

Neutrino astrophysics

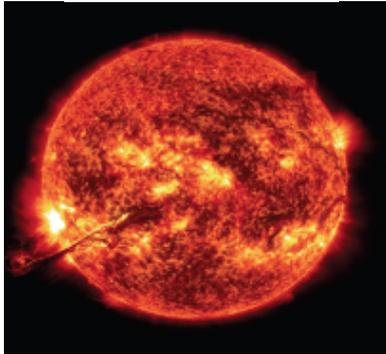
Maria Cristina VOLPE

AstroParticule et Cosmologie (APC)
Paris, FRANCE

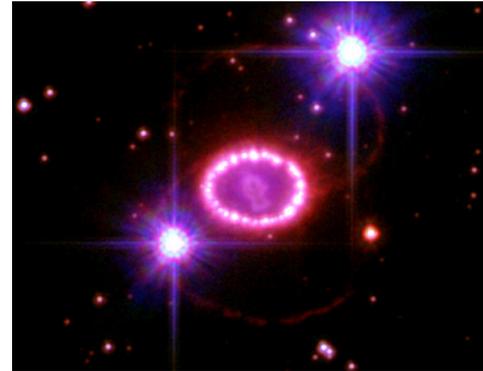
Outline

- ❖ Introduction
- ❖ Theoretical aspects of flavor evolution, recent developments and open issues
- ❖ Observational aspects of supernova ν
- ❖ Flavor evolution in binary neutron star merger (BNS) remnants
- ❖ Conclusions

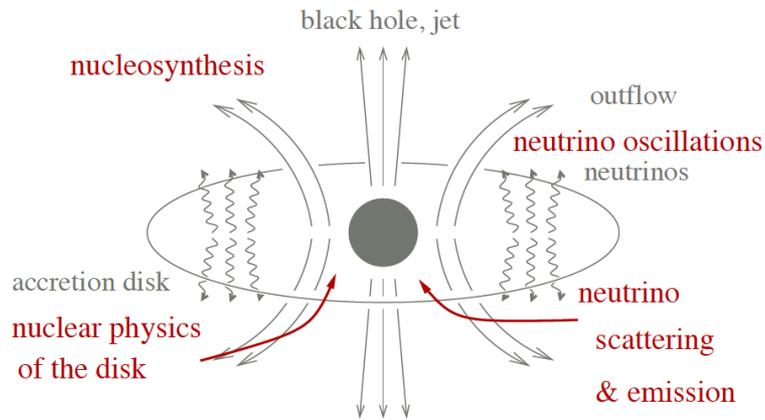
the Sun



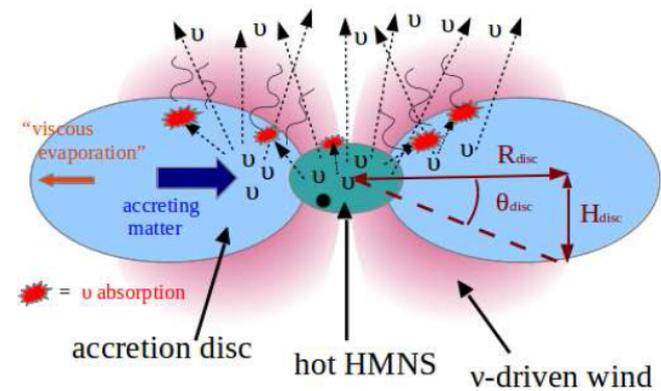
core-collapse Supernovae



accretion disks around black holes or neutron star mergers remnants



from McLaughlin

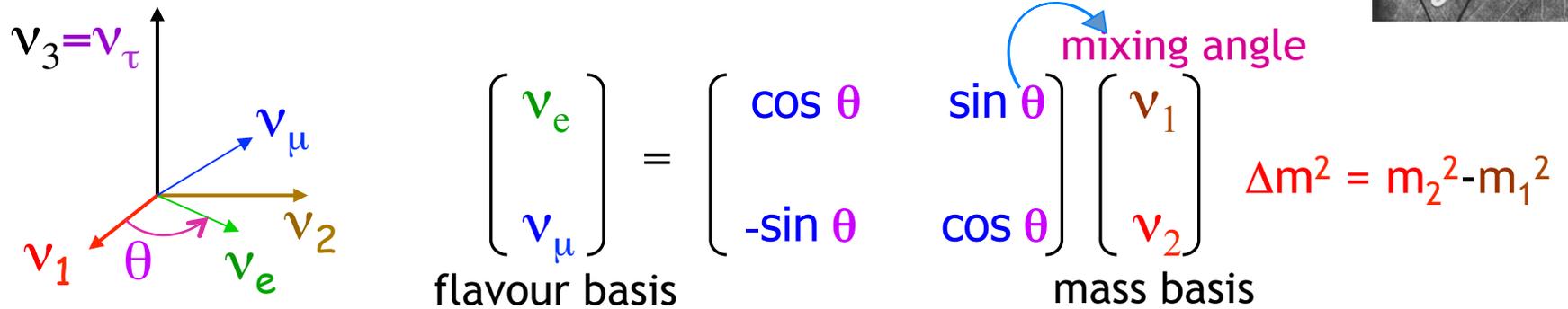


from Perego, A. 2014

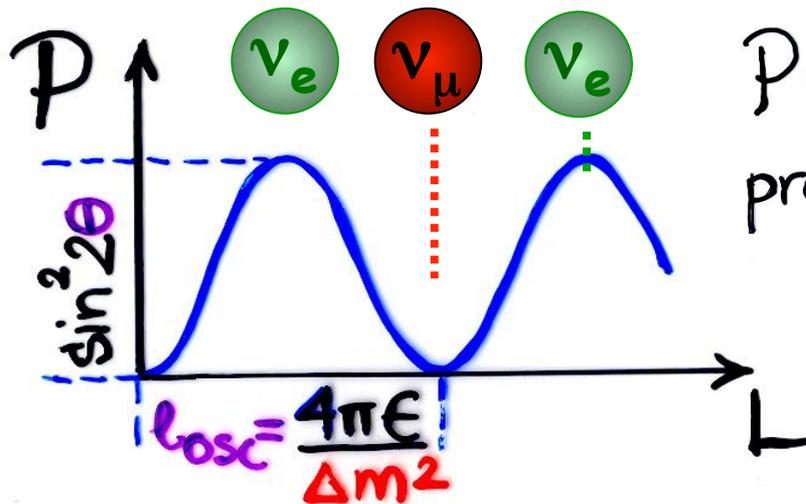
The oscillation phenomenon in vacuum

Neutrino oscillations occur if the neutrino interaction and propagation basis do not coincide, analogy with $K_0-\bar{K}_0$ systems.

Pontecorvo JETP 6 (1957)



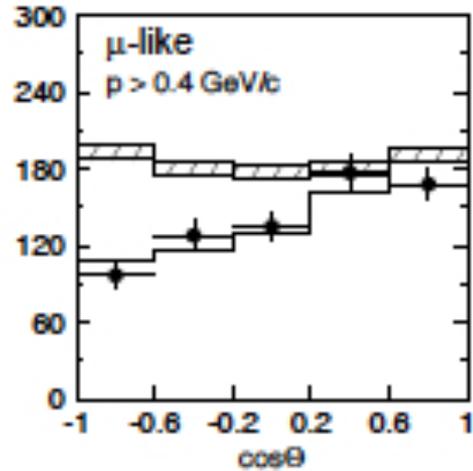
Time evolution : $|\nu_e(t)\rangle = \cos \theta \cdot e^{-iE_1 t} |\nu_1\rangle + \sin \theta e^{-iE_2 t} |\nu_2\rangle$



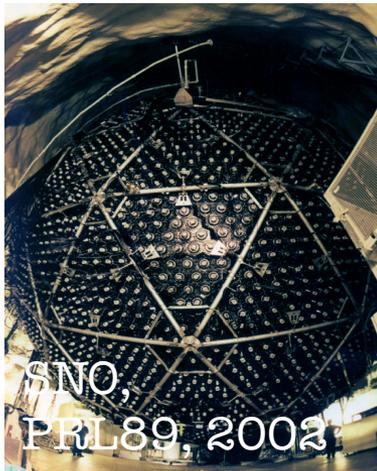
$P(\nu_e \rightarrow \nu_\mu) = \sin^2 2\theta \sin^2\left(\frac{L}{4E} \Delta m^2\right)$
probability for neutrino oscillations

AN INTERFERENCE
PHENOMENON

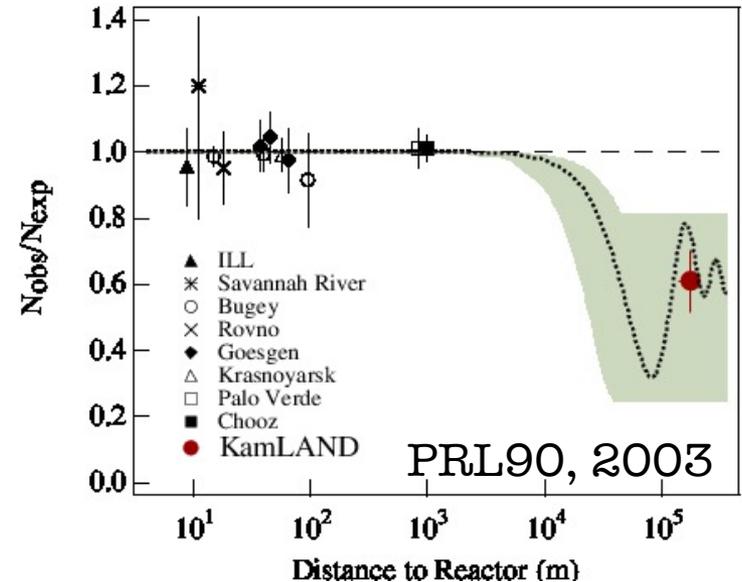
The neutrino oscillation discovery



Neutrinos : elementary massive particles with non-zero mixings



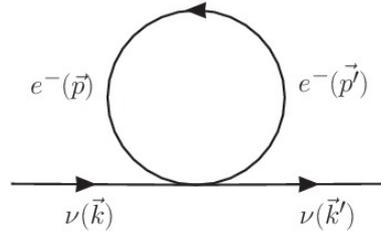
Total solar neutrino flux consistent with Standard Solar Model



