = 0) ) - E_{\text{sym}}(n_B, T, Y_p = 0.5)$$

[Typel et al.; 2010PRC81]  
[Kowalski et al.; 2007PRC75]

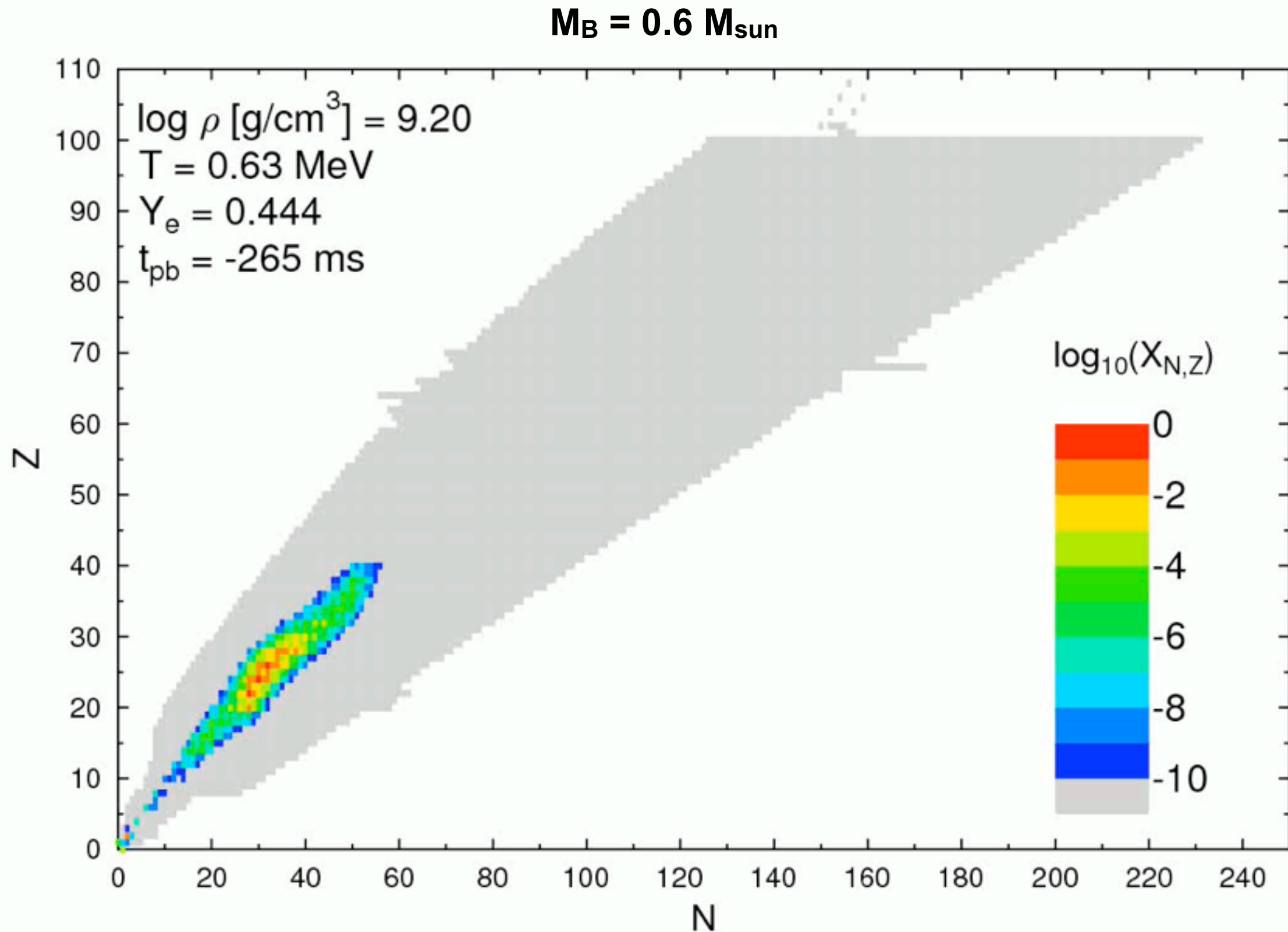
# Nuclear distributions



- $T=0.1 \text{ MeV}$ : only one nucleus
- peaks caused by neutron shell effects
- neutron magic numbers (40), 50, 82, 126, 184
- $T=1 \text{ MeV}$ : shell effects leading to narrow distributions with multiple peaks
- $T \geq 5 \text{ MeV}$ :  
broad distributions, weakening of shell effects