Mikheev-Smirnov-Wolfenstein large-mixing-angle solution

The Mikheev-Smirnov-Wolfenstein effect

Resonant flavour conversion due to ν -matter interactions



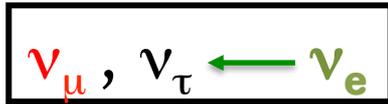
$$H_{\nu_e}(\rho_e) = \sqrt{2}G_F \rho_e$$

MEAN-FIELD

electron
number density

Wolfenstein PRD (1978)

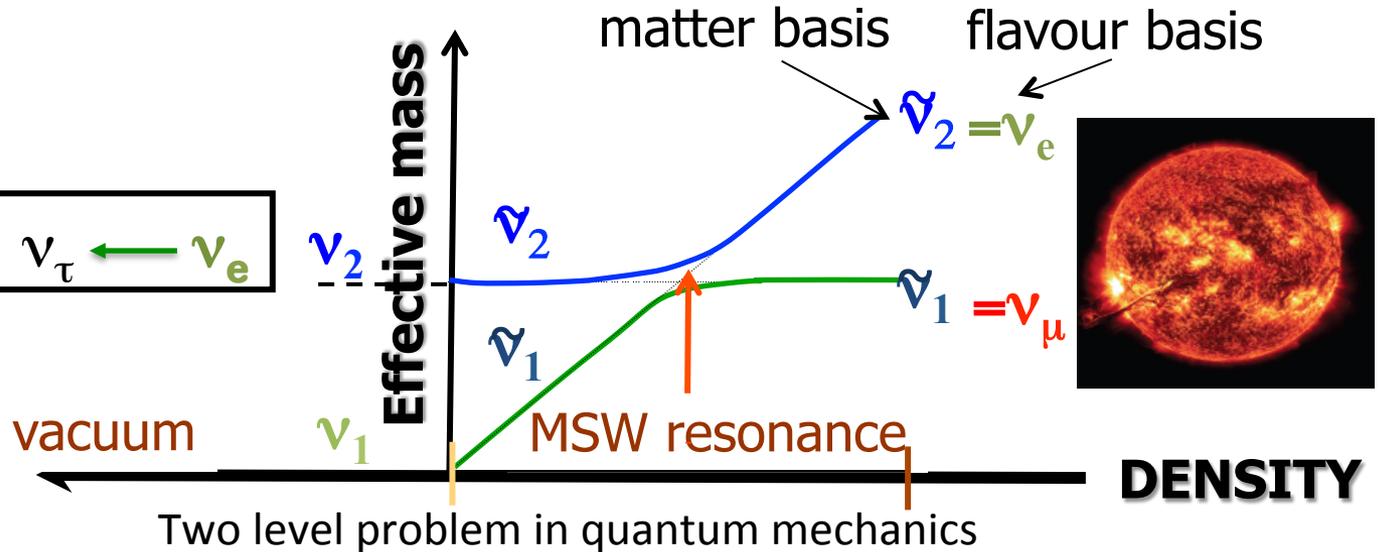
Mikheev, Smirnov(1985)



Effective mass

matter basis

flavour basis



In the Sun : evolution through the resonance adiabatic,
sign of one of the Δm^2 determined

Also in supernovae, in accretion disks around compact objects,
in the Earth and in the Early Universe (BBN epoch)

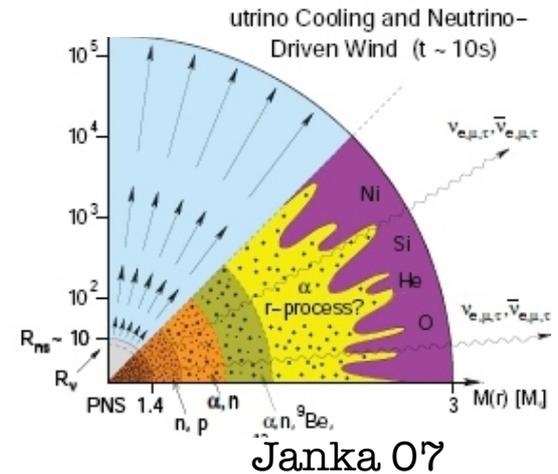
Neutrinos in dense environments

Linked to the key issues

□ How do massive stars explode ?

Current simulations :

multidimensional, realistic ν transport,
convection and turbulence, hydrodynamical instabilities (SASI).



□ What are the sites where heavy elements are made ?

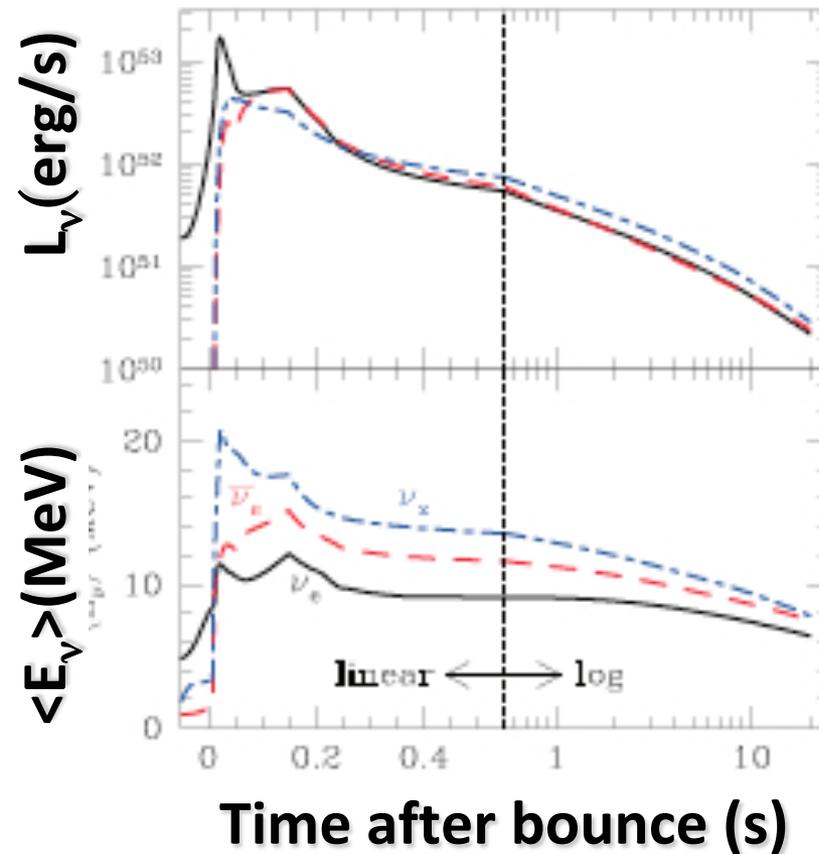
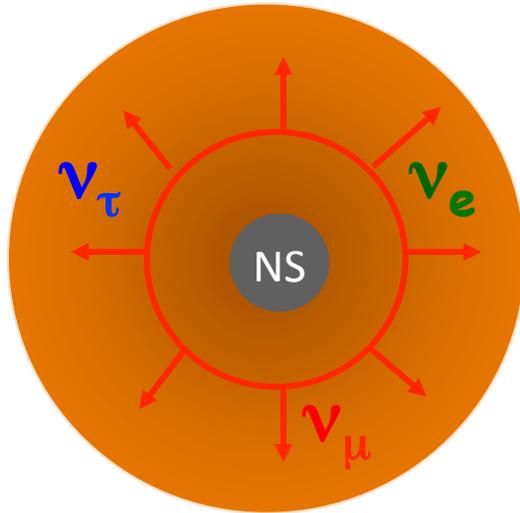
Candidate sites for heavy elements nucleosynthesis :

supernovae, accretion disks around compact objects (BH, NS-NS, BH-NS)

see e.g. Focus issue «Nucleosynthesis and the role of neutrinos»,
J.Phys. G 41 (2014)

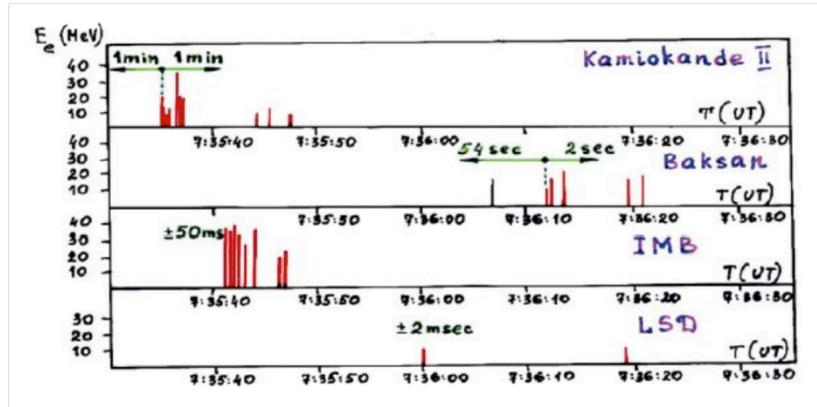
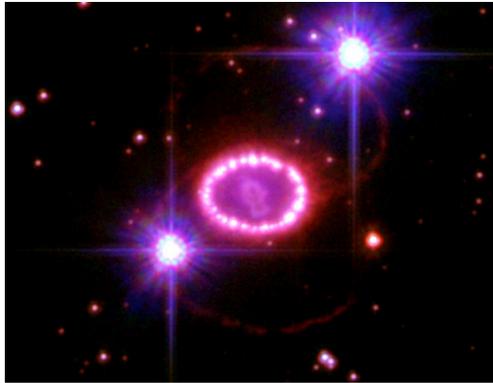
Neutrinos from gravitational core-collapse

99 % of the gravitational energy (10^{53} ergs) is radiated as neutrinos and anti-neutrinos of all flavours.