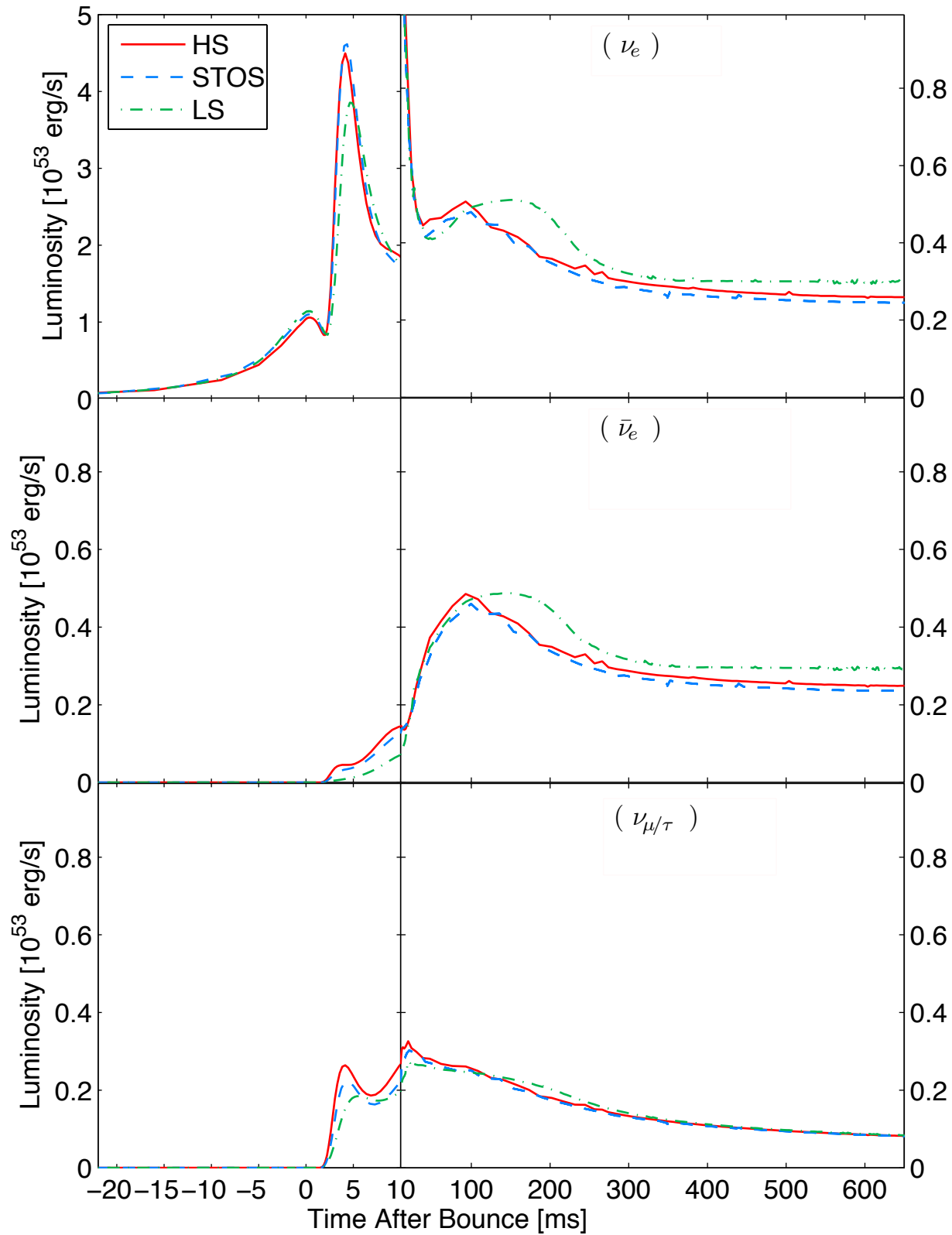


# Nuclear distributions in the SN





# Neutrino signal - 15 M<sub>sun</sub> progenitor

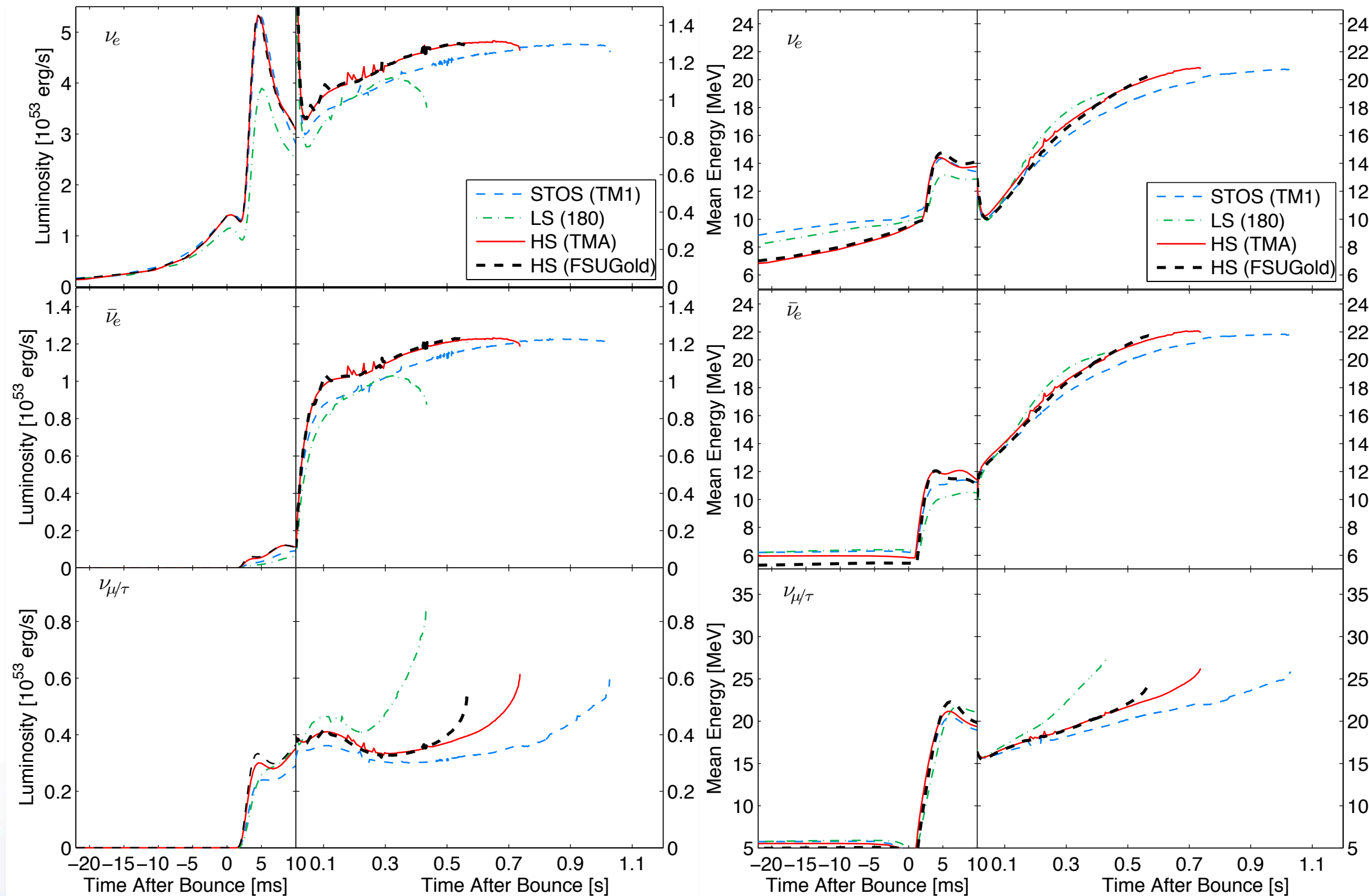


- STOS TM1 and HS TM1: negligible differences
  - softer LS EOS: hotter PNS, larger luminosities
- only slight impact of the low density EOS on the neutrino signal

# Results - high density

- simulations for  $40 M_{\text{sun}}$  progenitor of Woosley & Weaver ApJS 101 (1995)
- set of EOS tables: STOS, LS 180, LS 220, HS (TMA), HS(FSUgold)