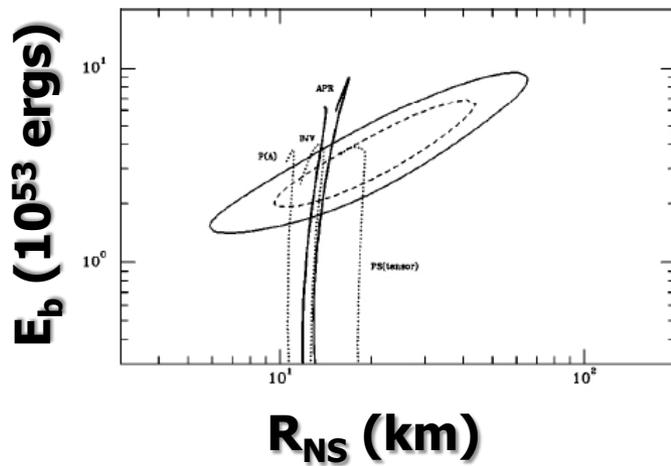
Nakazato et al, arXiv : 1210.6841

SN1987A

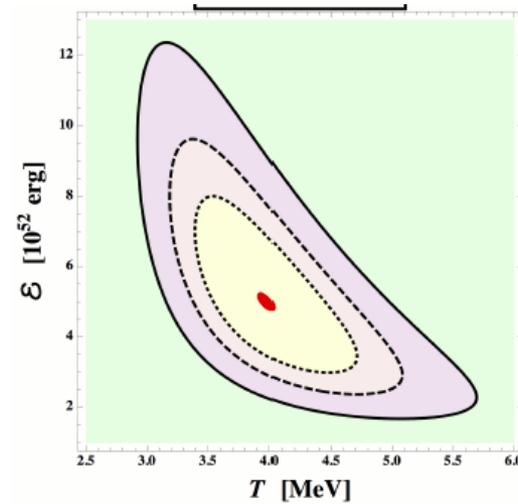


Suzuki, A. J. Conf. Phys. 120 (2008)

Sanduleak 69⁰202, a blue super-giant in Large Magellanic Cloud, at 50 kpc, no remnant found so far. Hirata et al, PRL58(1987)



Loredo, Lamb, PRD65 (2002)

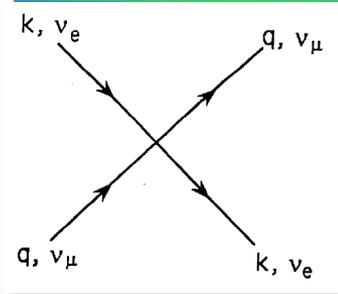
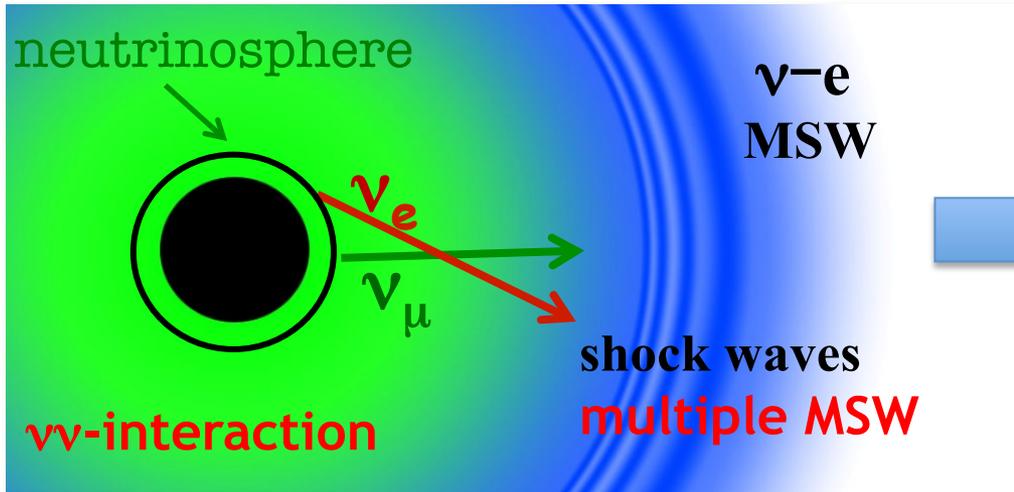


Vissani, JPG42 (2014)

SN1987A : Delayed explosion mechanism favored over the prompt one.

Flavor conversion in supernovae

Novel phenomena uncovered

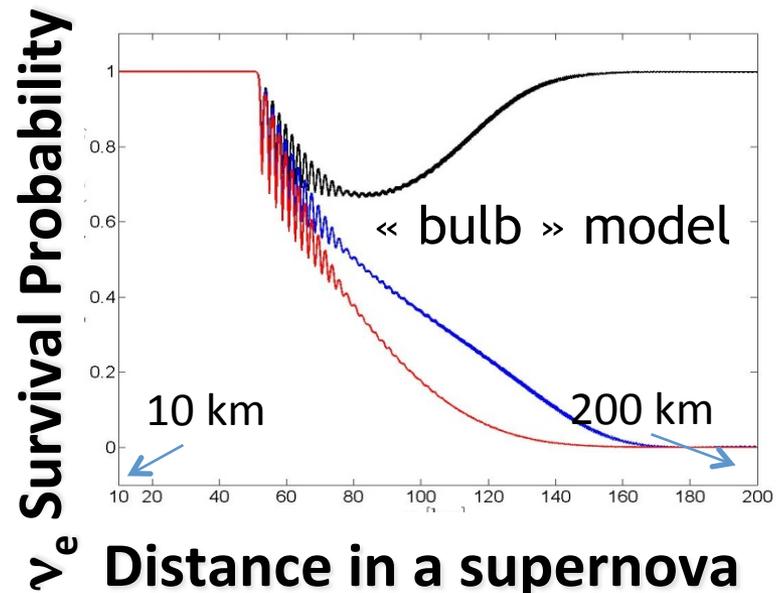


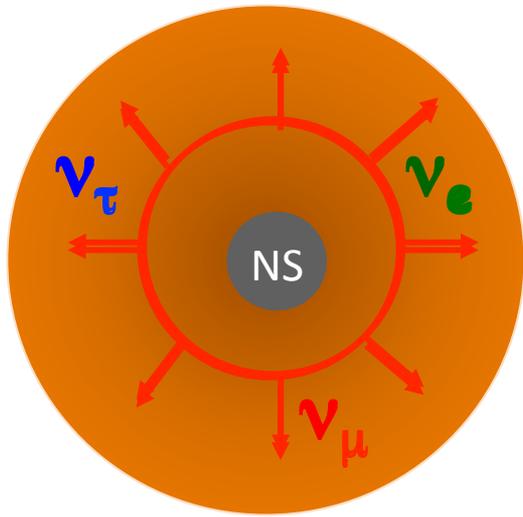
Pantaleone, PLB 1992

- Efficient conversion occurs close to the neutrinosphere

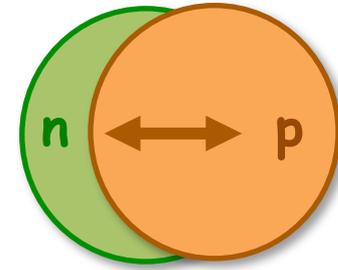
Models of increasing complexity studied since a decade

Duan, Fuller, Qian, Ann. Rev. 60 (2010)





ν in stars or accretion disks



atomic nucleus

10^{57}

N

200

weak

interaction

strong

unbound

system

bound

$$\rho_{ji} = \langle a_i^+ a_j \rangle \text{ neutrinos}$$

$$\bar{\rho}_{ji} = \langle b_i^+ b_j \rangle \text{ anti-neutrinos}$$

density

$$\rho = \langle a^+ a \rangle \text{ neutrons}$$

$$\text{protons}$$

Neutrino flavor evolution in dense environments :
a many-body problem

To determine the dynamics exactly:

$$\begin{array}{cccc} \rho_1 = \langle a^\dagger a \rangle & \rho_{12} = \langle a^\dagger a^\dagger a a \rangle & \rho_{123} = \langle a^\dagger a^\dagger a^\dagger a a a \rangle & \dots \\ \text{one-body density} & \text{two-body} & \text{three-body} & \text{N-body} \end{array}$$

To determine the dynamics exactly:

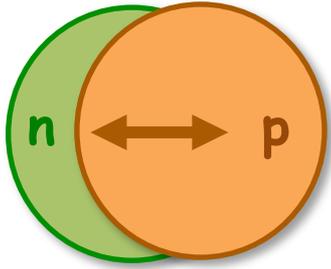
$$\rho_1 = \langle a^\dagger a \rangle \quad \rho_{12} = \langle a^\dagger a^\dagger a a \rangle \quad \rho_{123} = \langle a^\dagger a^\dagger a^\dagger a a a \rangle \quad \dots$$

one-body density two-body three-body N-body

$$\begin{aligned} i\dot{\rho}_1 &= [t_1, \rho_1] + \text{Tr}_{(2)} \{ [v_{12}, \rho_{12}] \} \\ i\dot{\rho}_{12} &= [t_1 + t_2 + v_{12}, \rho_{12}] + \text{Tr}_{(3)} \{ [v_{13} + v_{23}, \rho_{123}] \} \\ &\vdots \\ i\dot{\rho}_{1\dots n} &= \left[\sum_{i=1}^n t_i + \sum_{j>i=1}^n v_{ij}, \rho_{1\dots n} \right] \\ &\quad + \sum_{i=1}^n \text{Tr}_{(n+1)} \{ [v_{i(n+1)}, \rho_{1\dots(n+1)}] \} \end{aligned} \quad \begin{array}{l} \text{H} = \text{t} + \text{v} \\ \text{Hamiltonian} \end{array}$$

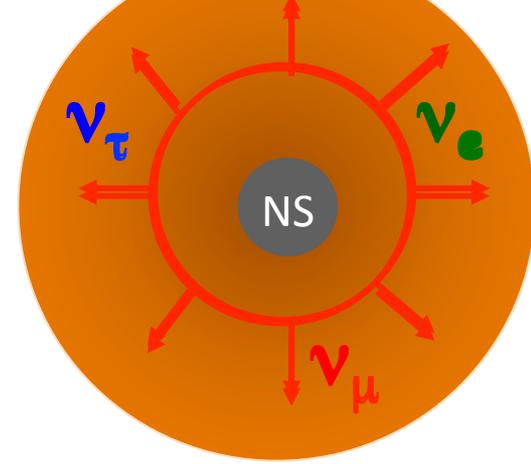
Born-Bogoliubov-Green-Kirkwood-Yvon (BBGKY) hierarchy

an infinite set of equations for a relativistic system



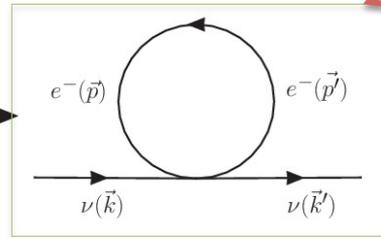
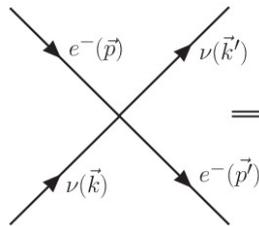
To determine the dynamics

$$\rho = \langle a^\dagger a \rangle \quad \text{one-body density matrix}$$



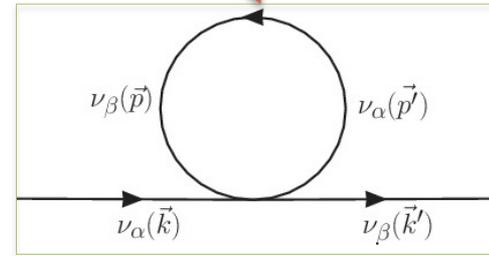
$$i\dot{\rho} = [h(\rho), \rho]$$

$$h(\rho) = h_0 + h_{\text{mat}} + h_{\nu\nu}(\rho)$$



neutrino-matter

$$h_{\text{mat}} = \sqrt{2}G_F \rho_e$$



neutrino self-interactions

non-linear term

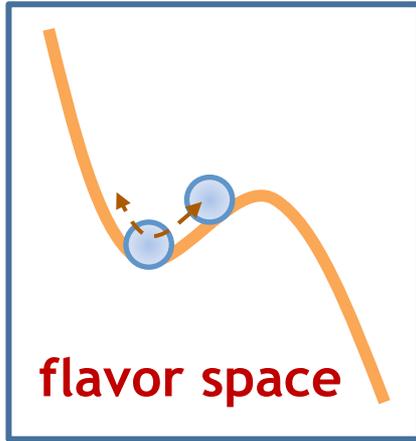
MEAN-FIELD approximation

BBGKY hierarchy : mean-field and beyond

Volpe, Väänänen, Espinoza. PRD 87 (2013)

- see Volpe, «Neutrino quantum kinetic equations », Int. J. Mod. Phys.E24(2015)

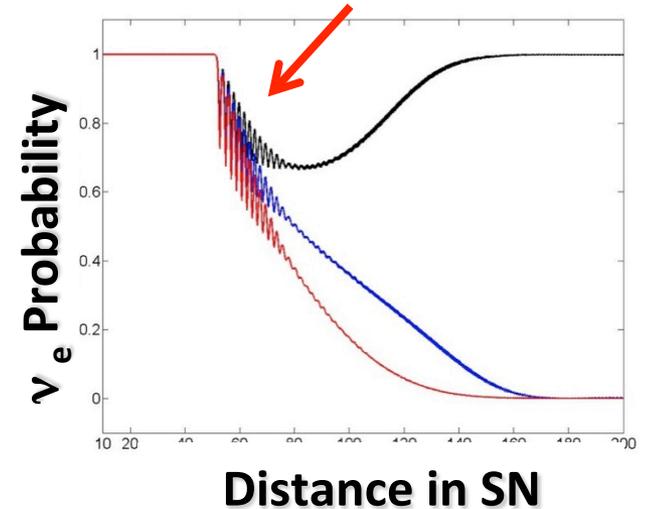
Collective neutrino modes and linearization



Small amplitude motion

Collective modes and instabilities can be studied with the linearization.

Banerjee, Dighe, Raffelt, PRD84 (2011)



Stability matrix

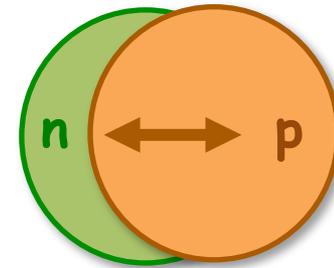
$$\begin{pmatrix} A & B \\ \bar{B} & \bar{A} \end{pmatrix} \begin{pmatrix} \rho' \\ \bar{\rho}' \end{pmatrix} = \omega \begin{pmatrix} \rho' \\ \bar{\rho}' \end{pmatrix}$$

S eigenvalues :

- > real : **stable collective**
- > imaginary : **instabilities**

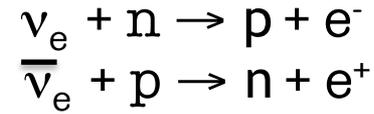
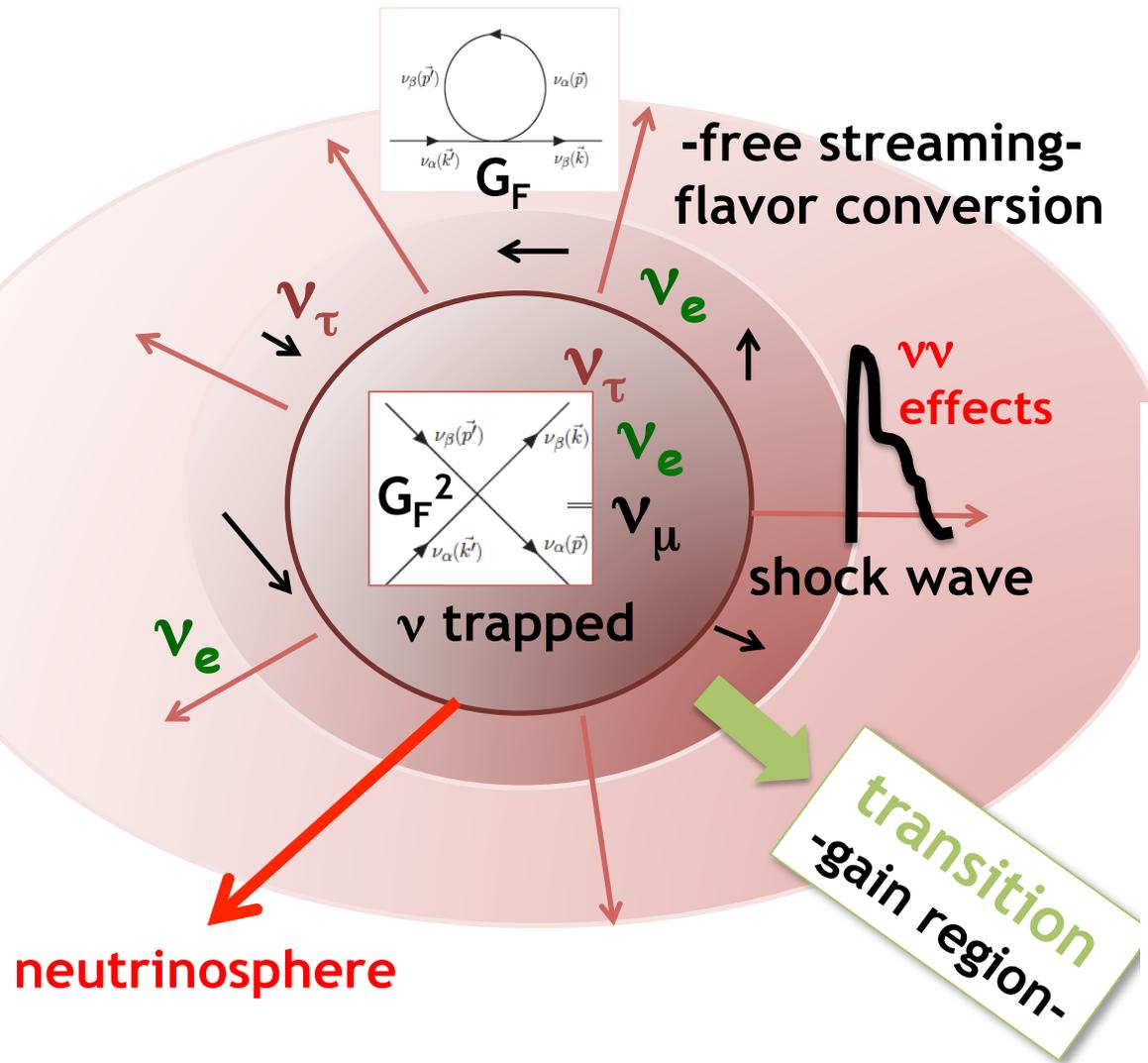
connection to collective modes in other many-body systems (nuclei, clusters, ...)

Väänänen and Volpe, PRD88 (2013)

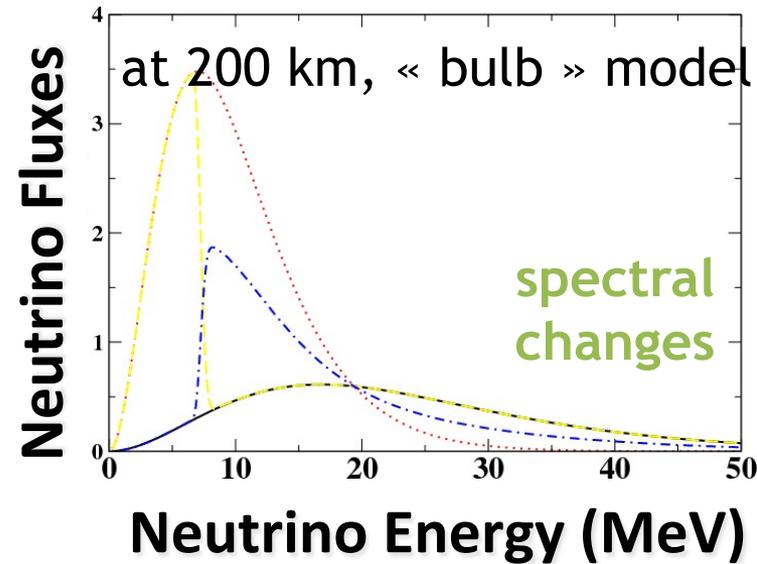


nuclear resonances

In supernovae explosions



Heating rate behind the shock could be enhanced by spectral changes of the neutrino fluxes.



Numerous aspects investigated in a decade...