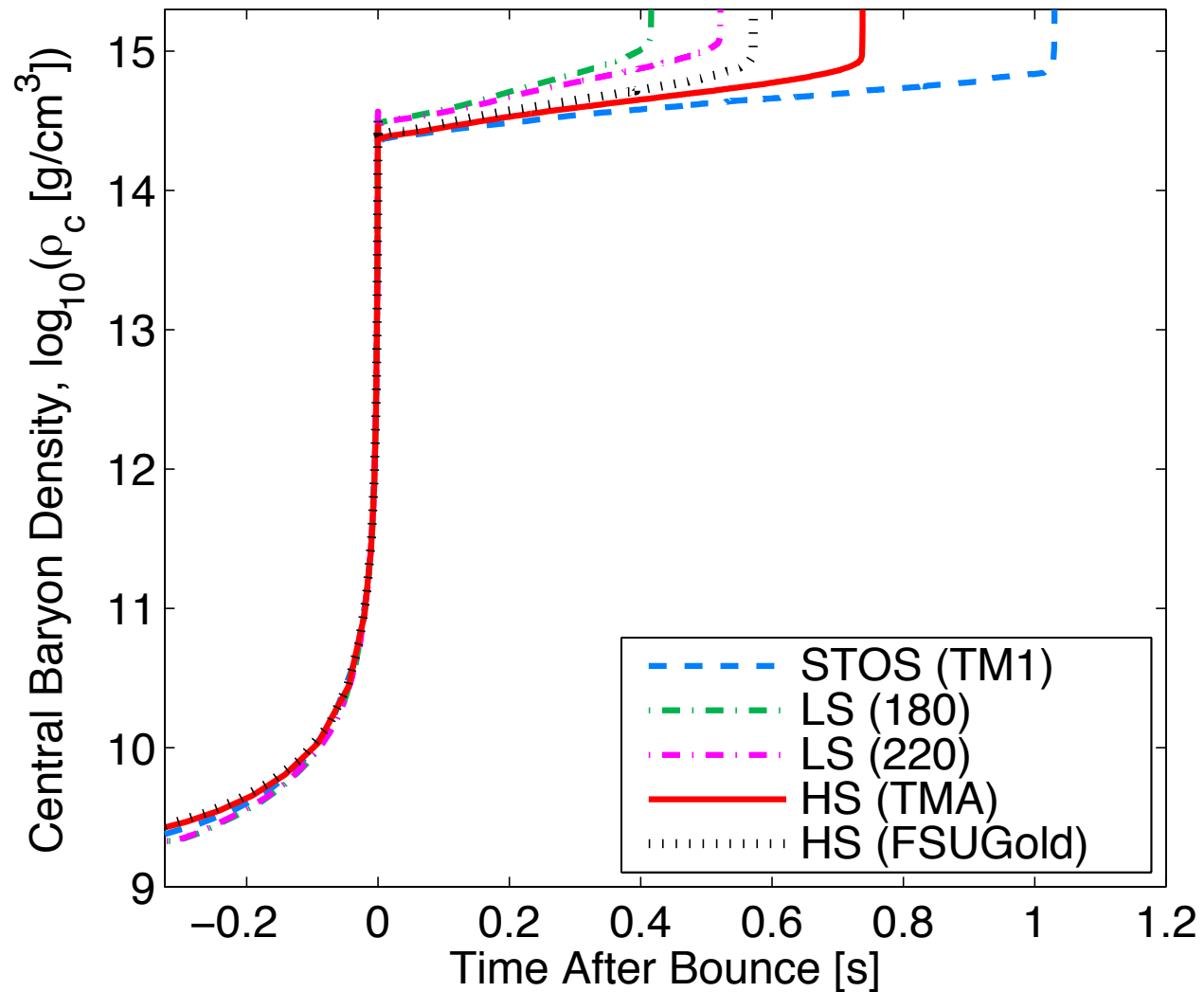
# Neutrino signal - 40 $M_{\text{sun}}$ progenitor



- different times until black hole formation
- $\mu/\tau$ -neutrinos most sensitive to EOS because emitted from deeper layers



# 40 $M_{\text{sun}}$ progenitor - time until black hole formation



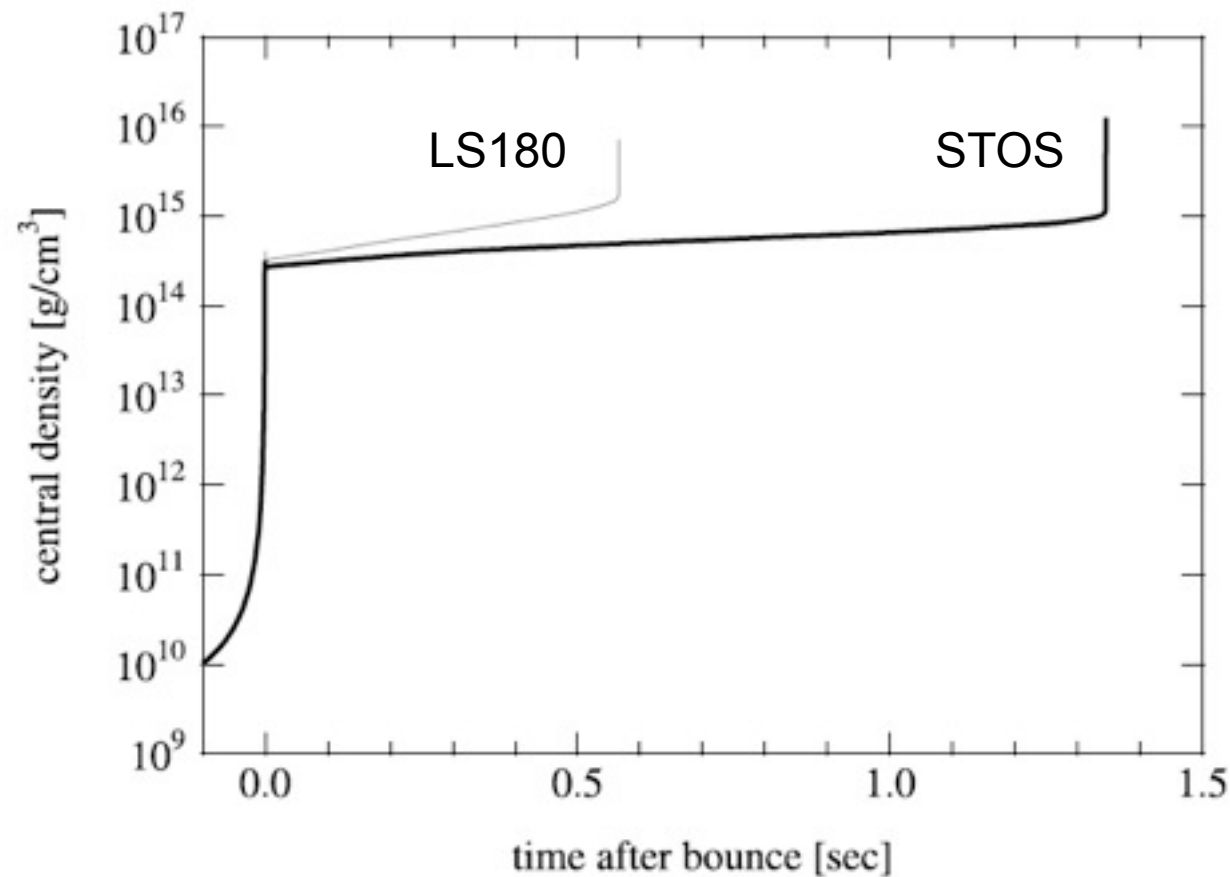
- surprising result: maximum mass of cold neutron stars is not directly correlated with  $t_{\text{BH}}$
- non-relativistic LS collapse earlier than expected
- note: accretion rate not much affected by EOS

	$M_{\text{max}}(T=0)$ [ $M_{\text{sun}}$ ]	$t_{\text{BH}}$ [ms]
TM1	2.22	1028
TMA	2.02	737
FSUGold	1.74	571
LS220	2.06	521
LS180	1.83	415