❖ Symmetry breaking

New instabilities found, broken symmetries.

see e.g. Raffelt, Sarikas, Seixas, PRL 111 (2013)

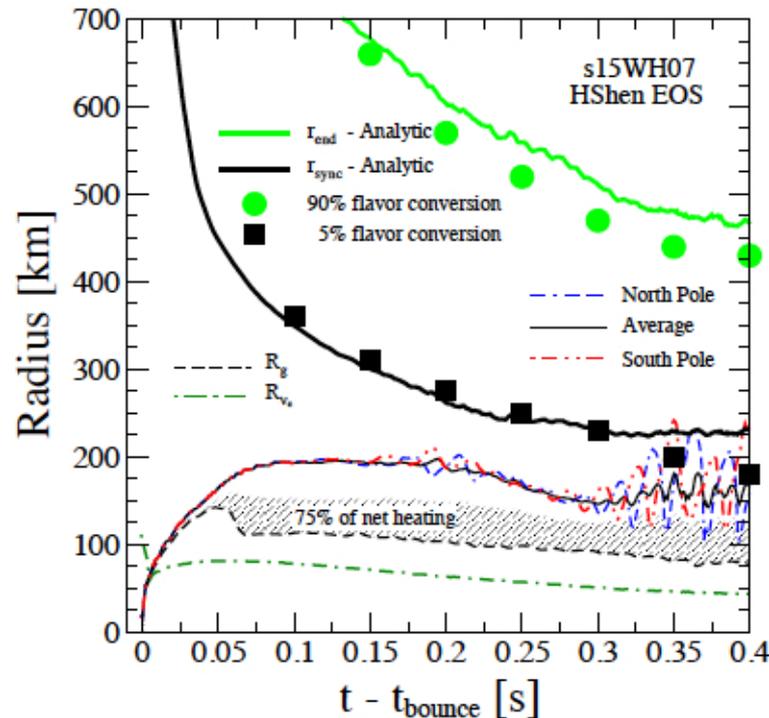
❖ Role of decoherence

❖ Impact on nucleosynthesis

Balantekin and Yuksel, Journ. Phys. 2005, Duan, McLaughlin, Surman JPG 2010, Meng-Ru et al, PRD91 (2015), ...

❖ Impact on the supernova dynamics

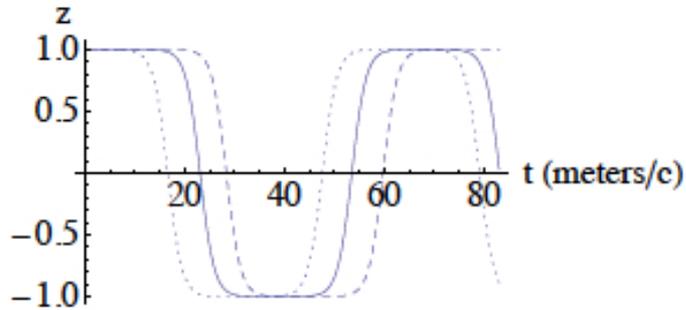
see the talk by H. Sasaki



Dasgupta et al. PRD 2012

Improved description of the transition region

- ❖ if different neutrinospheres considered

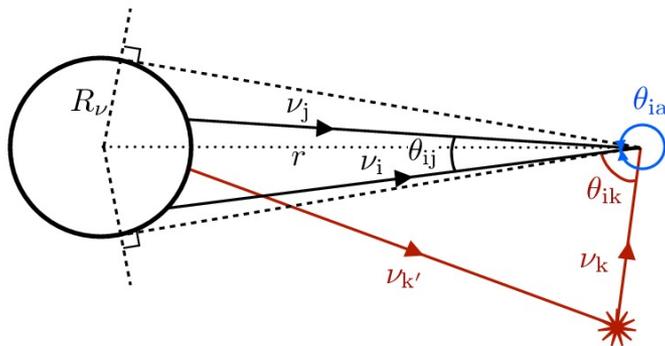


Conversion at short time scales

Sawyer, PRL108 (2016)

- ❖ collisions \rightarrow flavor patterns modified (schematic evaluation)

Cherry et al, PRL108 (2012),



Competition between the collision and flavor timescales ?

Improved description of the transition region

❖ mass terms contributions – corrections at the mean-field level

$\xi = \langle a_+^+ a_- \rangle$ correlators with helicity change

$$\mathcal{R} = \begin{pmatrix} \rho & \xi \\ \xi^* & \bar{\rho} \end{pmatrix} \quad \mathcal{H} = \begin{pmatrix} h & \Phi \\ \Phi^* & \bar{h} \end{pmatrix}$$

\mathcal{R} and \mathcal{H} have helicity and flavor structure ($2\mathcal{N}_f \times 2\mathcal{N}_f$).

Φ couples ν with $\bar{\nu}$
helicity (or spin) coherence

$$\Phi \sim (h_{\text{mat}}^{\text{perp}} + h_{\nu\nu}^{\text{perp}}) \times m/2E$$

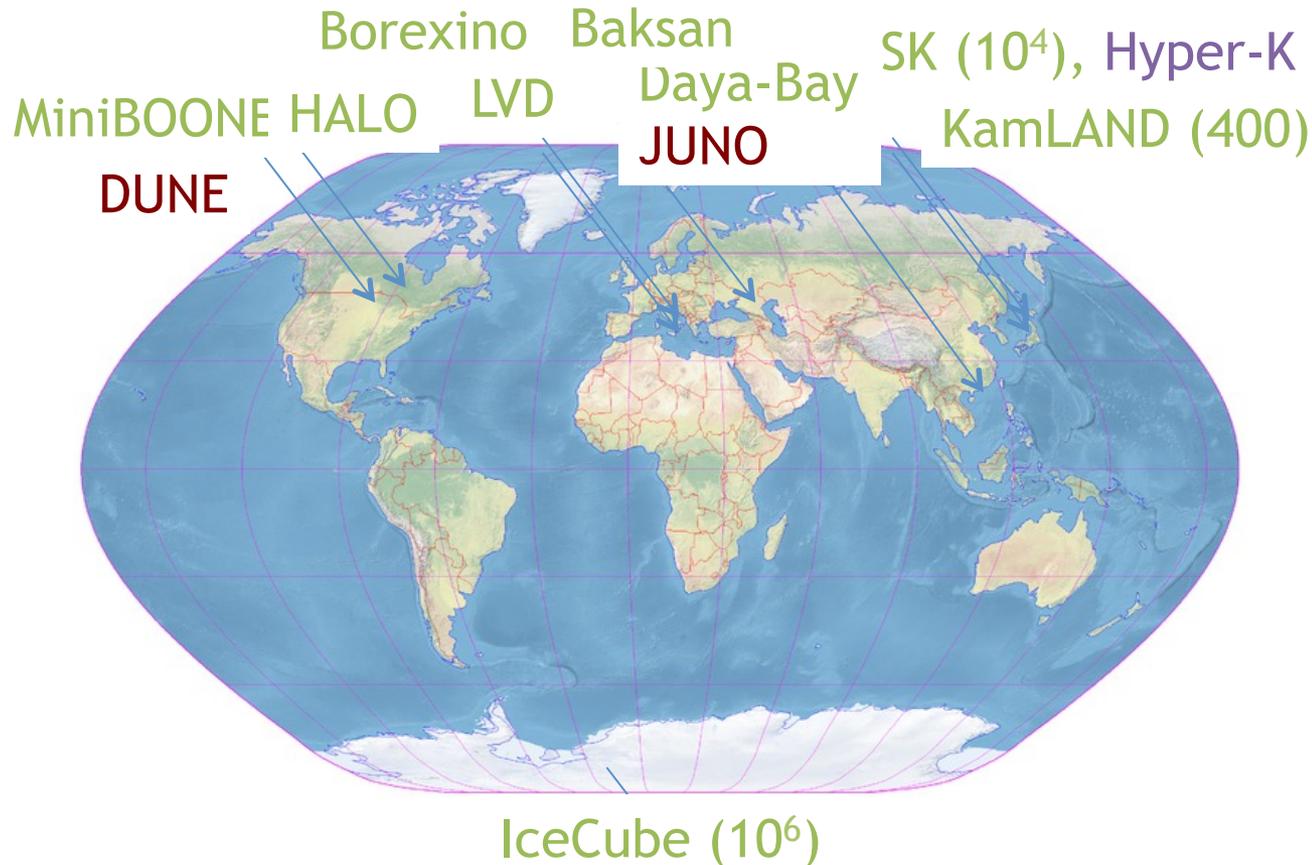
Vlasenko, Fuller, Cirigliano,
PRD89 (2014)

Serreau, Volpe, PRD90 (2014)

First calculation in a toy model has shown significant impact.

Supernova Early Warning System and SNe observatories

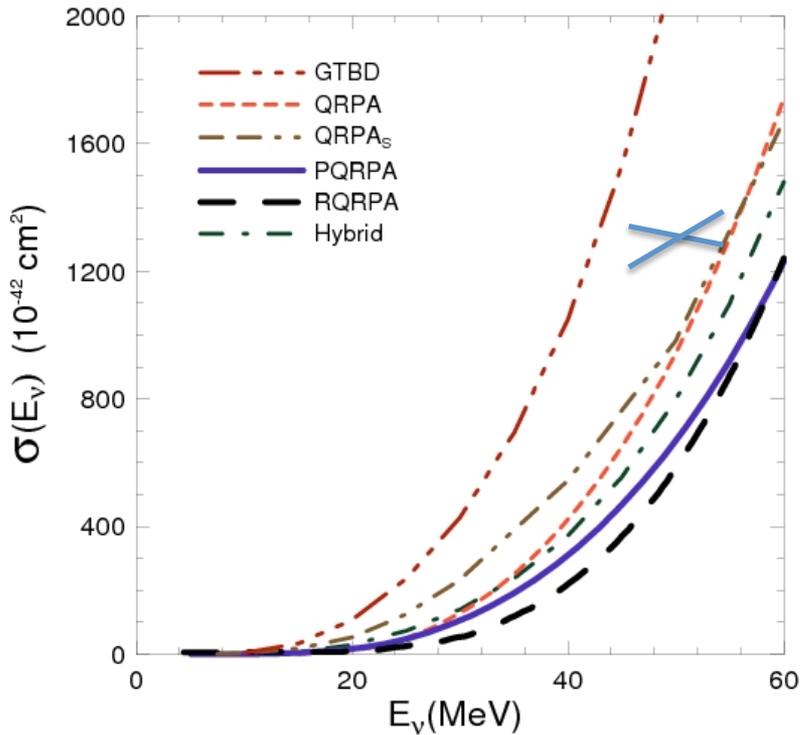
Events for a supernova explodes in our galaxy (10 kpc), up to 10^6 events



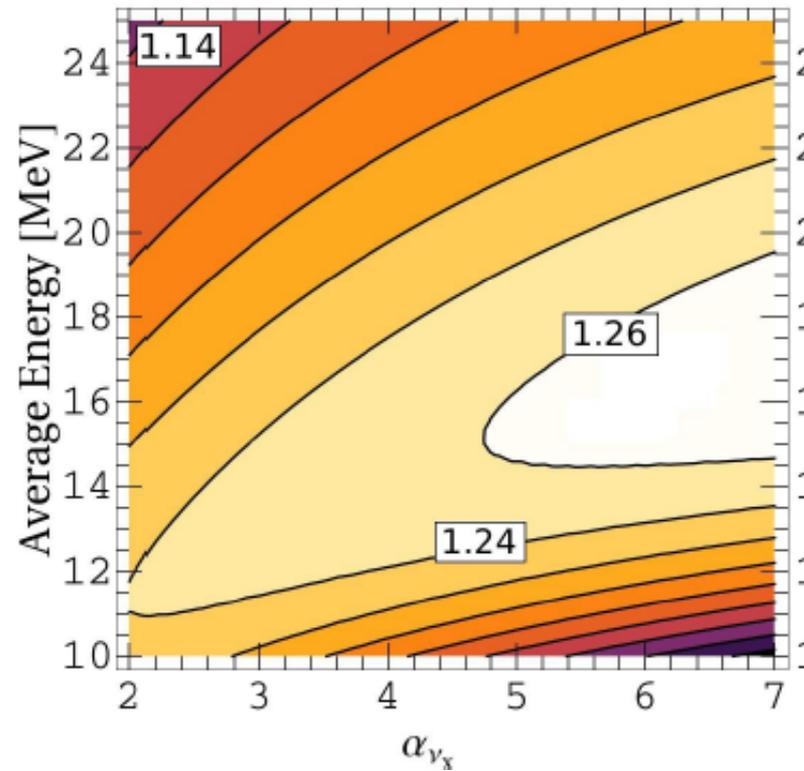
Detection channels : scattering of anti- ν_e with p , ν_e with nuclei, ν_x with e , p

We will measure the time, energy signal from future (extra)galactic explosions.

Neutrino-nucleus cross sections



Bertulani, Samana, 0802.1553



Vaananen, Volpe, JCAP 1011 (2011)

Several projects have been proposed over the years... ORLAND at SNS (2000), low energy beta-beam (Volpe, J. Phys. G30, 2003), ν at JPARC (H.Ejiri, 2003).

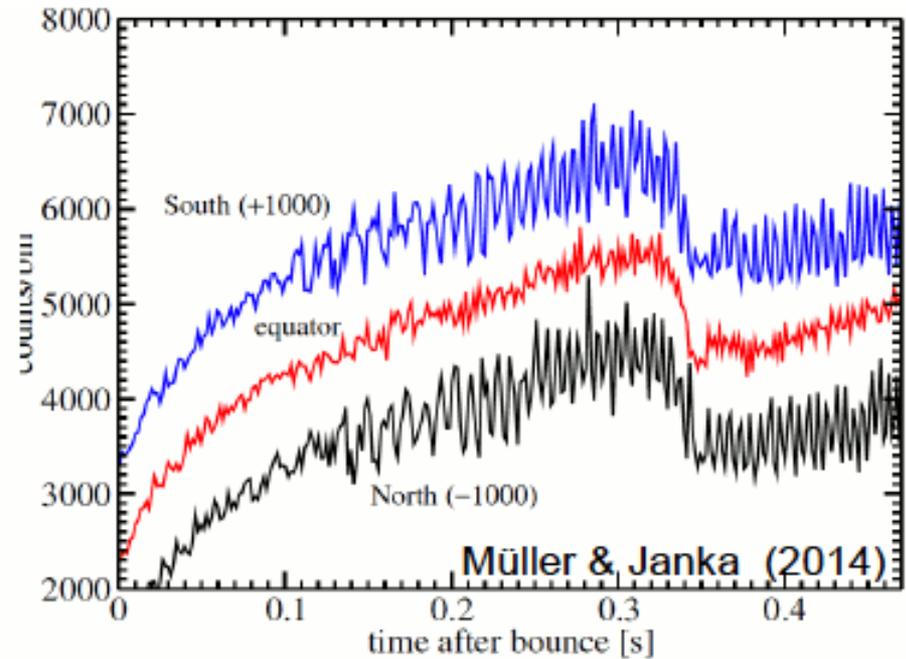
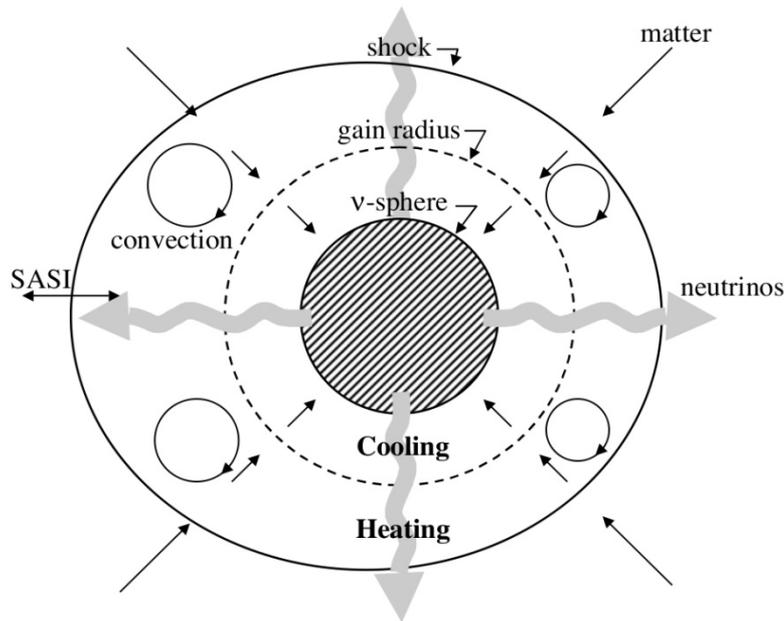
Measurements on ν -nucleus cross sections planned at SNS by the COHERENT collaboration ! (Pb, O, Ar, Fe)

Sensitivity to the quenching of g_A , also of forbidden transitions

The explosion mechanism and the neutrino signal

Now : Two-dimensional simulations agree for a set supernova progenitors, from three groups, accord with SN1987A explosion energy. Three-dimensional running, hints for explosions.

A. Mezzacappa's talk to «8th TPC Symposium », Paris



Time supernova neutrino signals encode the SASI signature (ICECUBE).

Gravitational binding energy of the newly formed neutron star

How well can we reconstruct the gravitational binding energy in a galactic explosion ?

Most of the analysis make assumptions - ex. equipartition hypothesis, or pinching parameter fixed.

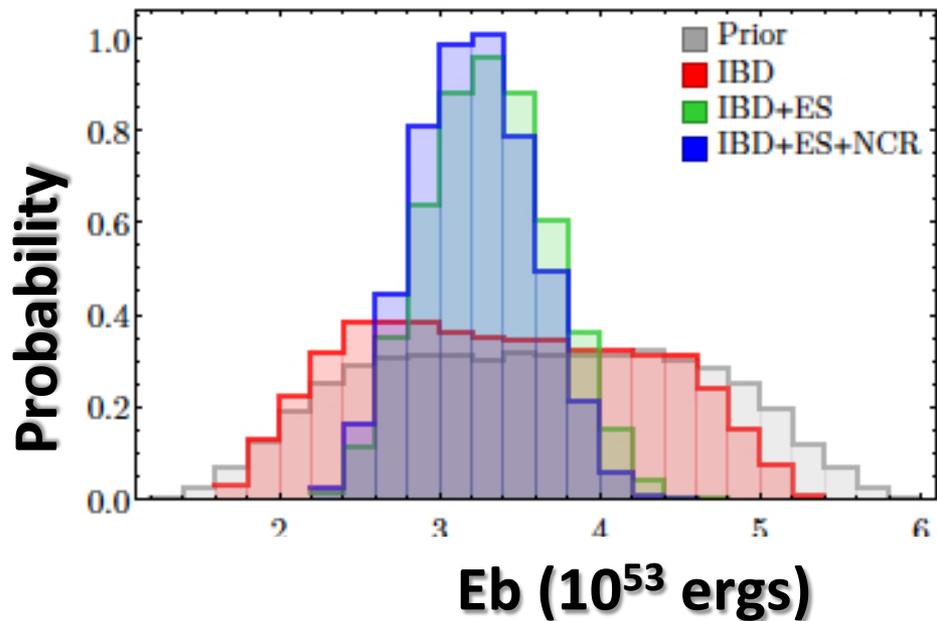
An et al. J. Phys. G43 (2016), Lu, Li, Zhou, PRD94 (2016) 023006, ...

Minakata et al, arXiv:0802.1489, only inverse beta-decay in Hyper-Kamiokande. Conclusions not optimistic...

Reconstructing the gravitational binding energy of the neutron star

Gallo Rosso, Vissani, Volpe, JCAP1711 (2017)
 Gallo Rosso, Vissani, Volpe, arXiv:1712.05584

For a galactic supernova at 10 kpc.



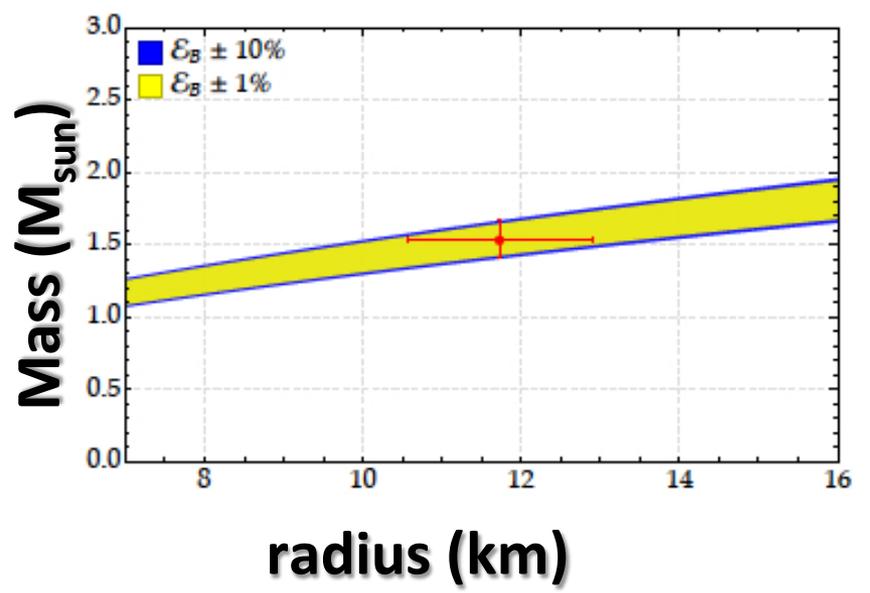
- ✓ Likelihood without any ansatz
- ✓ Fluence : a power-law, MSW, NH
- ✓ Combined inverse beta-decay, elastic scattering and neutral current on oxygen