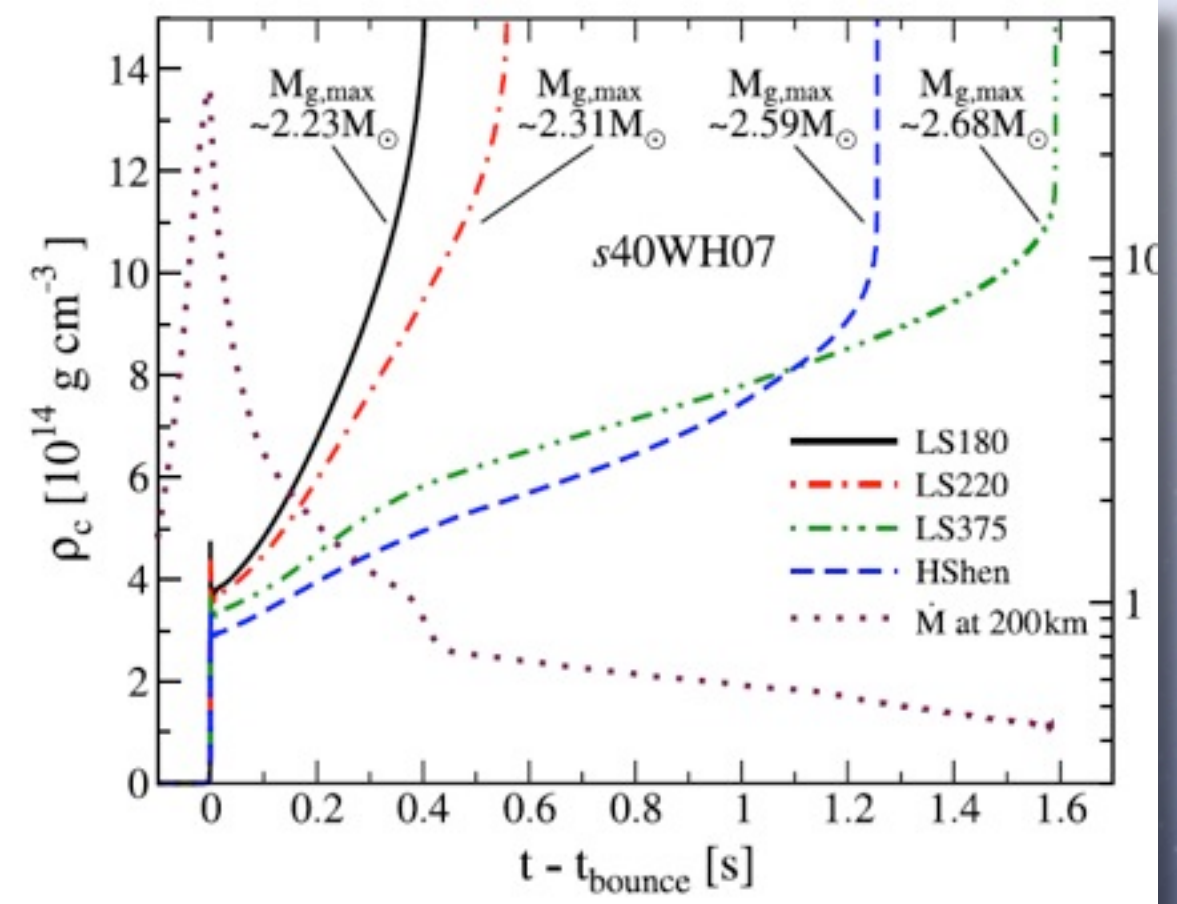
# 40 $M_{\text{sun}}$ progenitor - time until black hole formation

- this was not found in previous studies, because only STOS, the stiffest RMF EOS with TM1 was available:

[Sumiyoshi et al.; 2007ApJ667]