**E_b reconstructed with 11% accuracy in Super-Kamiokande
 3% in Hyper-Kamiokande**

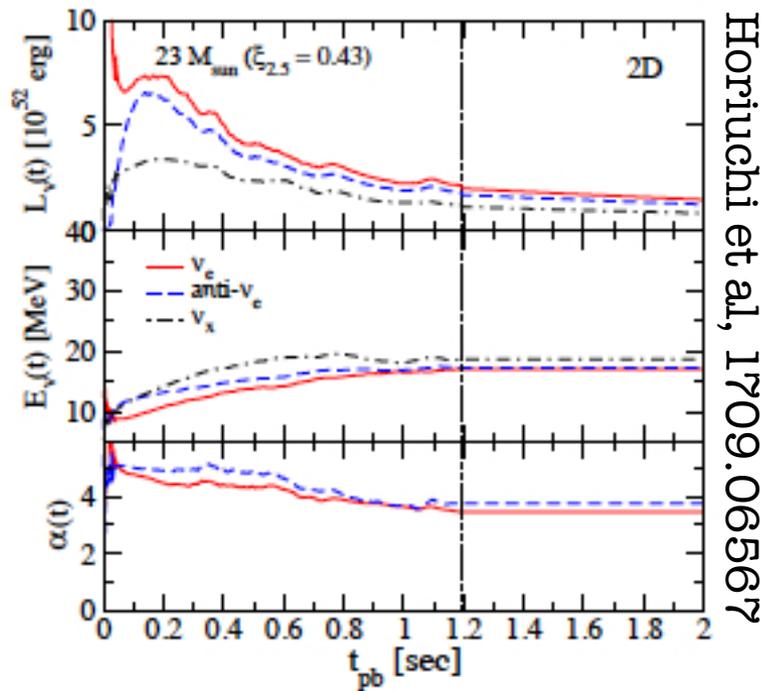
Fit to numerous EOS for NS

$$\frac{\mathcal{E}_B}{Mc^2} \approx \frac{(0.60 \pm 0.05) \beta}{1 - \beta/2}, \quad \beta = \frac{GM}{Rc^2}$$

From Lattimer, Prakash,



The diffuse supernova neutrino background



The integrated ν neutrino flux from supernovae at different redshifts :

$$F_{\alpha}(E_{\nu}) = \int dz \left| \frac{dt}{dz} \right| (1+z) R_{\text{SN}}(z) \frac{dN_{\alpha}(E'_{\nu})}{dE'_{\nu}},$$

$$E'_{\nu} = (1+z)E_{\nu},$$

Upper limits on DSNB fluxes :

1.4-1.9 anti- n_e /cm²/s

73-154 n_e /cm²/s

Lunardini and Peres, JCAP (2008)

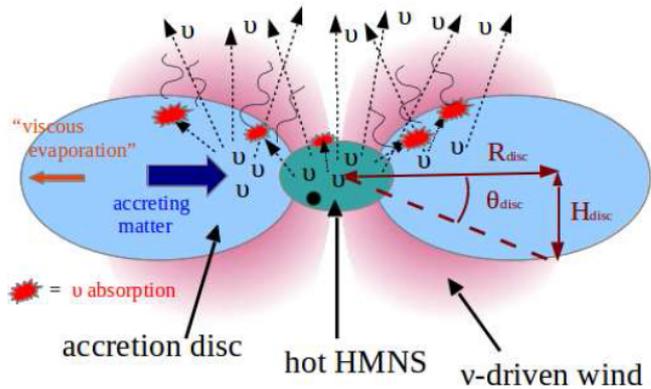
	Events (10 y)	window	detector
$\bar{\nu}_e$	90	9-25 MeV	50 kton scintillator
$\bar{\nu}_e$	300	19-30 MeV	440 kton Hyper-K
ν_e	30	17-41 MeV	50 kton liquid argon

Galais et al PRD 81(2010)

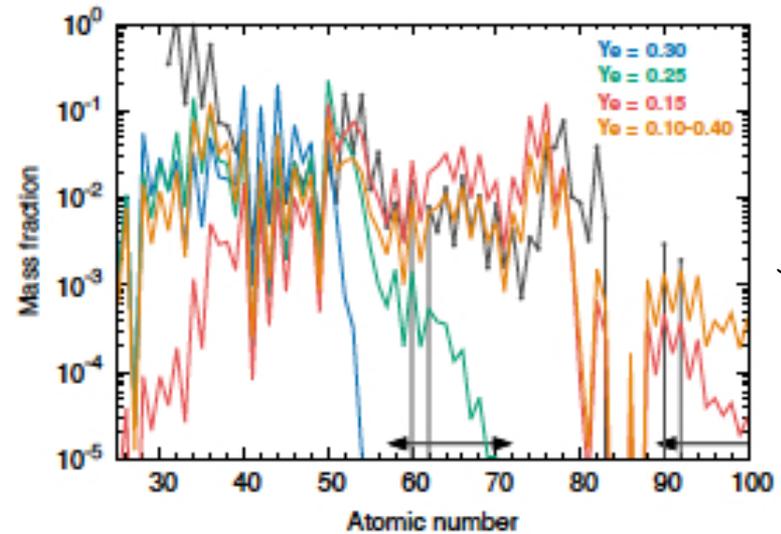
EGADS -Super-Kamiokande with Gd

DSNB discovery in the coming years

Flavor evolution in accretion disks around compact objects and kilonova observations



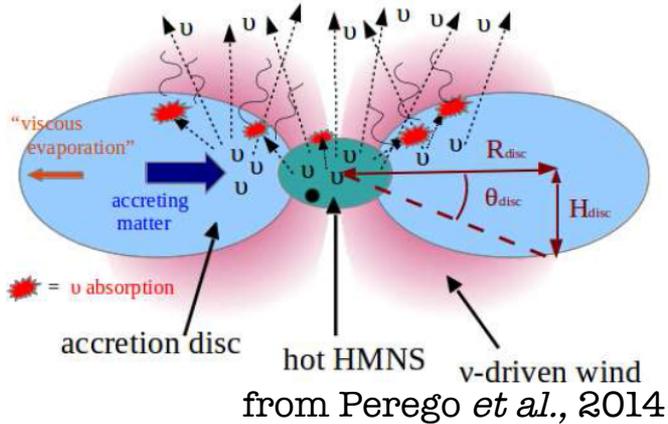
from Perego *et al.*, 2014



Tanaka et al., 1710.05850

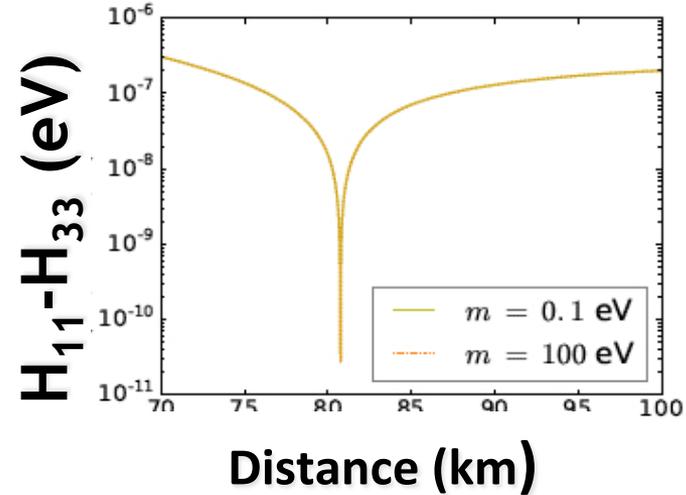
Helicity coherence (mass terms) effects

neutron star mergers remnant



$$\mathcal{H} = \begin{pmatrix} h & \Phi \\ \Phi^* & \bar{h} \end{pmatrix} \quad \text{In two flavors}$$

\mathcal{H} is a 4 x 4 matrix.

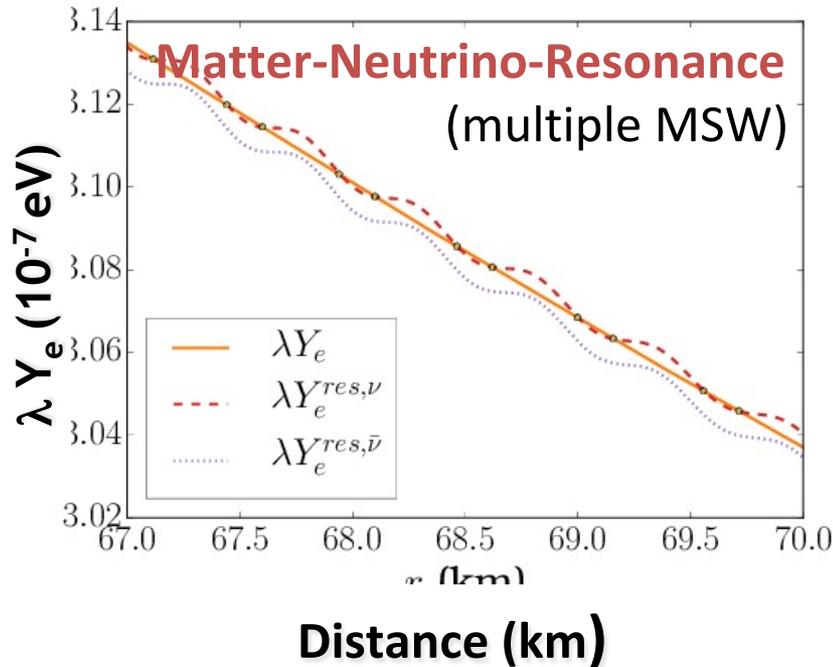


- resonance condition fulfilled
- adiabaticity not sufficient

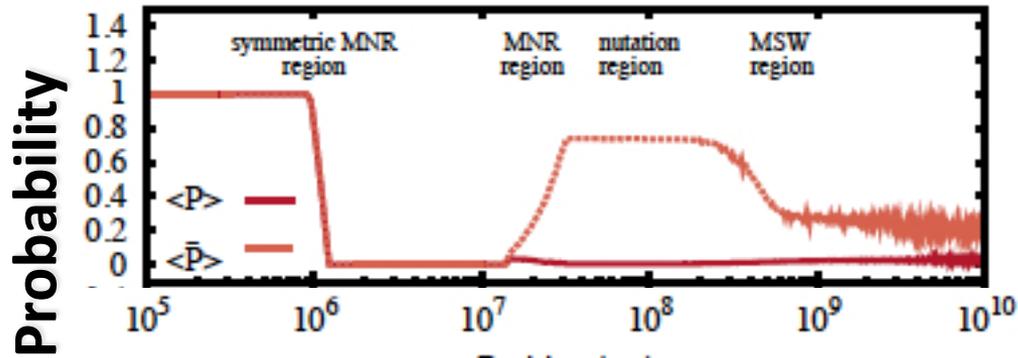
Non-linear feedback does not enhance mass effects.