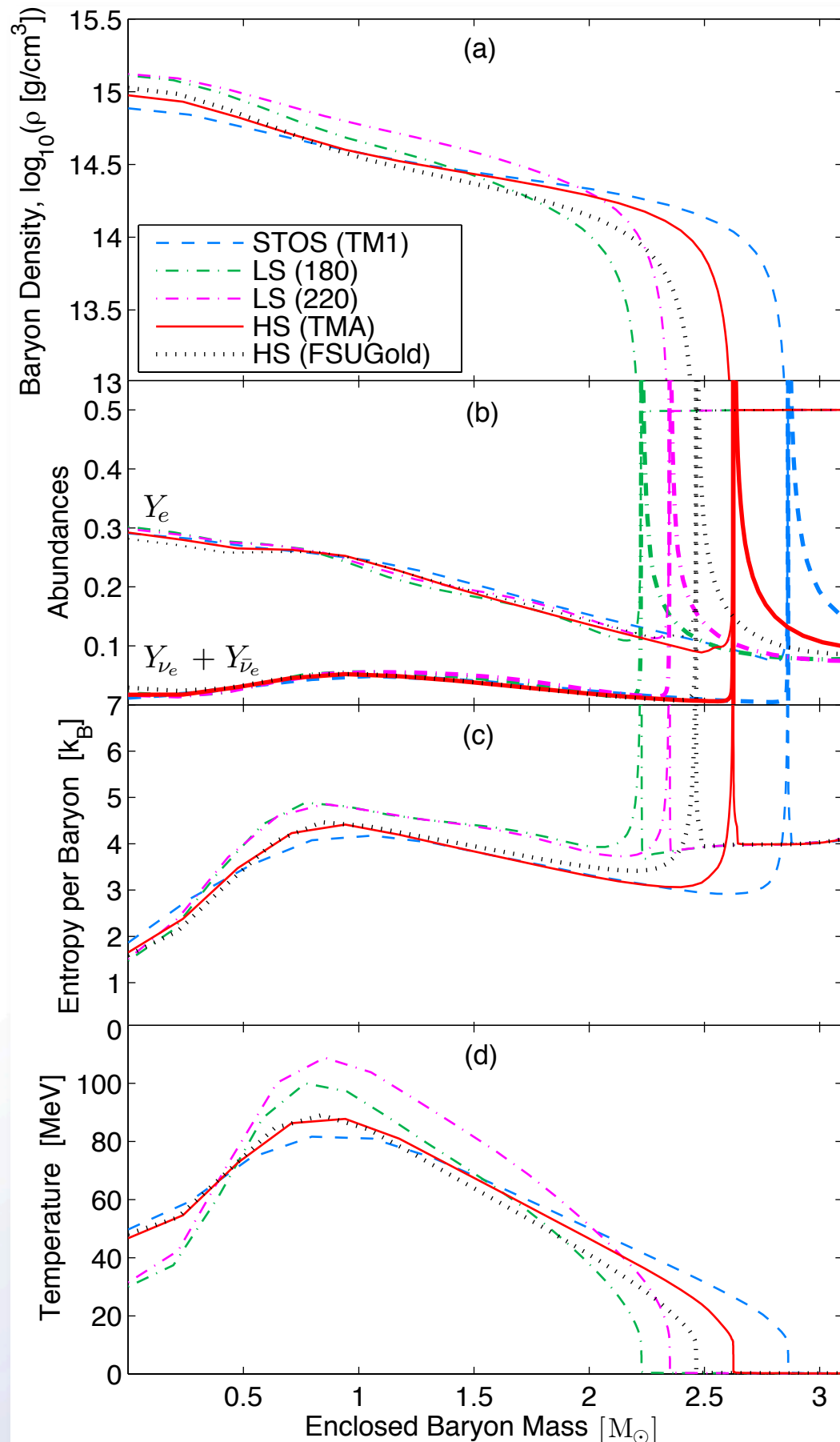


[O'Connor & Ott.; 2011ApJ730]



- why do the non-relativistic LS EOS collapse so early, respectively why do they behave so soft in the CCSN?