Chatelain, Volpe, PRD(2017)
1611.01862

contrary to the findings in Vlasenko, Fuller, Cirigliano, 1406.6724

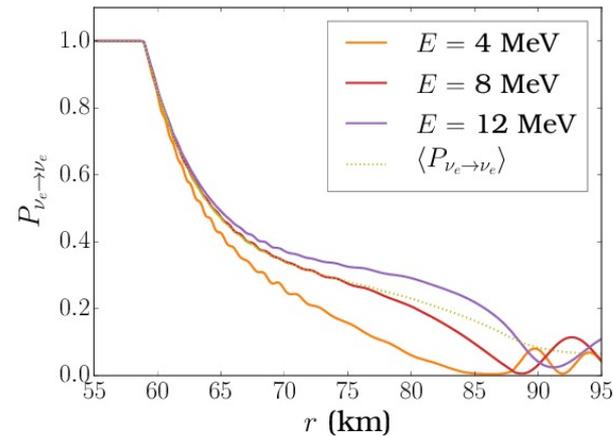
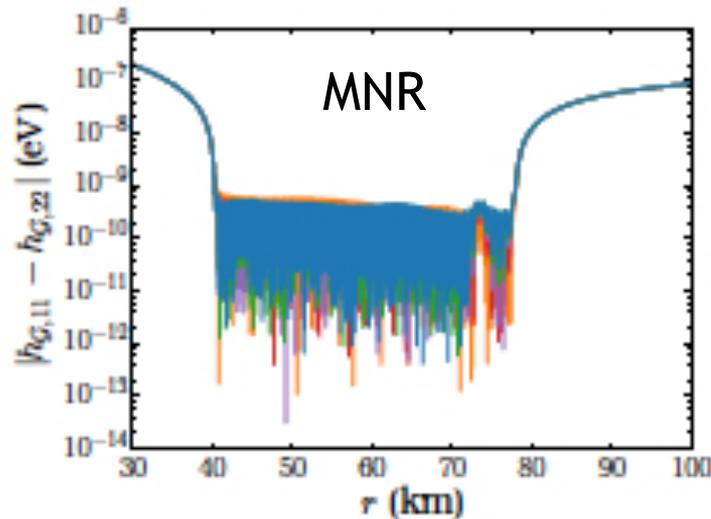


Flavor phenomena in accretion disks around compact objects



Excess of $\bar{\nu}_e$ over ν_e
 The self-interaction and matter potentials cancel.

Malkus et al, PRD86 (2012)

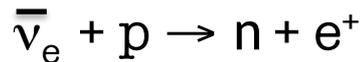
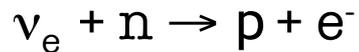


Chatelain, Volpe, PRD(2017)
 1611.01862

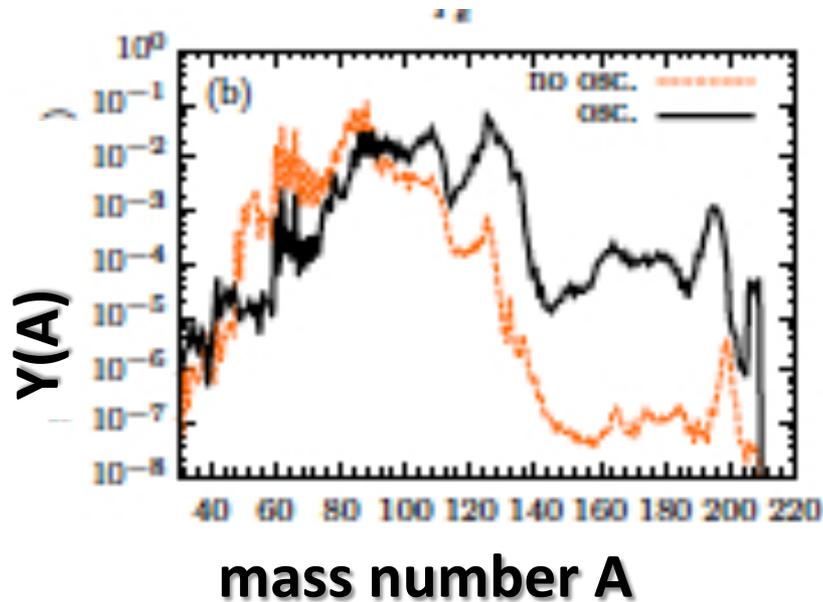
Resonance conditions met, adiabatic evolution,
 ν_e can be modified by flavor evolution

Nucleosynthesis in neutrino-driven winds and kilonovae

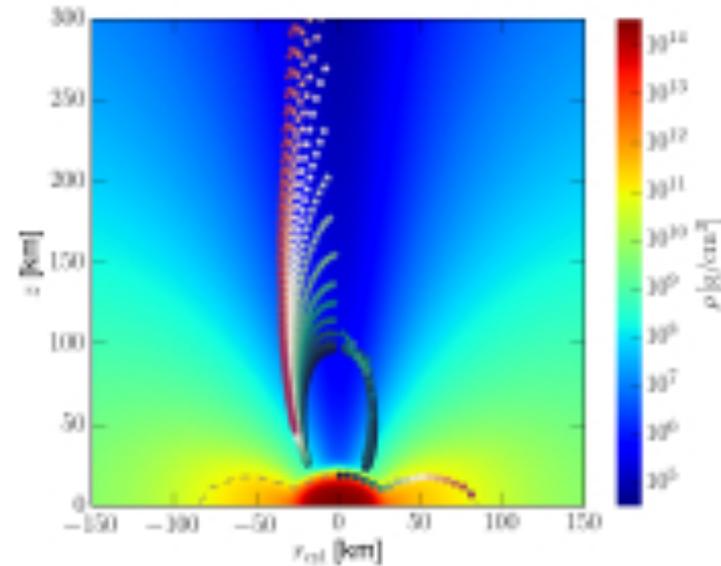
Neutrinos influence the neutron richness and determine Y_e



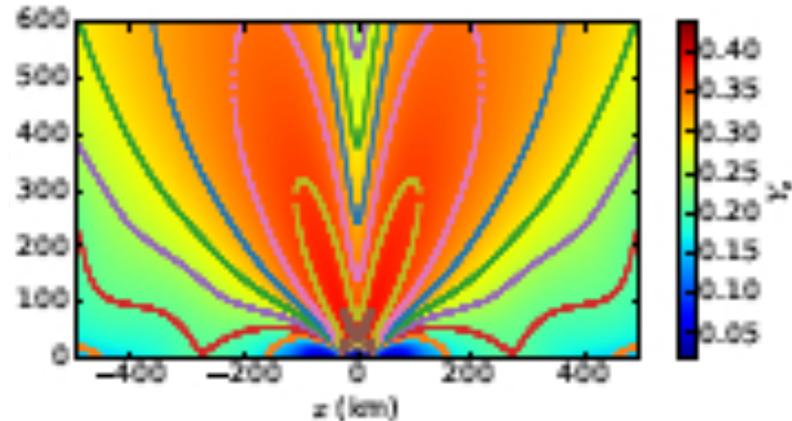
Flavor evolution can impact Y_e .



Wu et al., PRD96(2017)



Frensel et al., PRD95(2017)



Chatelain and Volpe, PRD98 (2018)

Kilonova observations - assess the role of flavor evolution on Y_e and nucleosynthesis in neutrino-driven winds

Conclusions and perspectives



Neutrino flavor evolution in dense astrophysical environments is a complex problem that still needs to be fully understood.

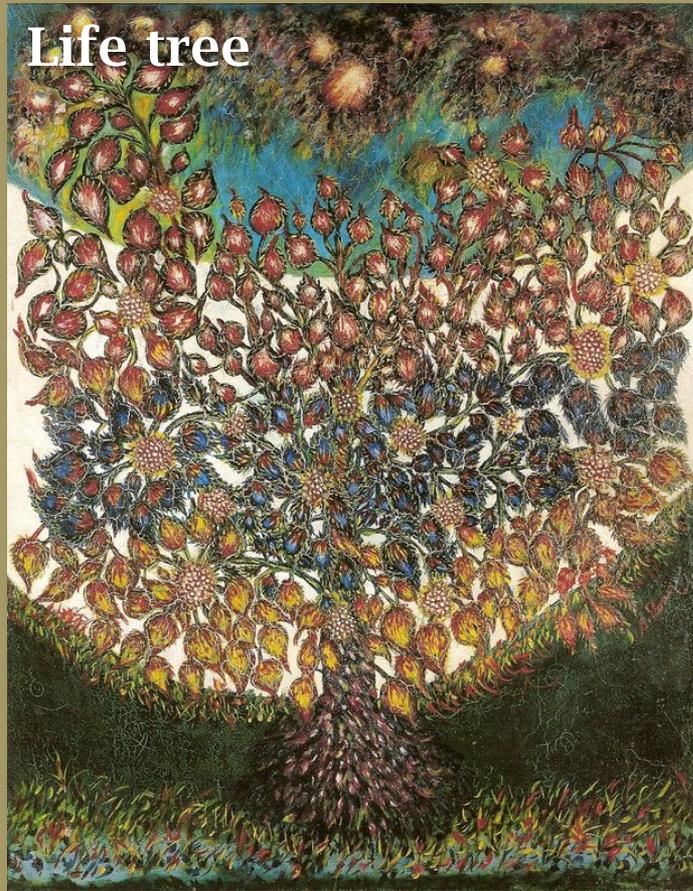


Fast modes bring equilibration of the spectra on short distance scales. They could influence the supernova dynamics. These and other phenomena could impact nucleosynthesis in neutrino-driven winds in kilonovae.



The total gravitational binding energy will be precisely reconstructed for a galactic supernova and also the neutrino spectra, if effects from the flavor modification simple enough (MSW or fast modes).

Life tree



Seraphine de Senlis