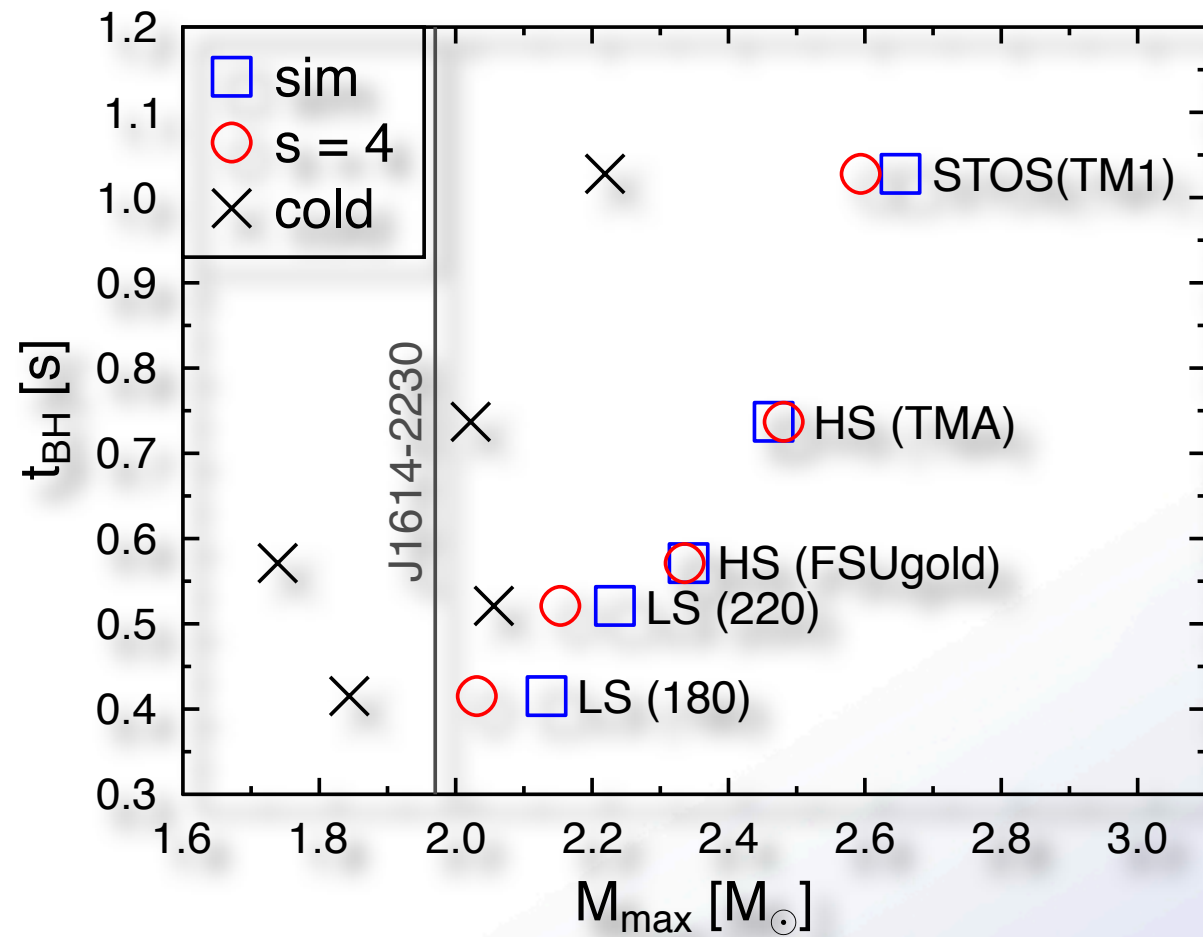
# 40 $M_{\text{sun}}$ progenitor - state before black hole formation



- similar configurations of the five models at the state before the collapse to black hole
- main differences in the temperature profiles
- idea: use  $s=4 k_B$ , beta-equilibrium and the TOV equations to analyze the maximum mass



# 40 $M_{\text{sun}}$ progenitor - correlation of $t_{\text{BH}}$ with maximum mass



- the maximum gravitational masses found in the simulation are significantly enhanced (up to  $0.6 M_{\text{sun}}$ ) compared to  $T=0$
- the  $s=4$  configuration reproduces the results of the simulations

- the LS EOS show a significantly weaker stiffening by temperature than the RMF EOS
- due to: effective mass, symmetry energy, non-relativistic description
- new correlation: the time until black hole formation gives information about the finite entropy EOS and the maximum mass of proto-neutron stars!

# Conclusions

- phenomenological, thermodynamic consistent model for an ensemble of nuclei in equilibrium with an interacting nucleon gas
- EOS tables and routines for composition available for NL3, TM1, TMA, FSUgold, DD2:  
<http://phys-merger.physik.unibas.ch/~hempel/eos.html>
- results from SN simulations (low densities):
  - broad distributions of heavy nuclei during collapse
  - large yields of light clusters in the shock heated matter
  - notable impact of the low-density EOS model→ role of light nuclei and distributions on neutrino transport?
- high densities:
  - SN of massive progenitors give information about the finite entropy EOS
  - temperature effects are EOS dependent→ extract information about nucleon effective mass or symmetry energy?



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